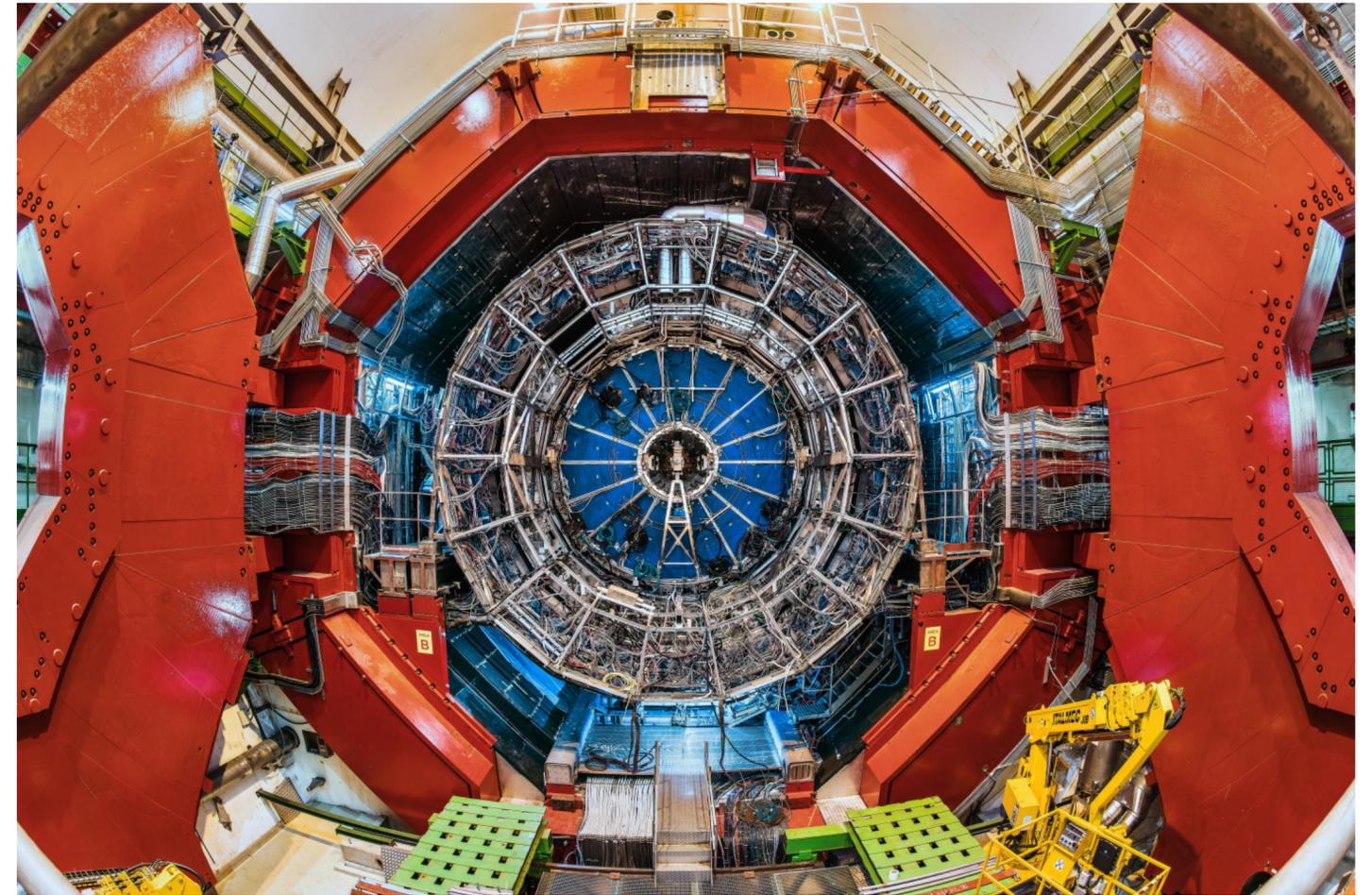


Understanding the Quark-Gluon Plasma with ALICE at the LHC

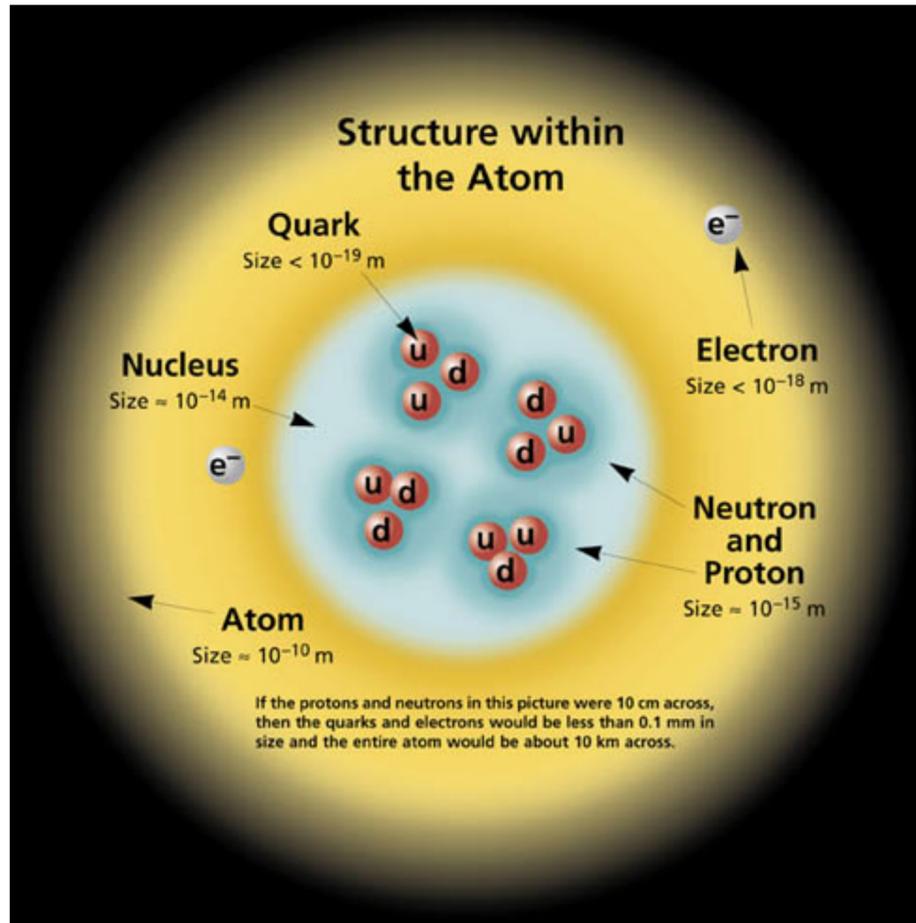
Marco van Leeuwen, Nikhef and CERN



Wigner 121 Symposium, 18-20 September 2023, Budapest

Particles and interactions

Atomic scale: electromagnetic interactions



Electrons, protons carry 1 unit of electric charge

Standard Model: quantum field theory of fundamental particles

	mass → 2.4 MeV/c ² charge → 2/3 spin → 1/2 u up	mass → 1.27 GeV/c ² charge → 2/3 spin → 1/2 c charm	mass → 171.2 GeV/c ² charge → 2/3 spin → 1/2 t top	mass → 0 charge → 0 spin → 1 γ photon	mass → ≈126 GeV/c ² charge → 0 spin → 0 H Higgs boson
QUARKS	mass → 4.8 MeV/c ² charge → -1/3 spin → 1/2 d down	mass → 104 MeV/c ² charge → -1/3 spin → 1/2 s strange	mass → 4.2 GeV/c ² charge → -1/3 spin → 1/2 b bottom	mass → 0 charge → 0 spin → 1 g gluon	
	mass → 0.511 MeV/c ² charge → -1 spin → 1/2 e electron	mass → 105.7 MeV/c ² charge → -1 spin → 1/2 μ muon	mass → 1.777 GeV/c ² charge → -1 spin → 1/2 τ tau	mass → 91.2 GeV/c ² charge → 0 spin → 1 Z Z boson	
LEPTONS	mass → <math>< 2.2</math> eV/c ² charge → 0 spin → 1/2 ν_e electron neutrino	mass → <math>< 0.17</math> MeV/c ² charge → 0 spin → 1/2 ν_μ muon neutrino	mass → <math>< 15.5</math> MeV/c ² charge → 0 spin → 1/2 ν_τ tau neutrino	mass → 80.4 GeV/c ² charge → ±1 spin → 1 W W boson	GAUGE BOSONS

Leptons (e, μ, τ) and photon are the fundamental particles of the electromagnetic interactions

