### Anisotropic flow fluctuation as a possible signature of clustered nuclear geometry in O-O collisions at the Large Hadron Collider

G.G. Barnaföldi, N. Mallick, S. Prasad, R. Sahoo

Support: Hungarian OTKA grants, K135515, NEMZ\_KI-2022-00031, 2024-1.2.5-TÉT-2024-00022, Wigner Scientific Computing Laboratory Ref.: Physics Letters B 860 (2025) 139145

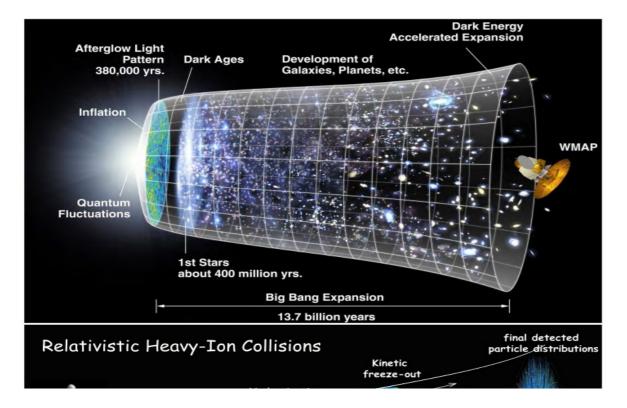
Triggering Discoveries in HEP Vysoke Tatry, Slovakia, 10<sup>th</sup> December 2024

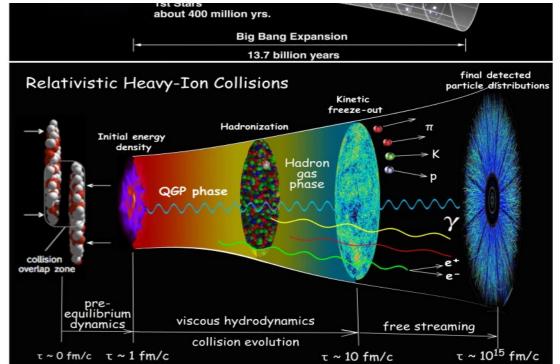


### Motivation & definitions

### Primordial matter in heavy-ion collisions

• Quark-Gluon Plasma (QGP) research



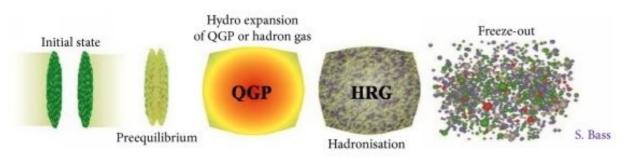


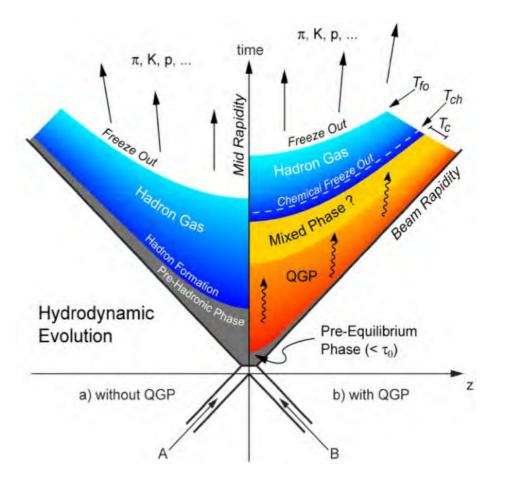
G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Primordial matter in heavy-ion collisions

#### • QGP in experimental vs theory points

- By colliding heavy-ions we can form small drop of the hot & dense primordial matter
- No direct observations, just signatures: jet-quenching, correlations, collective effects, (anisotropic) flow...
- Need a complex description, including QCD phenomenology, hydrodynamics, (non-equilibrium) thermodynamics

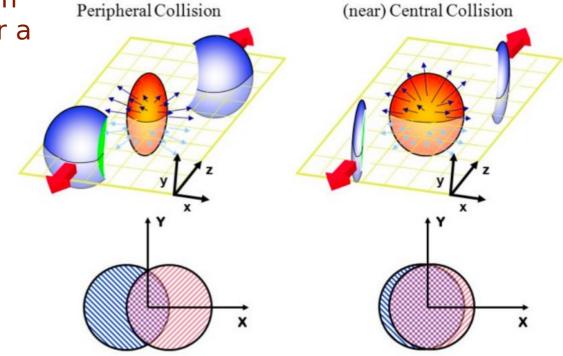




G.G. Barnafoldi: Triggering HEP, Tatry 2024

#### • Experimental point:

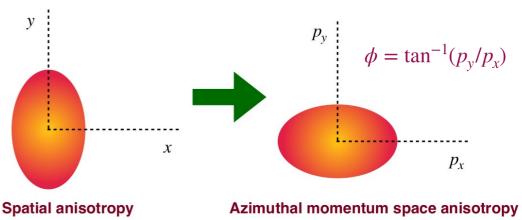
 Flow describes the azimuthal momentum space anisotropy of particle emission for a non-central heavy-ion collision.