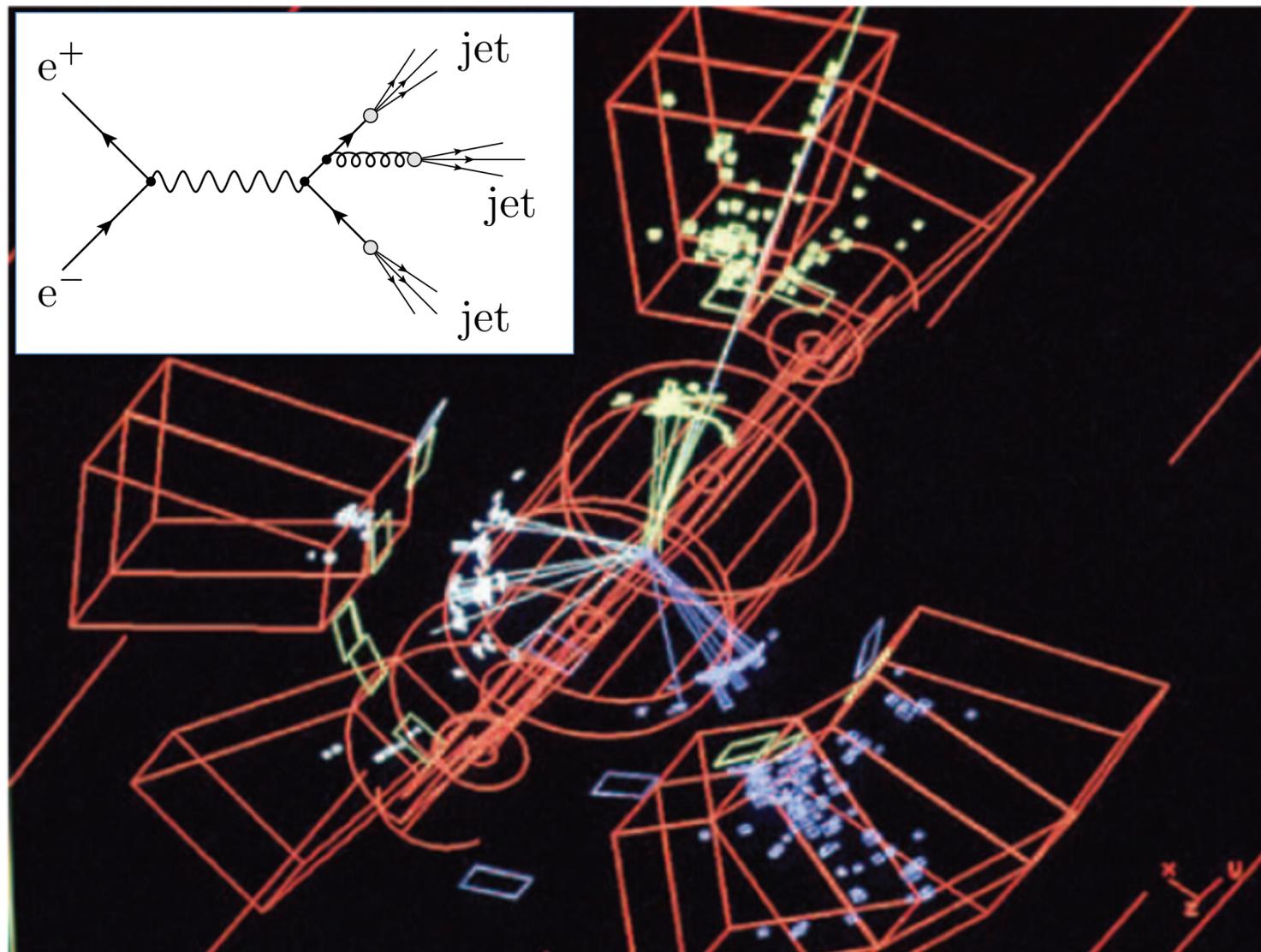
Quarks and gluons are the fundamental particles of the strong interaction

Dominant interaction on the subatomic scale (> MeV, < 1 fm)

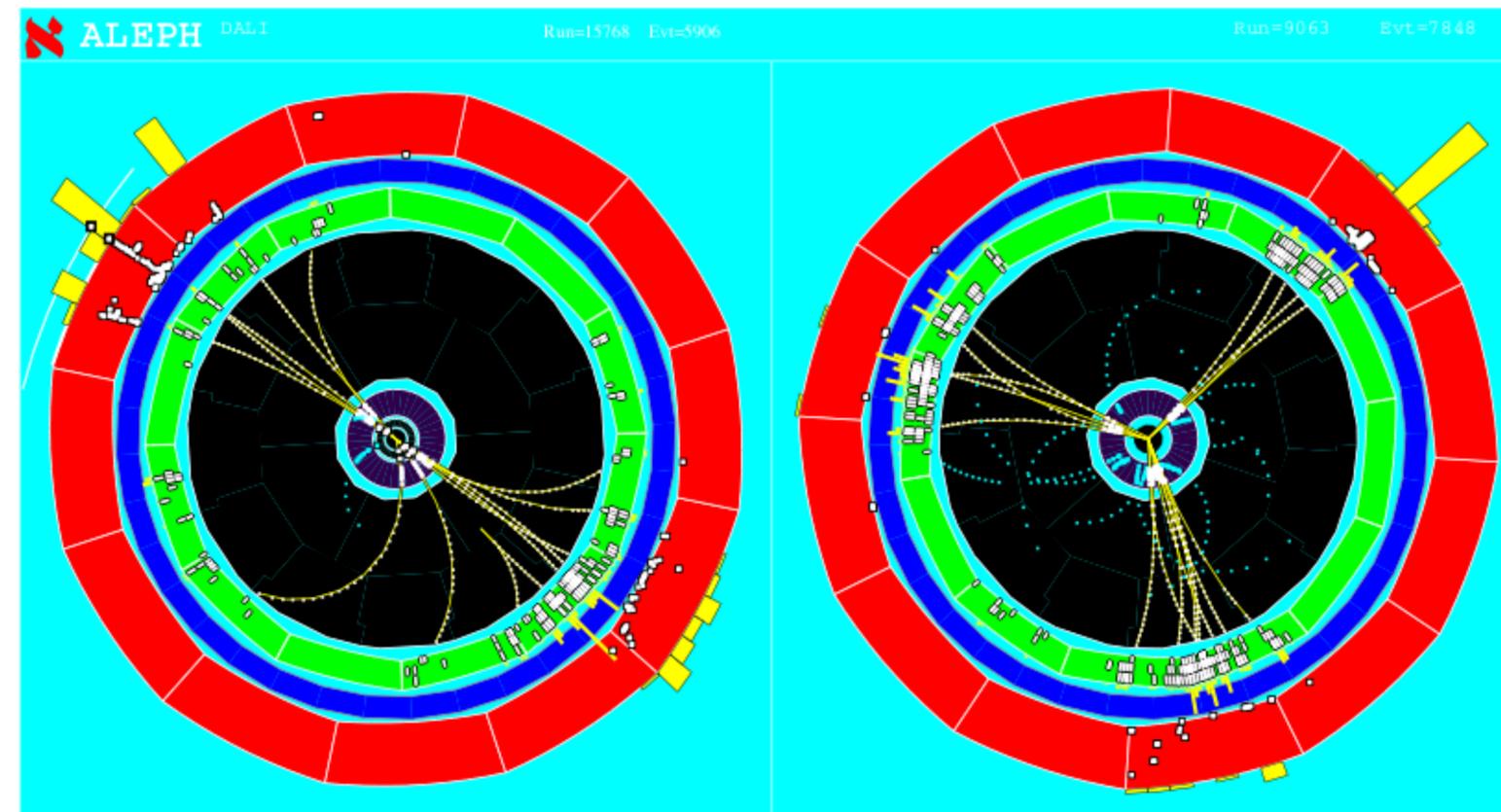
Experimental signatures of quarks and gluons: three-jet events

TASSO experiment @ PETRA @ DESY

ALEPH @ LEP (CERN)



collision energy: $\sqrt{s} = 13\text{-}31 \text{ GeV}$

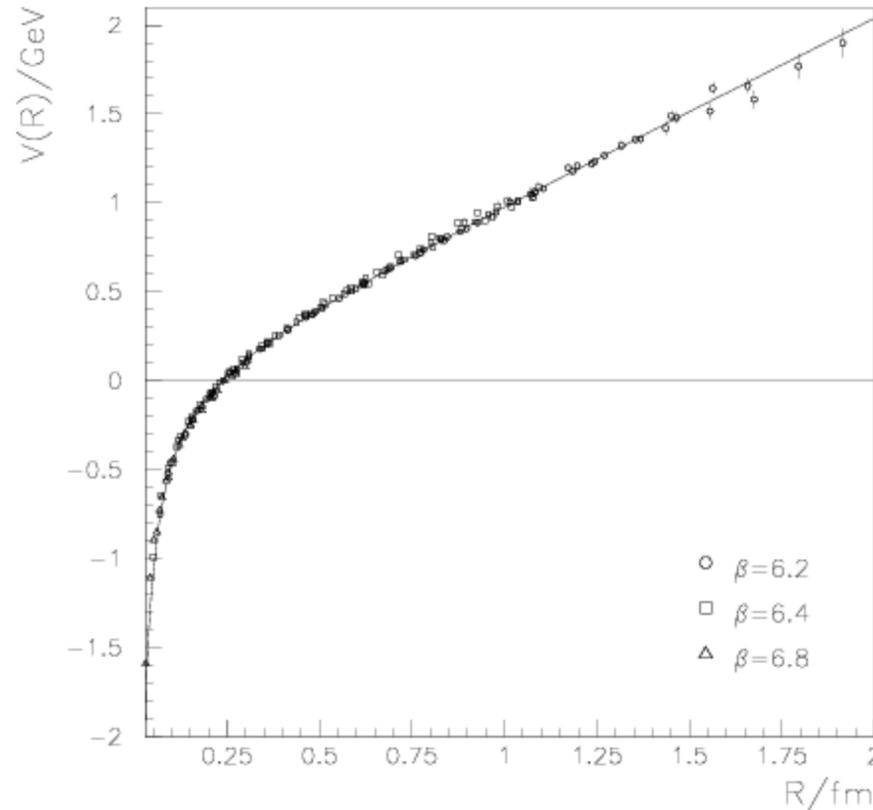


collision energy: $\sqrt{s} = 90\text{-}209 \text{ GeV}$

NB: quarks and gluons are not detected individually — parton shower and hadronisation