#### • Experimental point:

- Flow describes the azimuthal momentum space anisotropy of particle emission for a non-central heavy-ion collision.
- The n<sup>th</sup> harmonic coefficient of the Fourier expansion of azimuthal momentum distribution:

$$E\frac{d^{3}N}{dp^{3}} = \frac{d^{2}N}{p_{\rm T}dp_{\rm T}dy}\frac{1}{2\pi} \left(1 + 2\sum_{n=1}^{\infty}v_{n}\cos[n(\phi - \psi_{n})]\right)$$

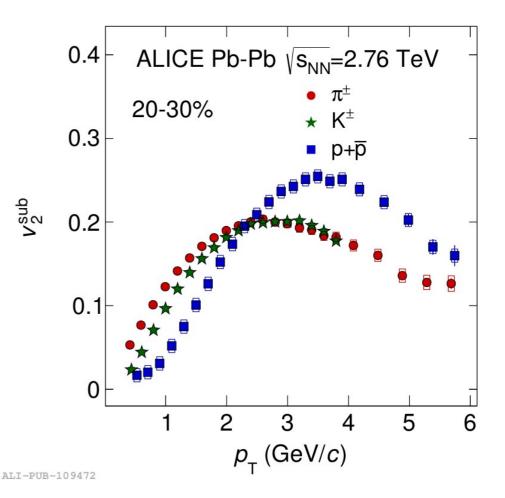


#### • Experimental point:

- Flow describes the azimuthal momentum space anisotropy of particle emission for a non-central heavy-ion collision.
- The n<sup>th</sup> harmonic coefficient of the Fourier expansion of azimuthal momentum distribution:

$$E\frac{d^{3}N}{dp^{3}} = \frac{d^{2}N}{p_{\mathrm{T}}dp_{\mathrm{T}}dy}\frac{1}{2\pi}\left(1 + 2\sum_{n=1}^{\infty}v_{n}\cos[n(\phi - \psi_{n})]\right)$$

- The  $v_2(p_T, y) = \langle \cos(2(\phi - \psi_2)) \rangle$  directly reflects the initial spatial anisotropy of the nuclear overlap region in the transverse plane.



G.G. Barnafoldi: Triggering HEP, Tatry 2024

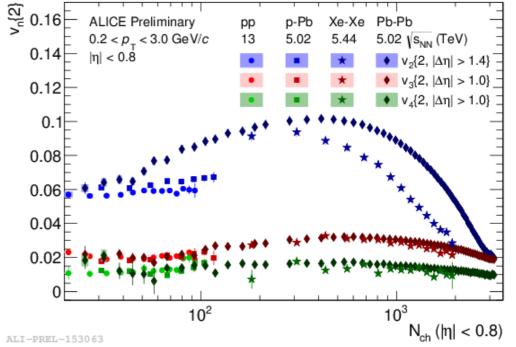
#### • Experimental point:

- Flow describes the azimuthal momentum space anisotropy of particle emission for a non-central heavy-ion collision.
- The n<sup>th</sup> harmonic coefficient of the Fourier expansion of azimuthal momentum distribution:

$$E\frac{d^{3}N}{dp^{3}} = \frac{d^{2}N}{p_{\mathrm{T}}dp_{\mathrm{T}}dy}\frac{1}{2\pi}\left(1 + 2\sum_{n=1}^{\infty}v_{n}\cos[n(\phi - \psi_{n})]\right)$$

- The  $v_2(p_T, y) = \langle \cos(2(\phi - \psi_2)) \rangle$  directly reflects the initial spatial anisotropy of the nuclear overlap region in the transverse plane.





### Future Nuclear Collisions at LHC

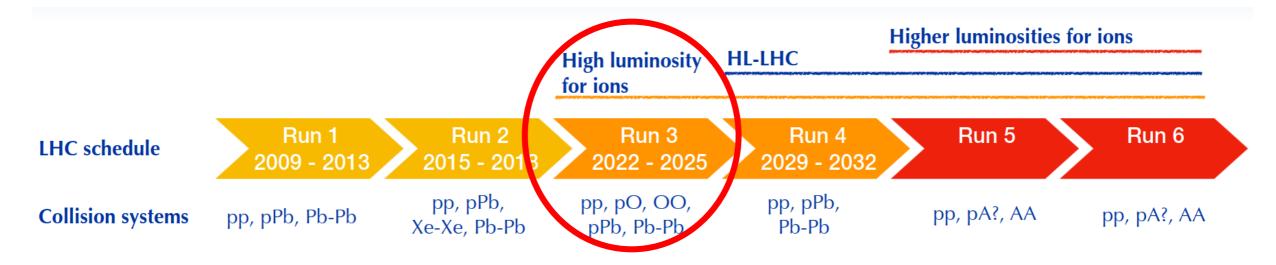
- LHC Schedule with new nuclear collisions
  - Run 2: XeXe
  - Run 3: pO & OO



### Future Nuclear Collisions at LHC

#### LHC Schedule with new nuclear collisions

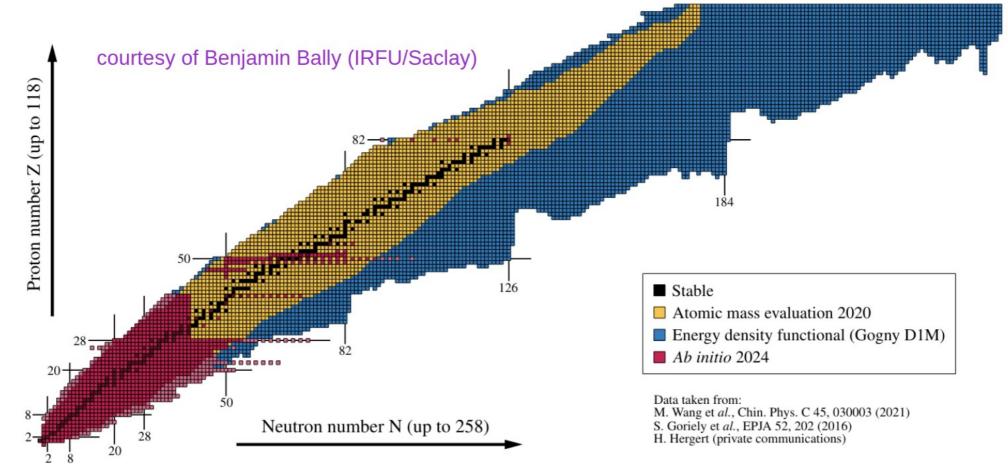
- Run 2: XeXe
- Run 3: pO & OO



### Nuclei & nuclear structure

### Nuclei for Future Nuclear Collisions

High-mass and deformed nuclei are in the focus:

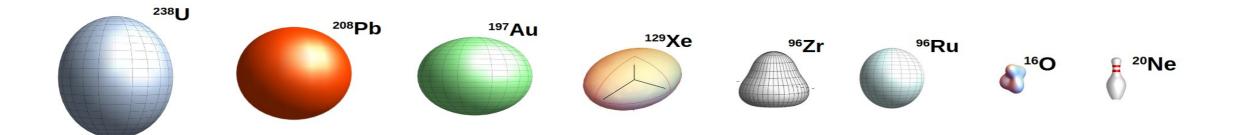


G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Nuclei for Future Nuclear Collisions

#### Experimental possibilities & interest

- Large deformed nuclei: uranium, gold, xenon
- Smaller zirconium, rubidium, oxygen, neon

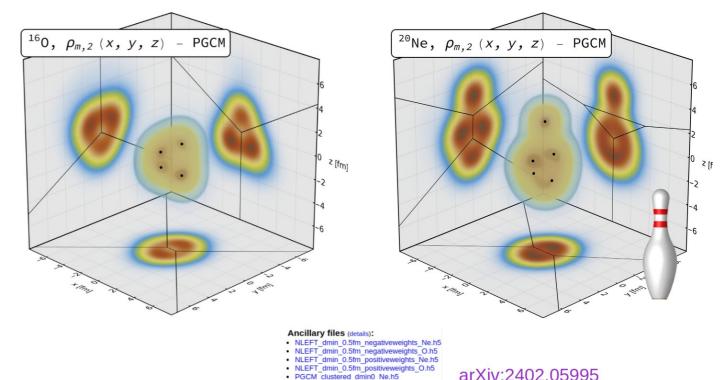


### Nuclei for Future Nuclear Collisions

#### Oxygen and Neon are unique

 Oxygen is a double magic nucleus, since both shells are closed shell. In cluster model Tetrahedron shape.

 Neon, has bowling pin shape, even more complicated geometry



PGCM\_clustered\_dmin0\_0.h5
PGCM\_uniform\_dmin0\_Ne.h5
PGCM\_uniform\_dmin0\_0.h5

### The shape of the oxygen

#### Modeling the oxygen

Woods-Saxon (WS)

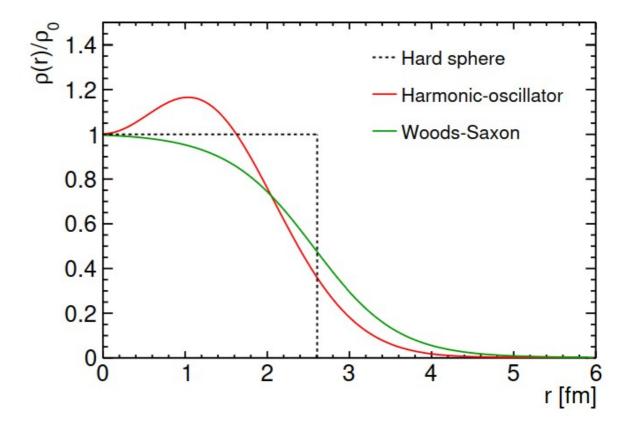
$$\rho(r) = \rho_0 \left[ 1 + \alpha \left( \frac{r}{a} \right)^2 \right] \exp \left( \frac{-r^2}{a^2} \right)$$

- Harmonic oscillator (HO)

$$\rho(r) = \frac{\rho_0 (1 + w(\frac{r}{r_0})^2)}{1 + \exp(\frac{r - r_0}{a})}$$

- Normalization:

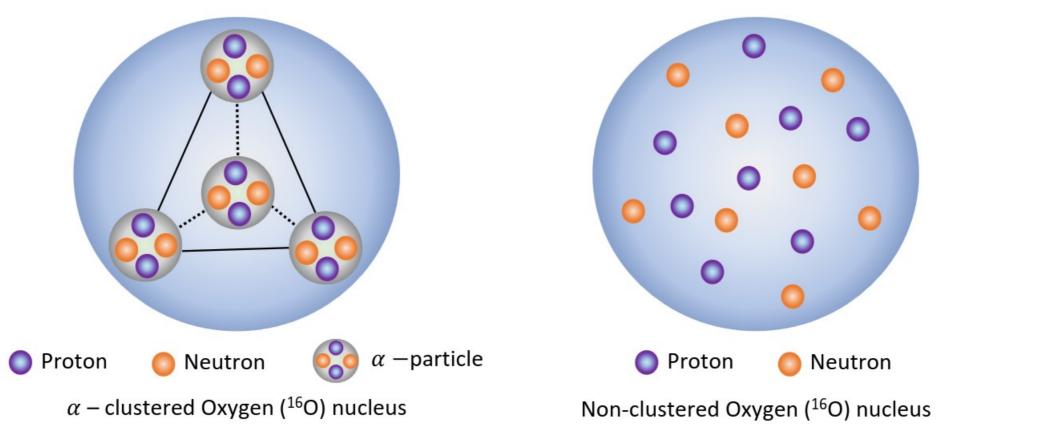
$$\int \rho(r)d^3r = 4\pi \int \rho(r)r^2dr = Ze$$