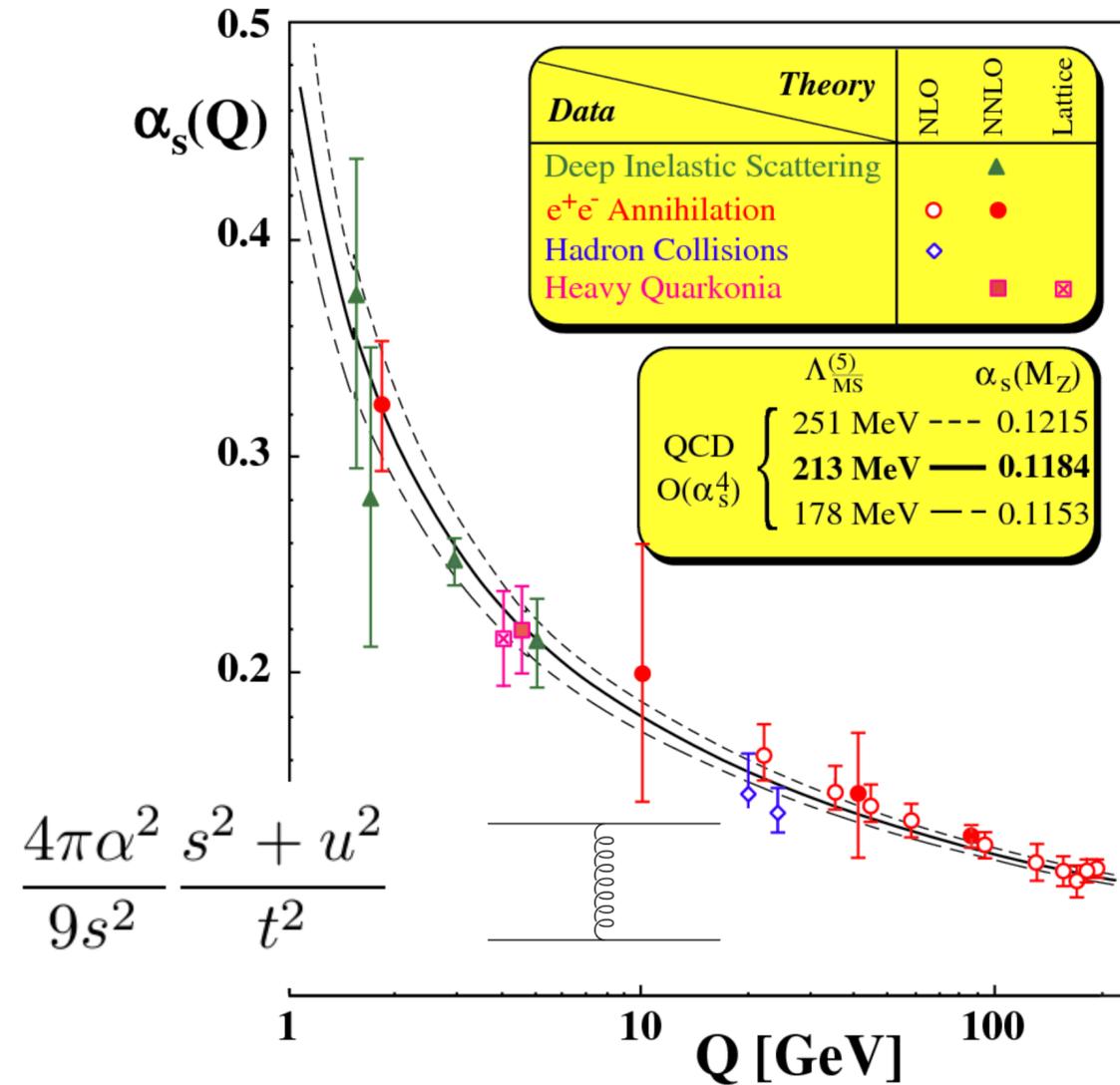
The running coupling of QCD

QCD potential
(lattice QCD)

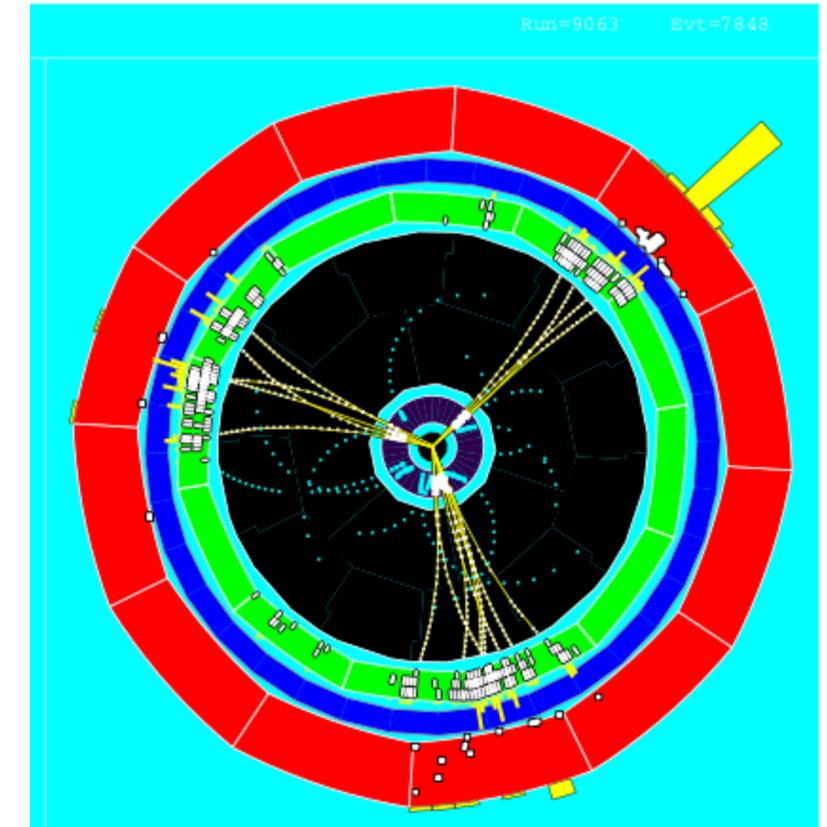


Bali, hep-lat/9311009

QCD coupling constant



Low-energy, large distances:
quarks are bound into hadrons



High energy, short distances:
quarks and gluons interact as
quasi-free particles

At large distance, small energy: perturbative calculations do not converge
Static QCD potential does not capture full dynamics

Understanding the interactions

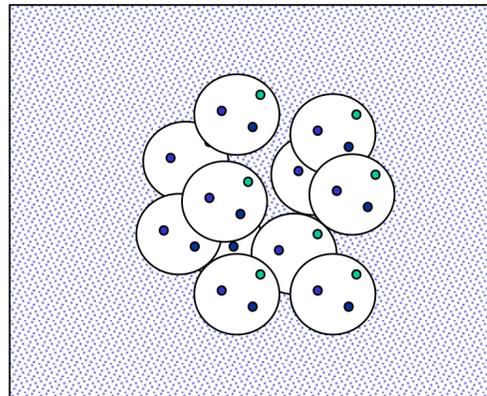
Three regimes giving rise to subfields of physics:

	Electromagnetism	Strong interaction
Free particles	Two-body scattering	
Bound states	Atomic physics	Hadronic, nuclear physics
Many-body physics	Condensed matter: thermal, electrical properties, superconductivity, etc	Heavy-ion physics: quark-gluon plasma

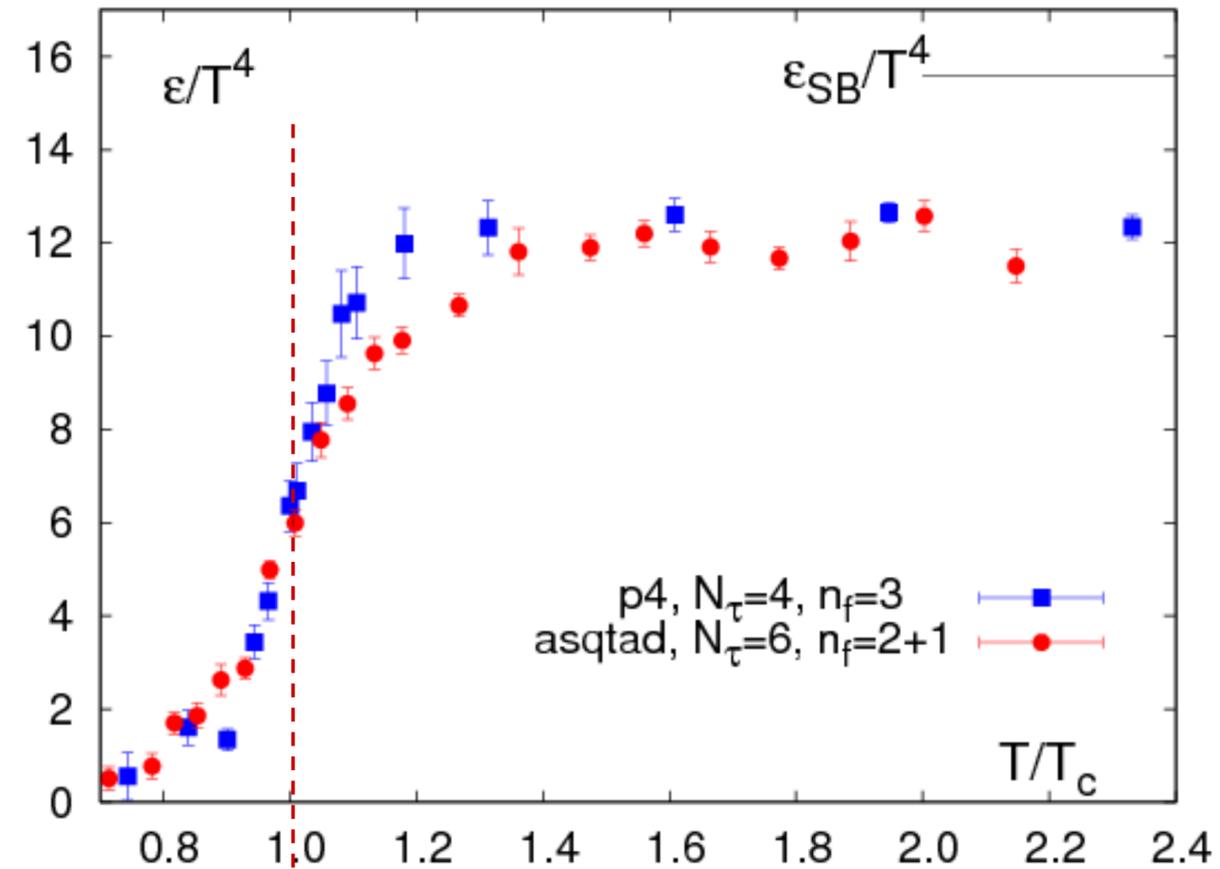
Heavy-ion collisions are used to study ‘condensed matter physics’ of QCD
Unique form of ‘quantum condensed matter’

Condensed matter of QCD: the quark-gluon plasma

Lattice QCD calculations: energy density vs temperature



Low temperature:
quarks and gluons confined in hadrons

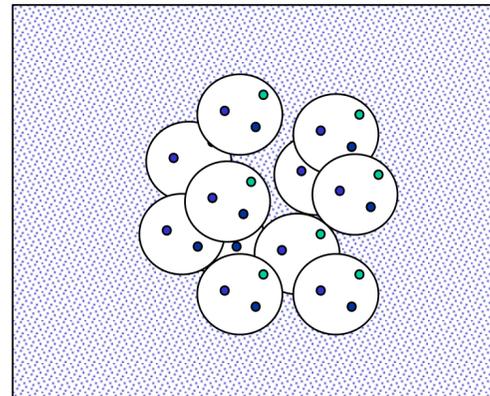


$T_c \sim 155$ MeV
 $\epsilon_c \sim 1$ GeV/fm³

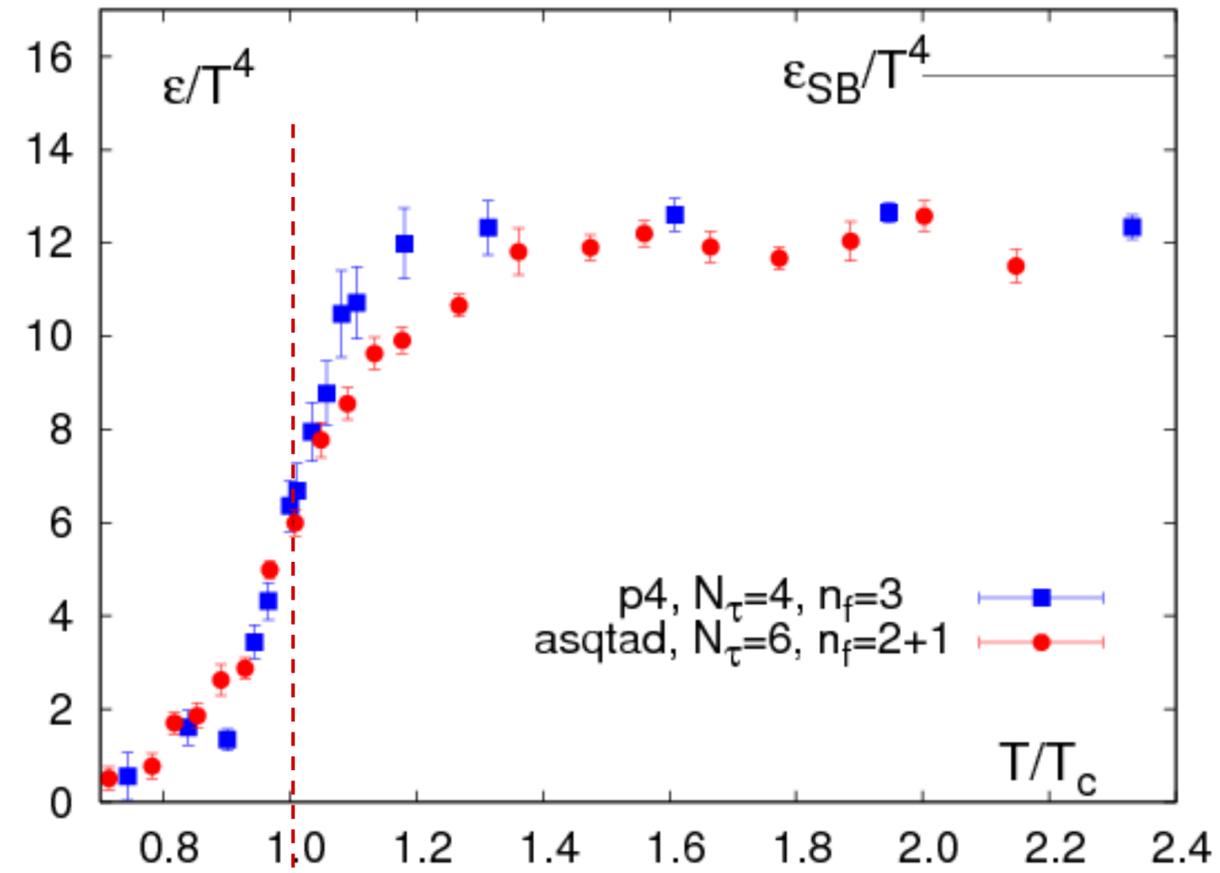
Bernard et al. hep-lat/0610017

Condensed matter of QCD: the quark-gluon plasma

Lattice QCD calculations: energy density vs temperature

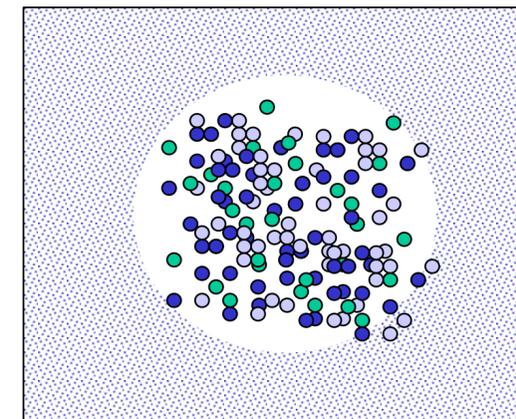


Low temperature:
quarks and gluons confined in hadrons



$T_c \sim 155 \text{ MeV}$
 $\epsilon_c \sim 1 \text{ GeV/fm}^3$

Bernard et al. hep-lat/0610017

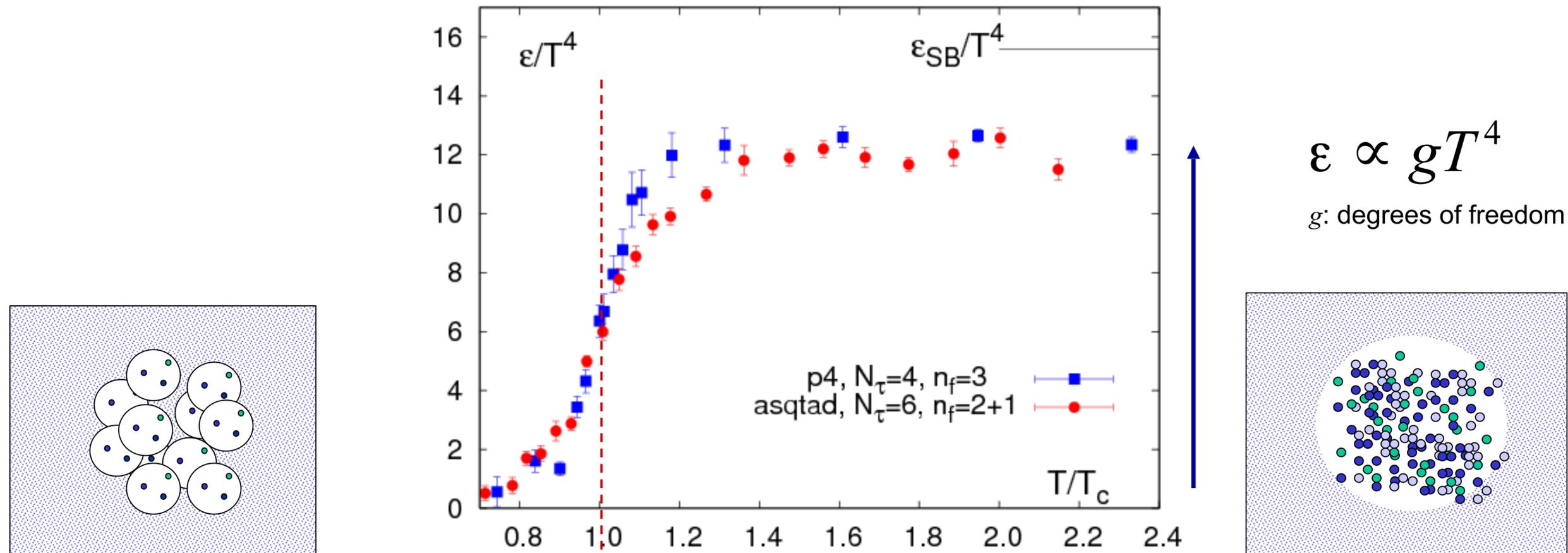


High temperature:
deconfined quark-gluon plasma

Phase transition at critical temperature $T_c \approx 155 \text{ MeV} \approx 10^{12} \text{ K}$

Condensed matter of QCD: the quark-gluon plasma

Lattice QCD calculations: energy density vs temperature



Low temperature:
quarks and gluons confined in hadrons

$T_c \sim 155$ MeV
 $\epsilon_c \sim 1$ GeV/fm³

Bernard et al. hep-lat/0610017

High temperature:
deconfined quark-gluon plasma

Phase transition at critical temperature $T_c \approx 155$ MeV $\approx 10^{12}$ K