### The shape of the oxygen

#### **Nuclear structure description**

- Cluster model vs.



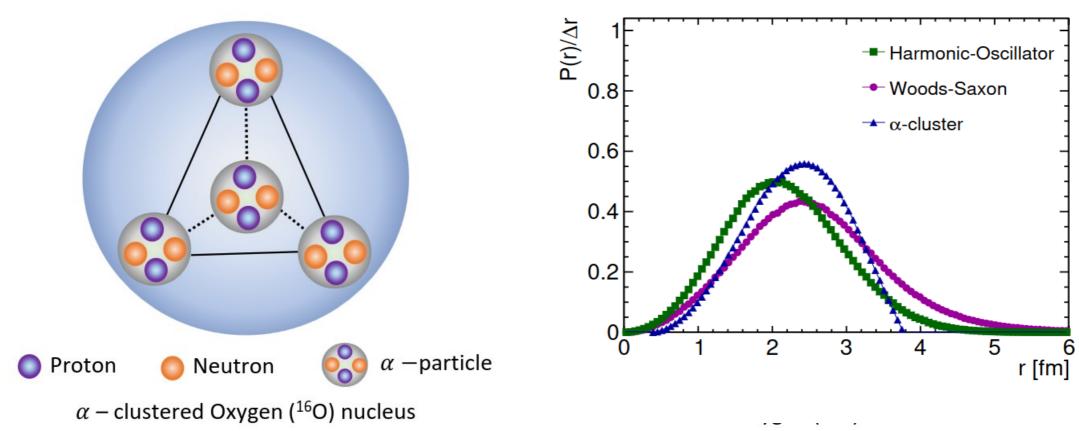
G.G. Barnafoldi: Triggering HEP, Tatry 2024

Non-cluster model (Woods-Saxon)

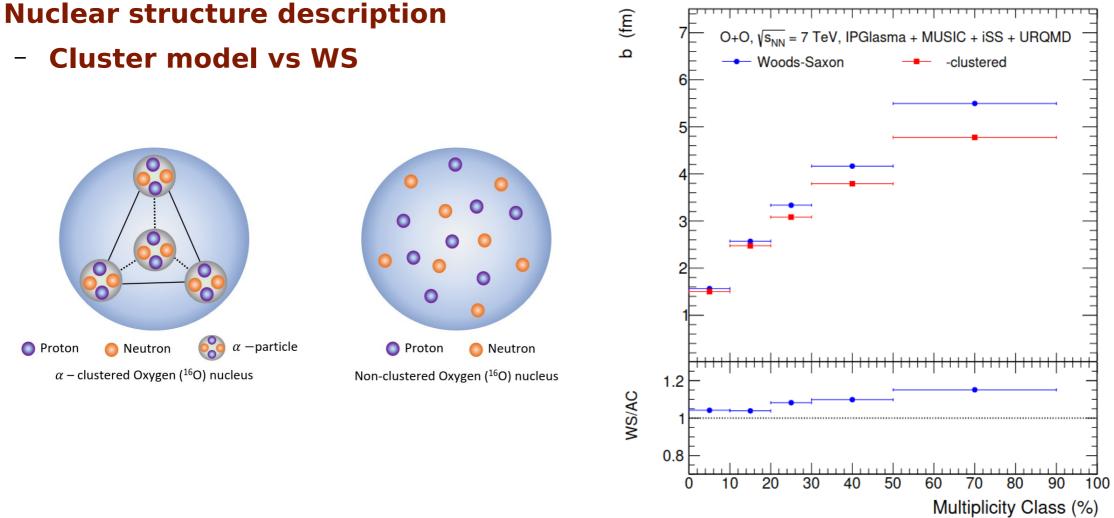
### The shape of the oxygen

#### **Nuclear structure description**

- Cluster model vs WS & HO



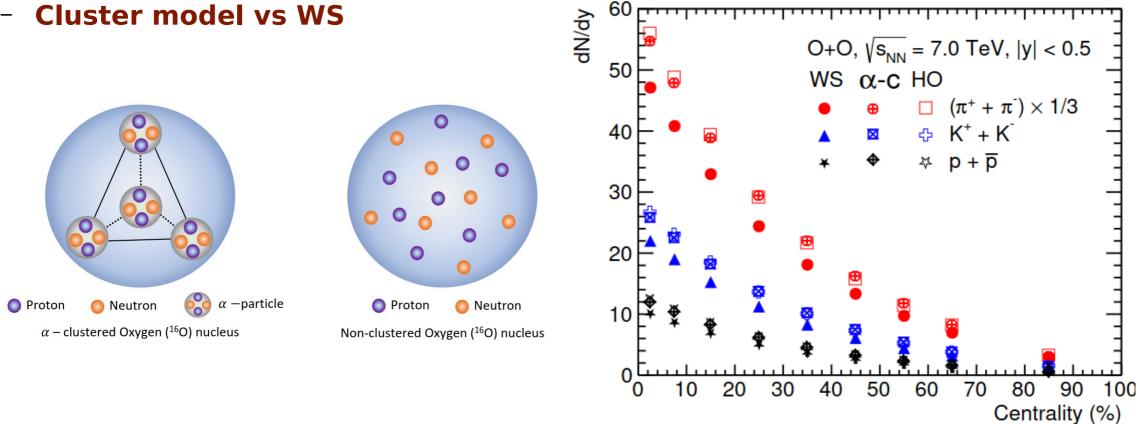
#### Probability of the radial position in O