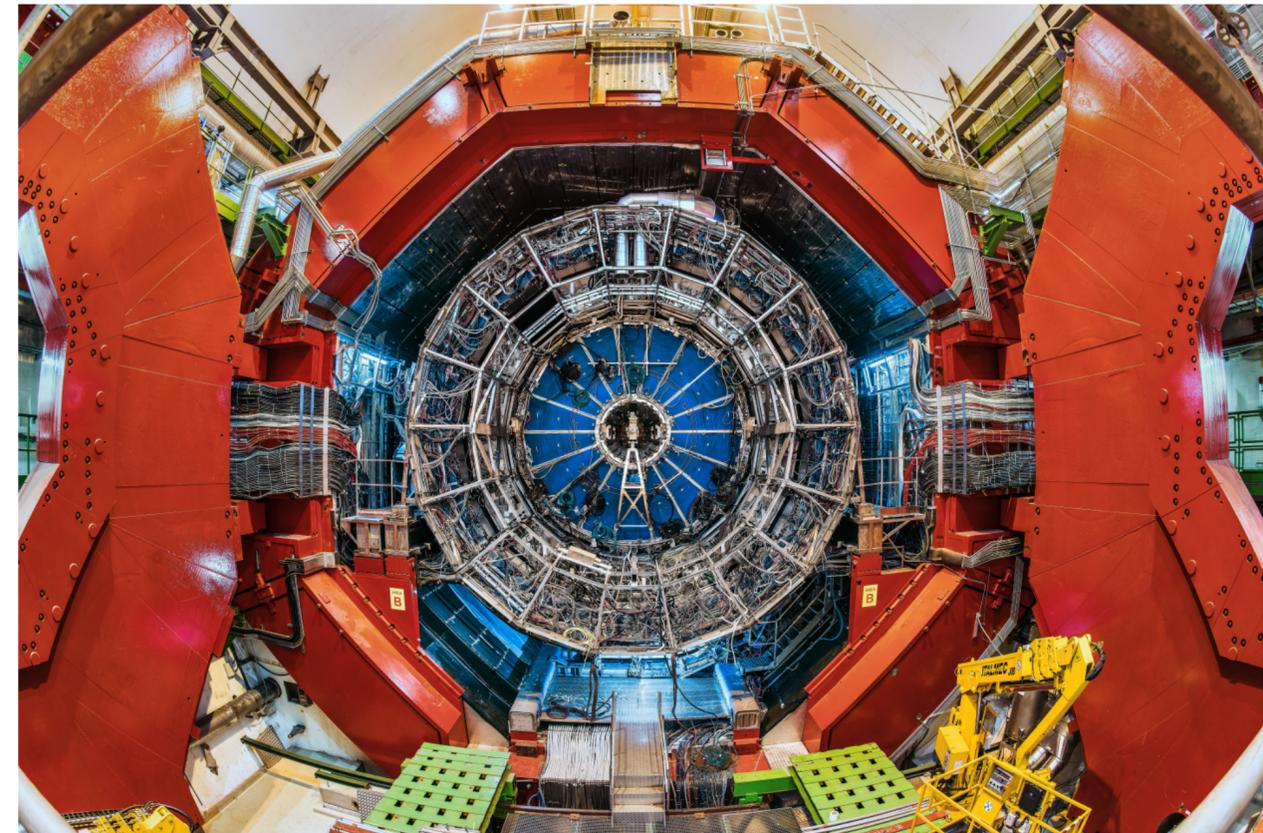
Increase of number of degrees of freedom: hadrons (3 pions) \rightarrow quarks+gluons (37)

The Large Hadron Collider and ALICE

The Large Hadron Collider

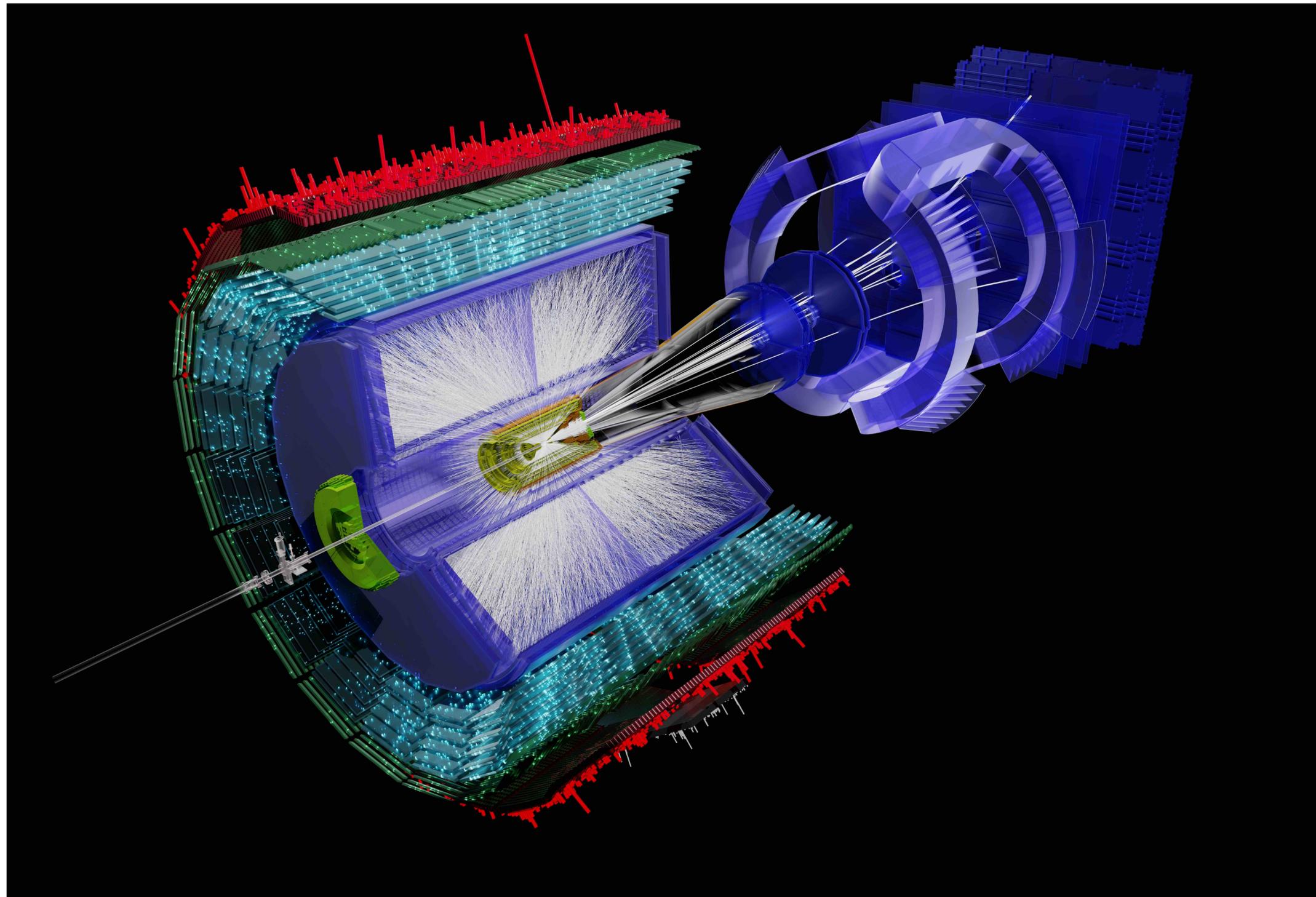


A Large Ion Collider Experiment



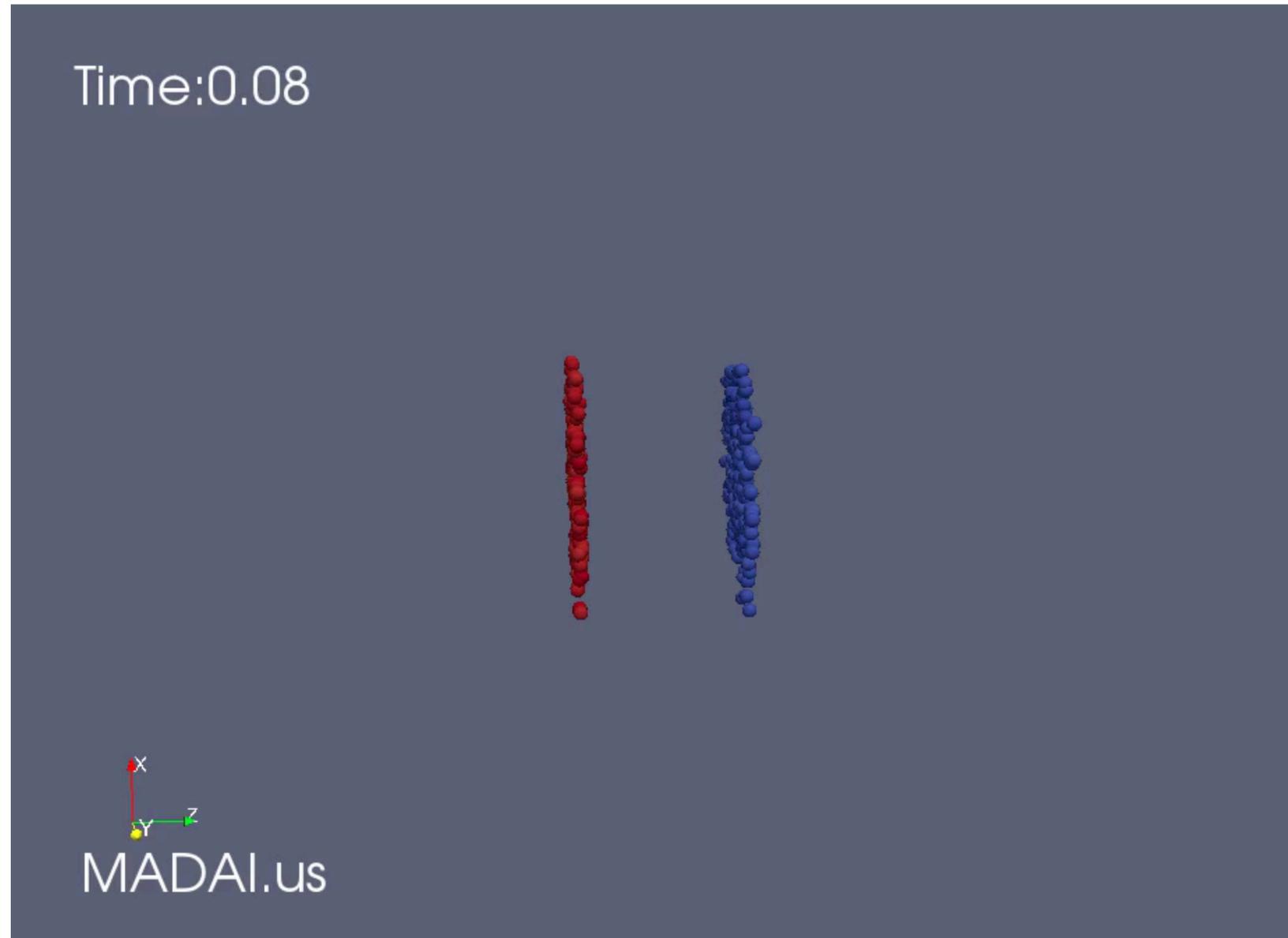
LHC: most powerful particle accelerator-collider in the world
pp collision: 13.6 TeV
Pb-Pb collisions: 5.36 TeV per nucleon pair

ALICE: one of the four large LHC experiments
Focus on strong interaction, heavy-ion collisions



**Thousands of particles are produced in each lead ion collision
— study momentum distributions, correlations, types of particles**

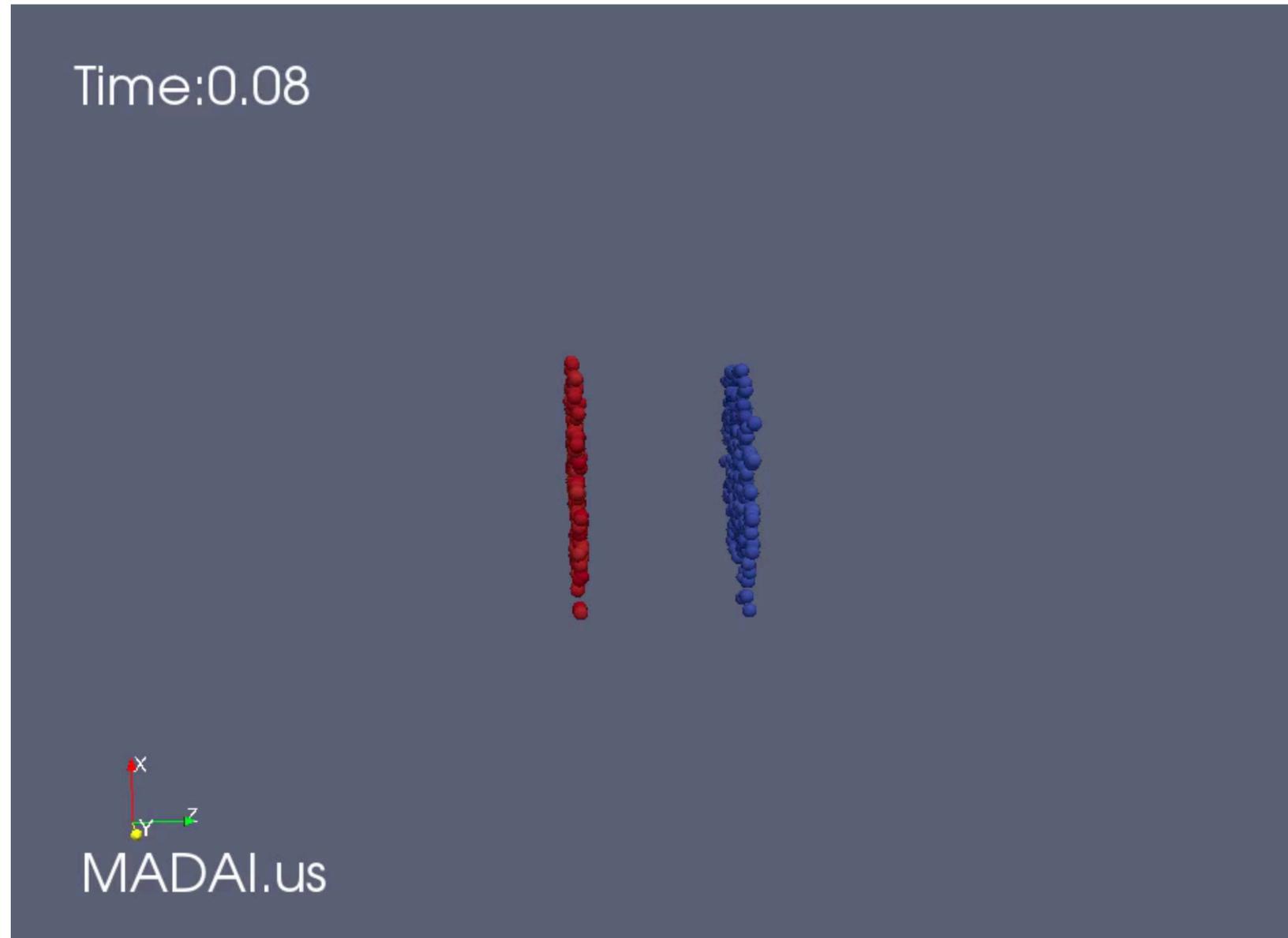
Heavy ion collisions: Little Bangs



Stages of the collision: initial stages — QGP/fluid stage — hadron formation (freeze out)

‘Little Bang’: recreate primordial matter in the laboratory

Heavy ion collisions: Little Bangs

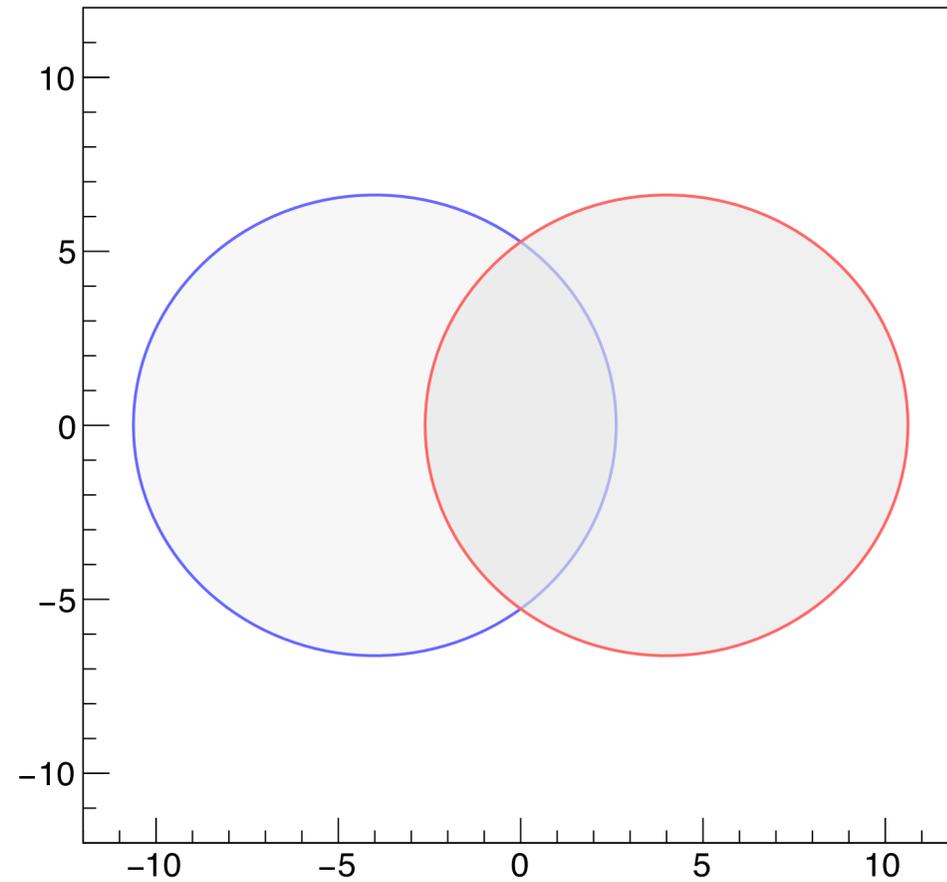


Stages of the collision: initial stages — QGP/fluid stage — hadron formation (freeze out)