G.G. Barnafoldi: Triggering HEP, Tatry 2024

#### **Nuclear structure description**

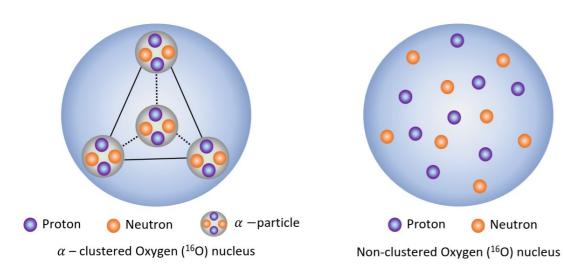
Cluster model vs WS



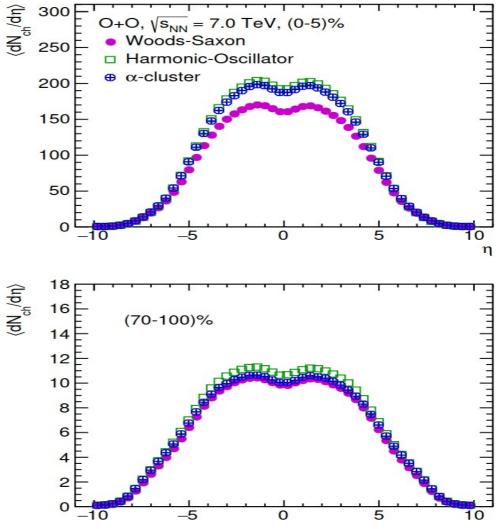
G.G. Barnafoldi: Triggering HEP, Tatry 2024

#### **Nuclear structure description**

- Cluster model vs WS

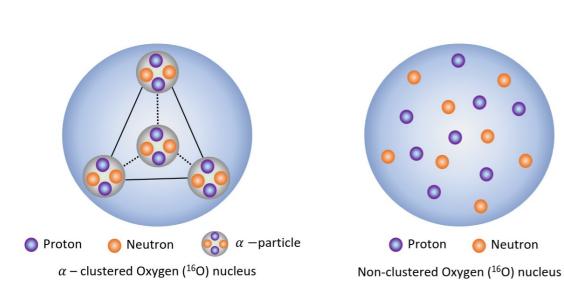


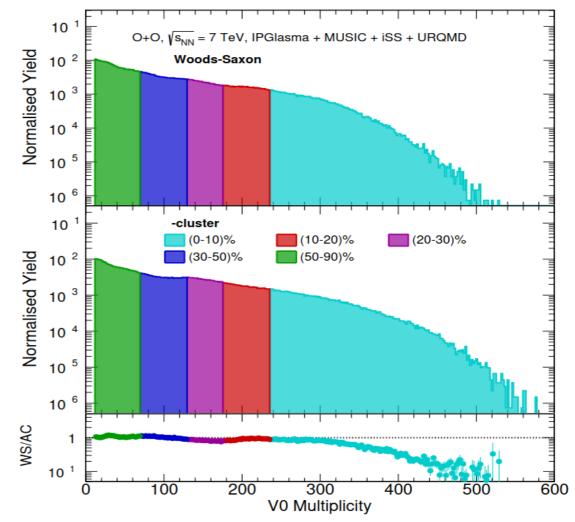
G.G. Barnafoldi: Triggering HEP, Tatry 2024



#### **Nuclear structure description**

- Cluster model vs WS





G.G. Barnafoldi: Triggering HEP, Tatry 2024

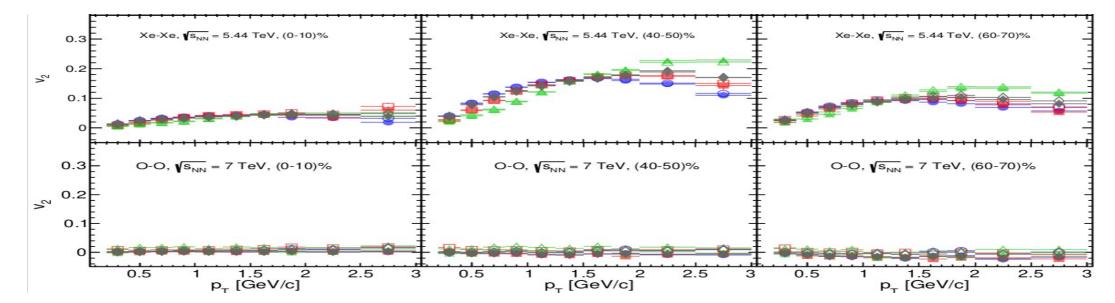
### Calculating the flow in small systems