‘Little Bang’: recreate primordial matter in the laboratory

Azimuthal anisotropy: initial and final states

Simulated event: location of nucleons

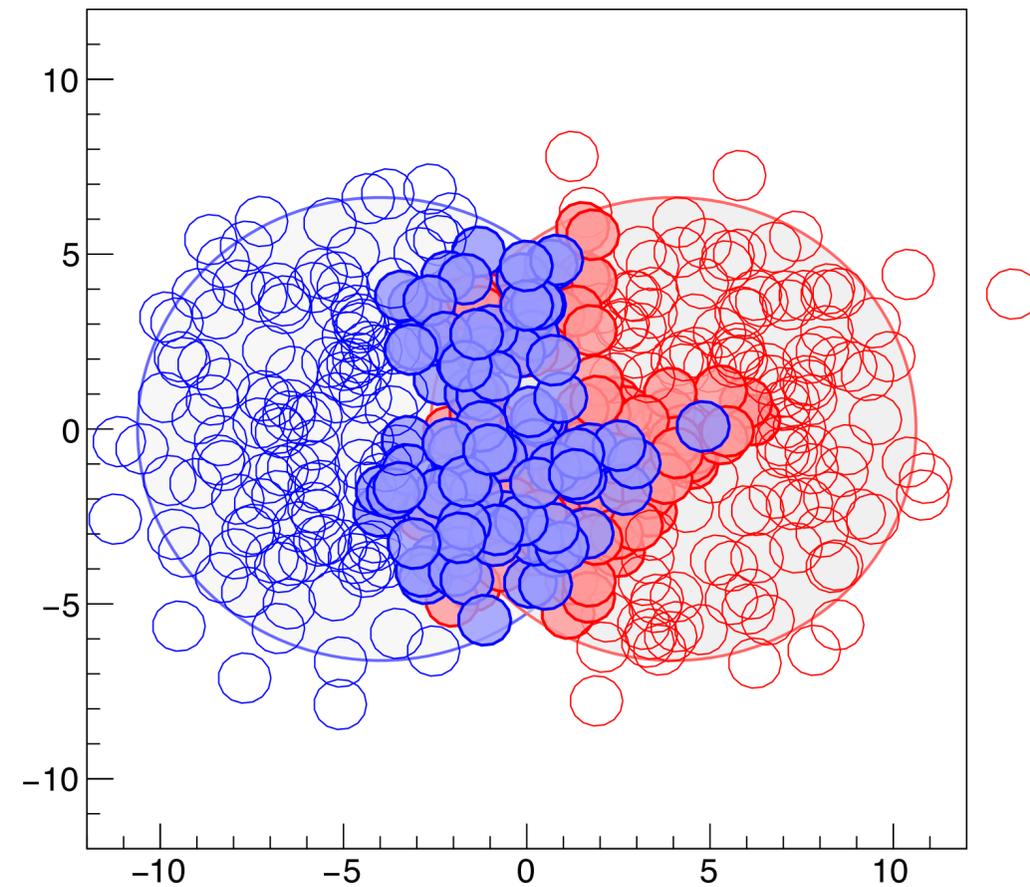


Characterise shape by harmonics:

$$\varepsilon_n = \frac{\sum r^2 (\cos^2 n\varphi + \sin^2 n\varphi)}{\sum r^2}$$

Azimuthal anisotropy: initial and final states

Simulated event: location of nucleons

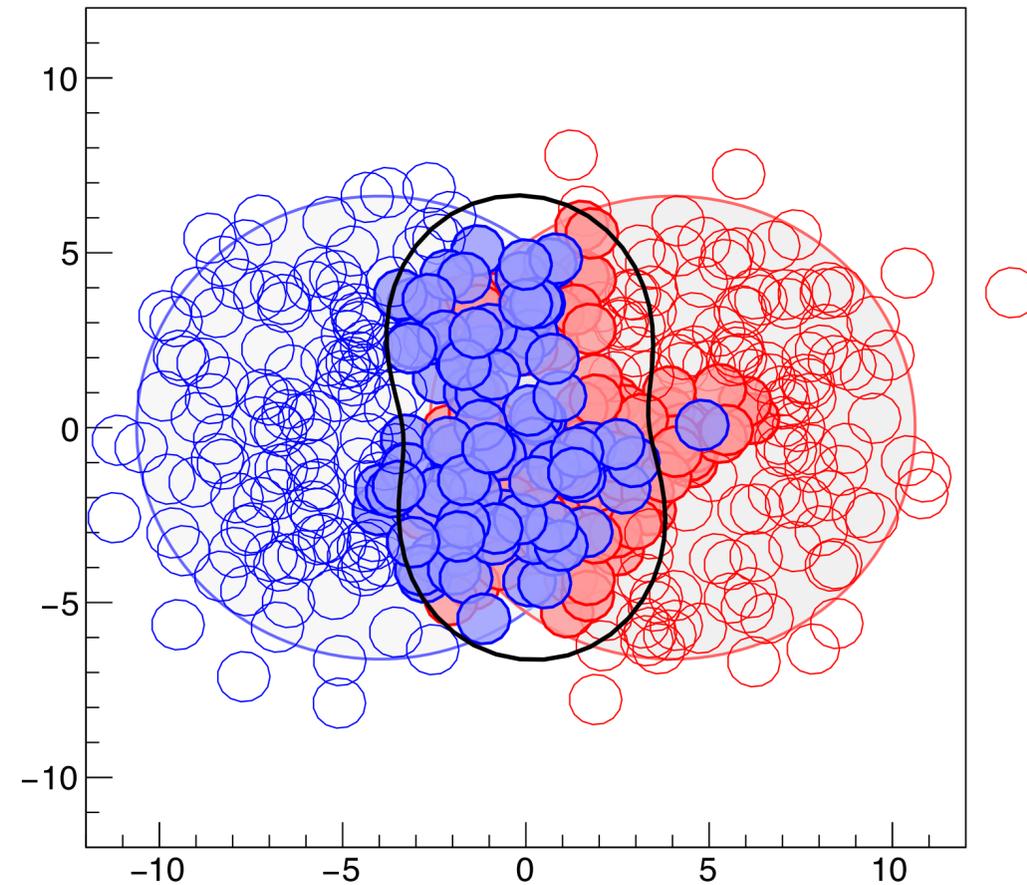


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Azimuthal anisotropy: initial and final states

Simulated event: location of nucleons

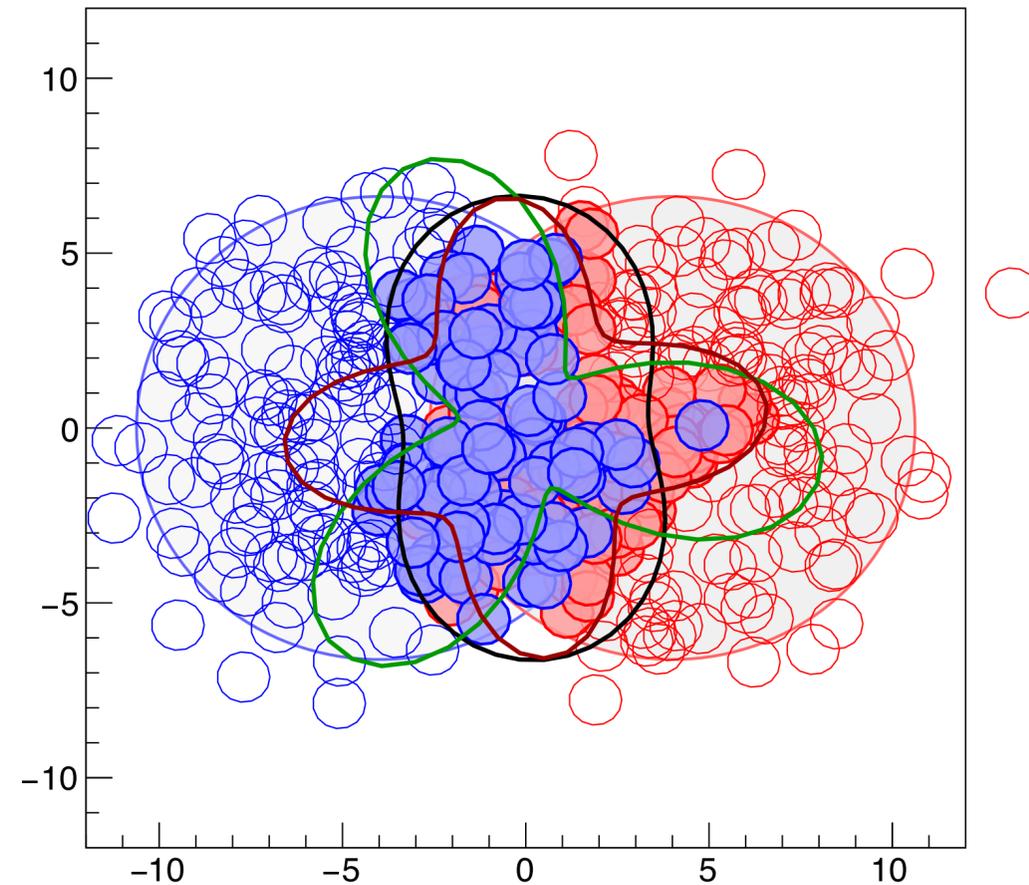


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Azimuthal anisotropy: initial and final states