G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Calculating the flow

#### **Event plane and average method**

- $v_n = \langle \cos[n(\phi \psi_n)] \rangle$
- Need to determine the event plain, which fails for small nuclei:



G.G. Barnafoldi: Triggering HEP, Tatry 2024

### The Model

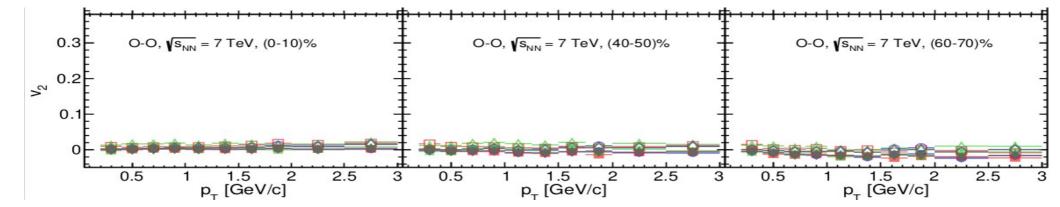
• A full hidro & Boltzmann transport with viscosity:

- Kinematical settings are:
  - Energy (c.m.): 7 TeV O+O
  - Pseudorapidity:  $|\eta| < 2.5$
  - Transverse momentum:  $0.2 < p_{\rm T} < 5.0 \ {\rm GeV/c}$
  - Pseudorapidity gap: ,  $|\Delta \eta| > 1.0$

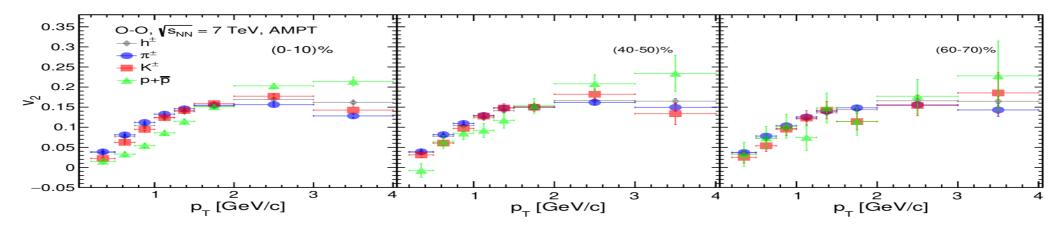
 $-\Delta \eta = 1 \rightarrow$   $A \qquad B$   $-2.5 \qquad -0.5 \quad 0 \quad 0.5 \qquad 2.5$   $Pseudorapidity (\eta) \rightarrow$ 

### Calculating the flow

#### **Event plane and average method**



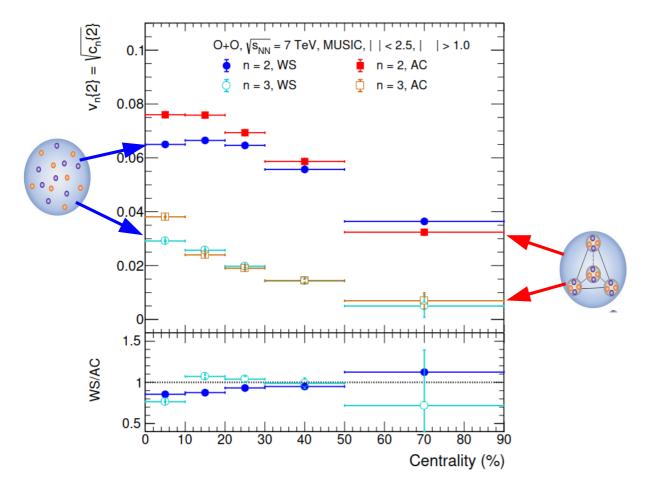
**Multiparticle Q-cummulant method** 



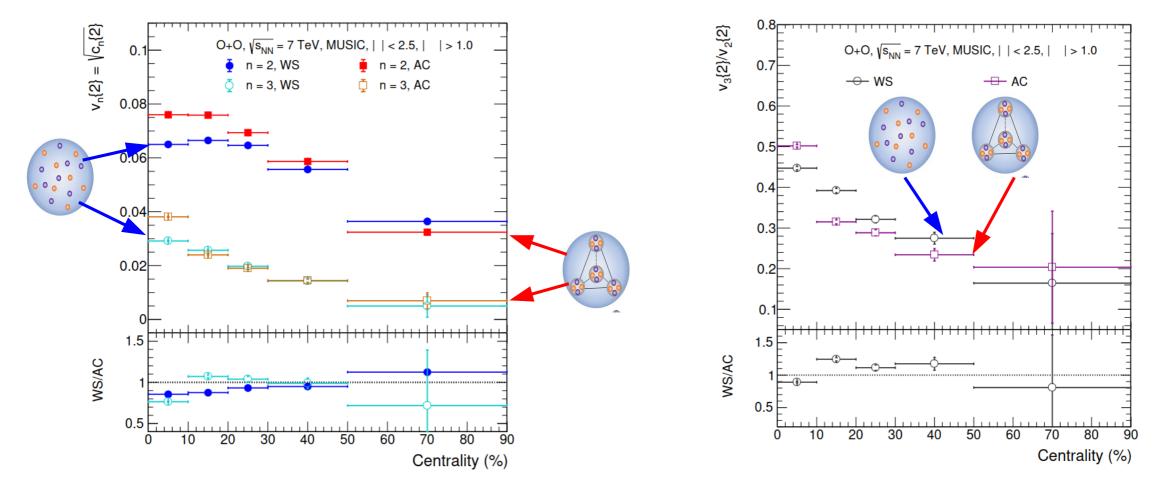
G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Flow in oxygen-oxygen (OO)

#### 2-cummulants based calculation of v<sub>2</sub> & v<sub>3</sub>

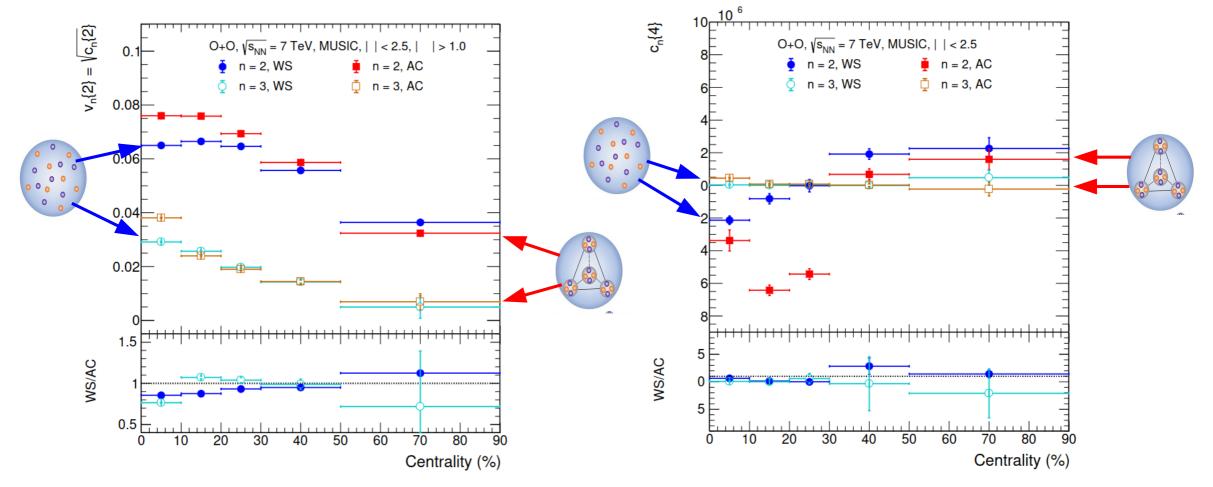


#### 2-cummulants based calculation of v<sub>2</sub> & v<sub>3</sub>



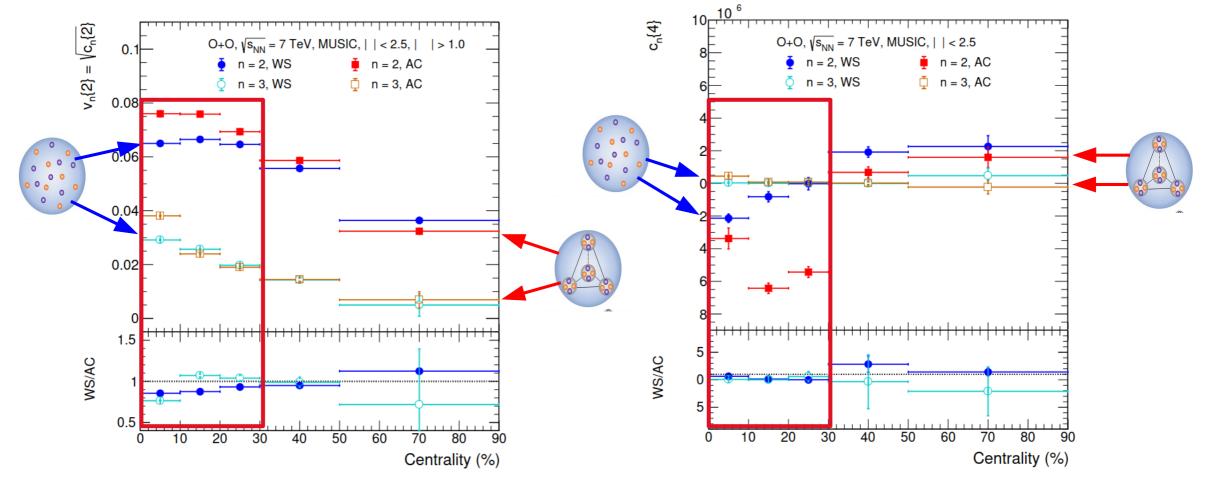
G.G. Barnafoldi: Triggering HEP, Tatry 2024

#### 2- & 4-cummulants based v<sub>n</sub> & c<sub>n</sub> calculations



G.G. Barnafoldi: Triggering HEP, Tatry 2024

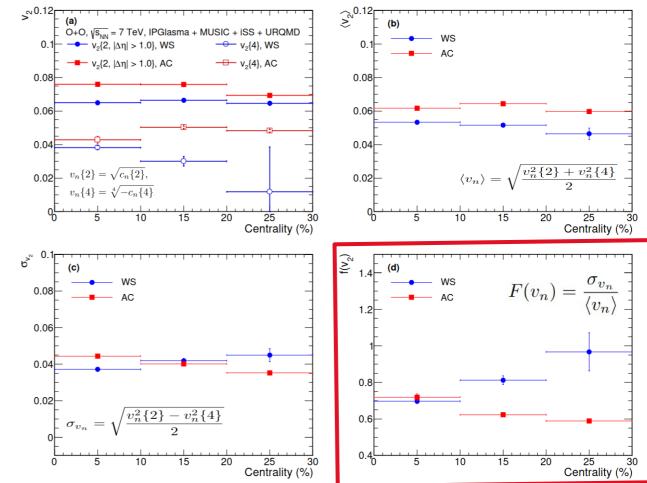
#### 2- & 4-cummulants based v<sub>n</sub> & c<sub>n</sub> calculations



G.G. Barnafoldi: Triggering HEP, Tatry 2024

#### 2- & 4-cummulants based calculations

- Flow and fluctuation measures changed significantly in the most central 0-30% regime
- Alpha-cluster has larger values, than Wood-Saxon profile
- Higher cummulants has higher effect at larger centrality
- Clearly visible on the relative measure:  $F(v_n) = \frac{\sigma_{v_n}}{\langle v_n \rangle}$



G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Conclusions

- In a IPGlasma+MUSIC+iSS+URQMD = "realistic model"
  - It is possible to calculate the flow for small system like OO
    - $\rightarrow$  event plane method fails
    - $\rightarrow$  2- & 4-cummulants can be calculated for v2
    - $\rightarrow$  v3 can not be calculated for 4-cummulant
    - → Need for a kinematical cut to reduce non-flow
- Nuclear structure has consequences on the flow
  - Nuclear structure matters in the calculations
    - $\rightarrow$  Alpha Cluster method is stronger than Woods-Saxon
    - $\rightarrow$  Relevant difference is in centra O+O collisions
    - → Comparable with the size of the alpha cluster

### Thank You!

# Can we prove the model' validity in heavy-ion collisions?

### Calculating the flow

#### **Event plane and average method**

$$v_n = \langle \cos[n(\phi - \psi_n)] \rangle$$

#### **Multiparticle Q-cummulant method**

- Flow vector  $Q_n = \sum_{j=1}^M e^{in\phi_j}$
- The 2- and 4-particle cummulants are:

$$\begin{aligned} \langle 2 \rangle &= \frac{|Q_n|^2 - M}{M(M-1)}, \\ \langle 4 \rangle &= \frac{|Q_n|^4 + |Q_{2n}|^2 - 2 \cdot \operatorname{Re}[Q_{2n}Q_n^*Q_n^*]}{M(M-1)(M-2)(M-3)} & \qquad c_n\{2\} = \langle \langle 2 \rangle \rangle, \\ c_n\{4\} &= \langle \langle 4 \rangle \rangle - 2 \cdot \langle \langle 2 \rangle \rangle^2 \\ &= \sqrt{c_n\{2\}}, \\ c_n\{4\} &= \langle \langle 4 \rangle \rangle - 2 \cdot \langle \langle 2 \rangle \rangle^2 \\ &= \sqrt{c_n\{4\}} = \sqrt[4]{-c_n\{4\}}, \end{aligned}$$

G.G. Barnafoldi: Triggering HEP, Tatry 2024

### Calculating the flow

#### Suppressing the non-flow contribution:

Kinematical cut: 2 sub-events, A&B are intoduced, with a rapidity gap:

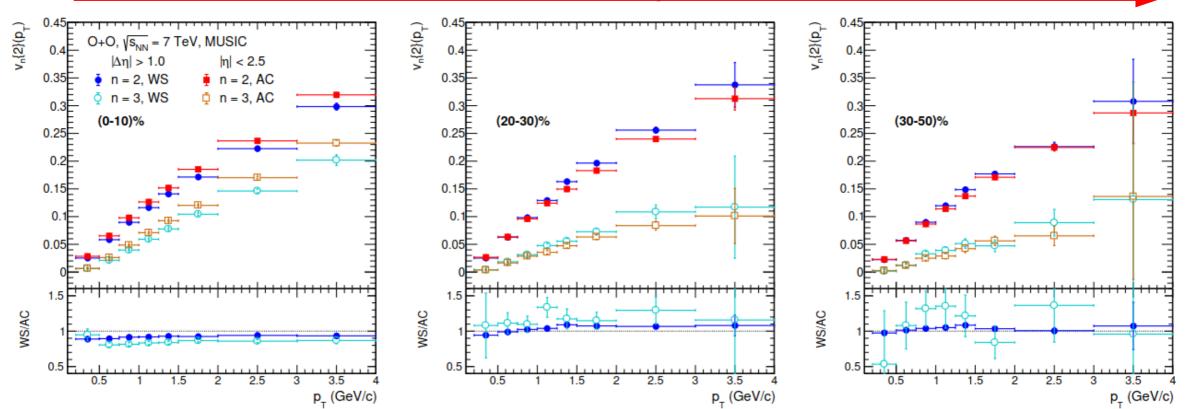
$$\langle 2 \rangle_{\Delta \eta} = \frac{Q_n^A \cdot Q_n^{B*}}{M_A \cdot M_B} \qquad \qquad \blacktriangleright \qquad v_n \{2, |\Delta \eta|\}(p_{\rm T}) = \frac{d_n \{2, |\Delta \eta|\}}{\sqrt{c_n \{2, |\Delta \eta|\}}}$$

Differential flow cummulants:

- Mean and the fluctuations of the flow & ratio:



#### 2-cummulants based $v_n(p_T)$ calculations



#### **Centrality**

G.G. Barnafoldi: Triggering HEP, Tatry 2024