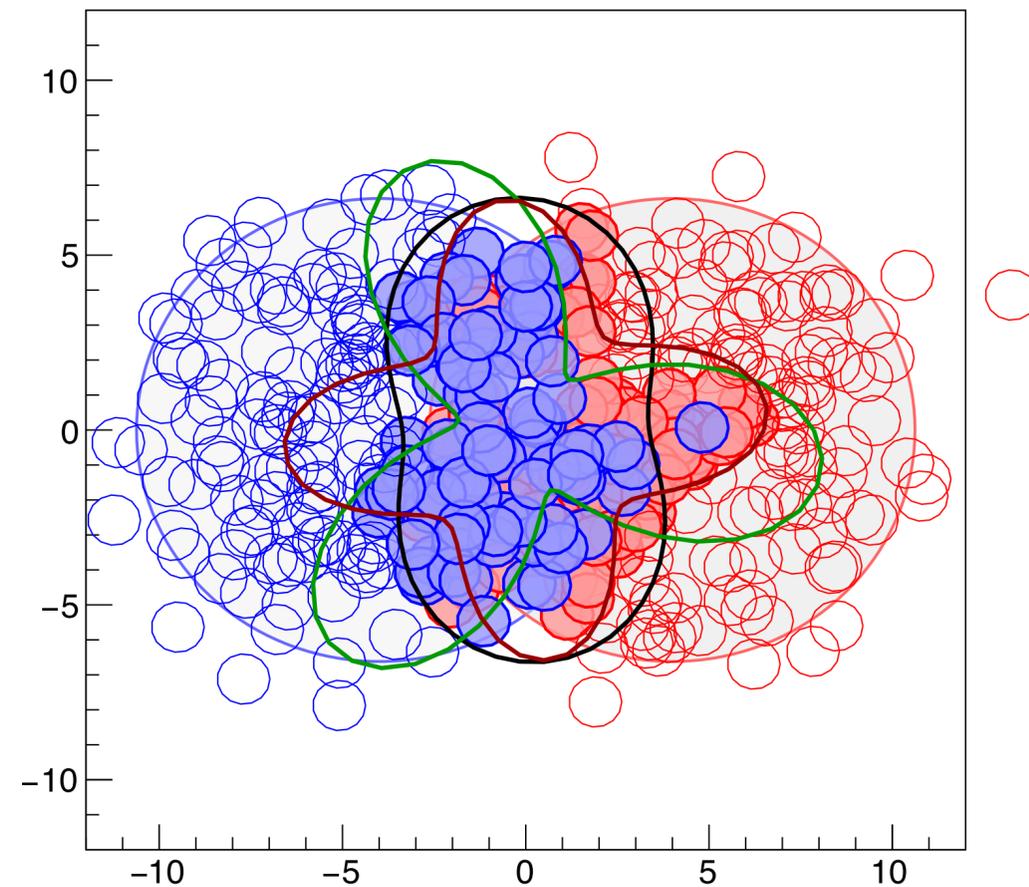
Simulated event: location of nucleons



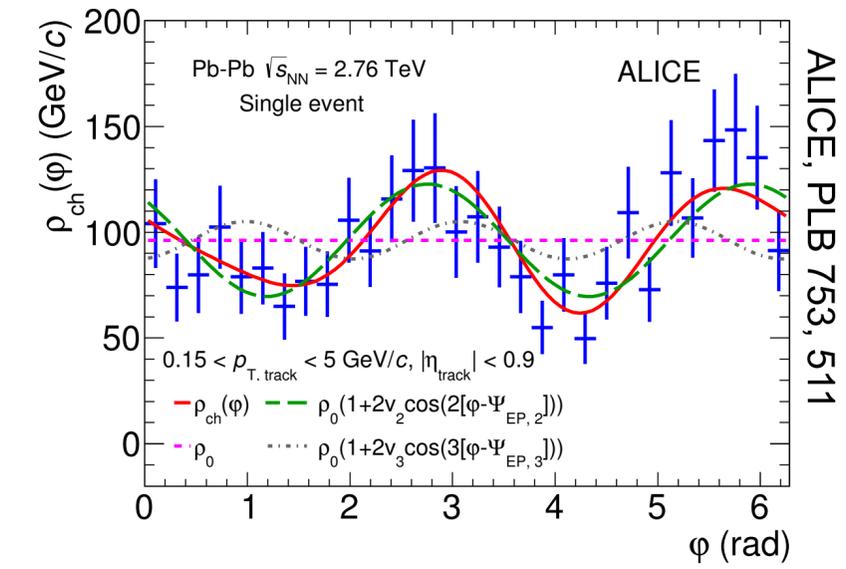
Initial state spatial anisotropies ε_n are transferred into
final state momentum anisotropies v_n
by pressure gradients, flow of the Quark Gluon Plasma

Azimuthal anisotropy: initial and final states

Simulated event: location of nucleons



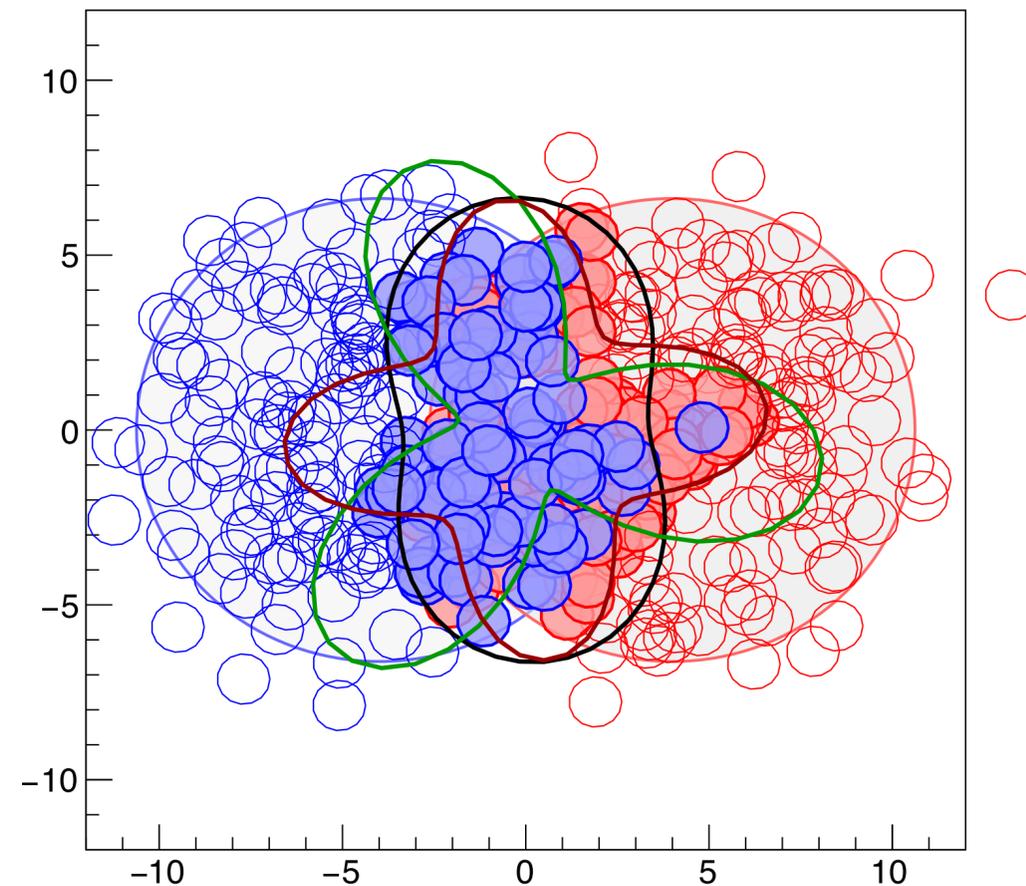
Azimuthal distribution single event



Initial state spatial anisotropies ε_n are transferred into
 final state momentum anisotropies v_n
 by pressure gradients, flow of the Quark Gluon Plasma

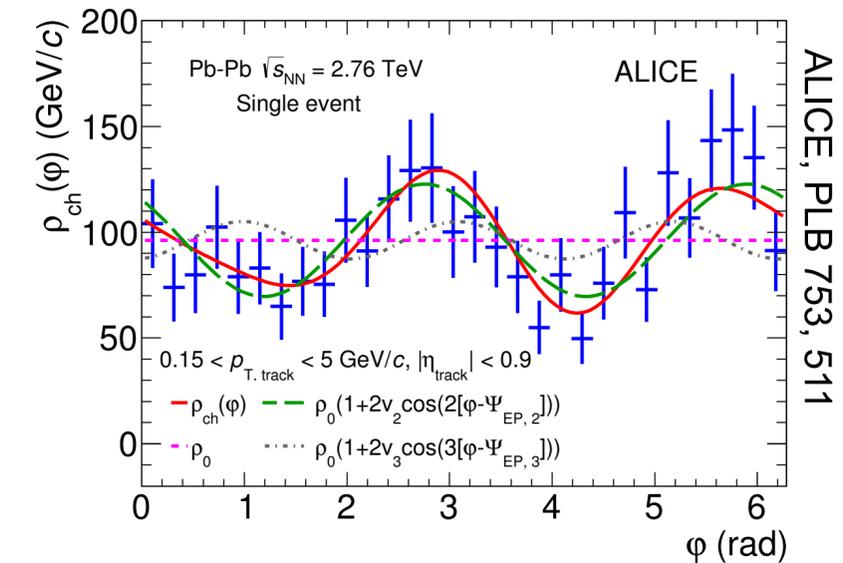
Azimuthal anisotropy: initial and final states

Simulated event: location of nucleons

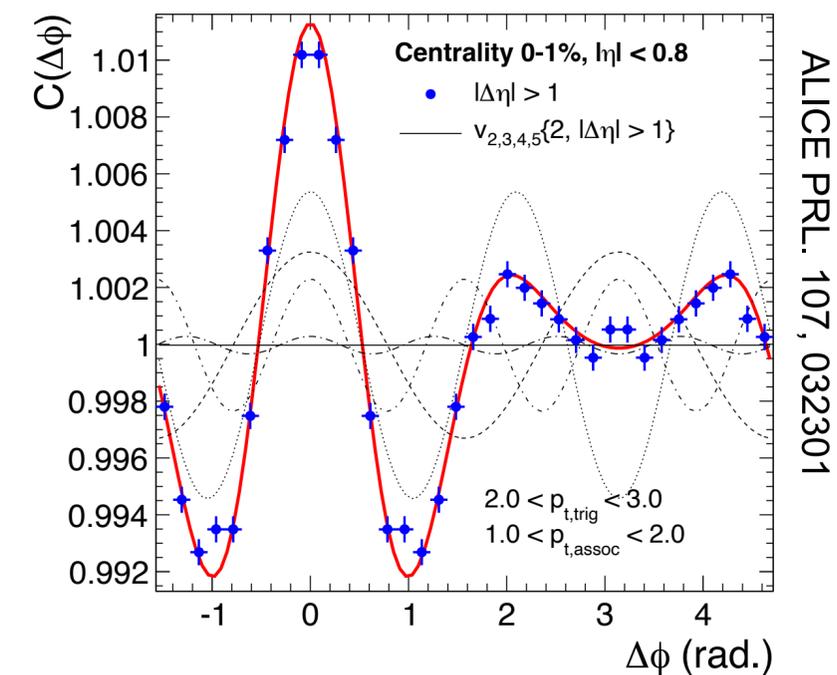


Initial state spatial anisotropies ϵ_n are transferred into final state momentum anisotropies v_n by pressure gradients, flow of the Quark Gluon Plasma

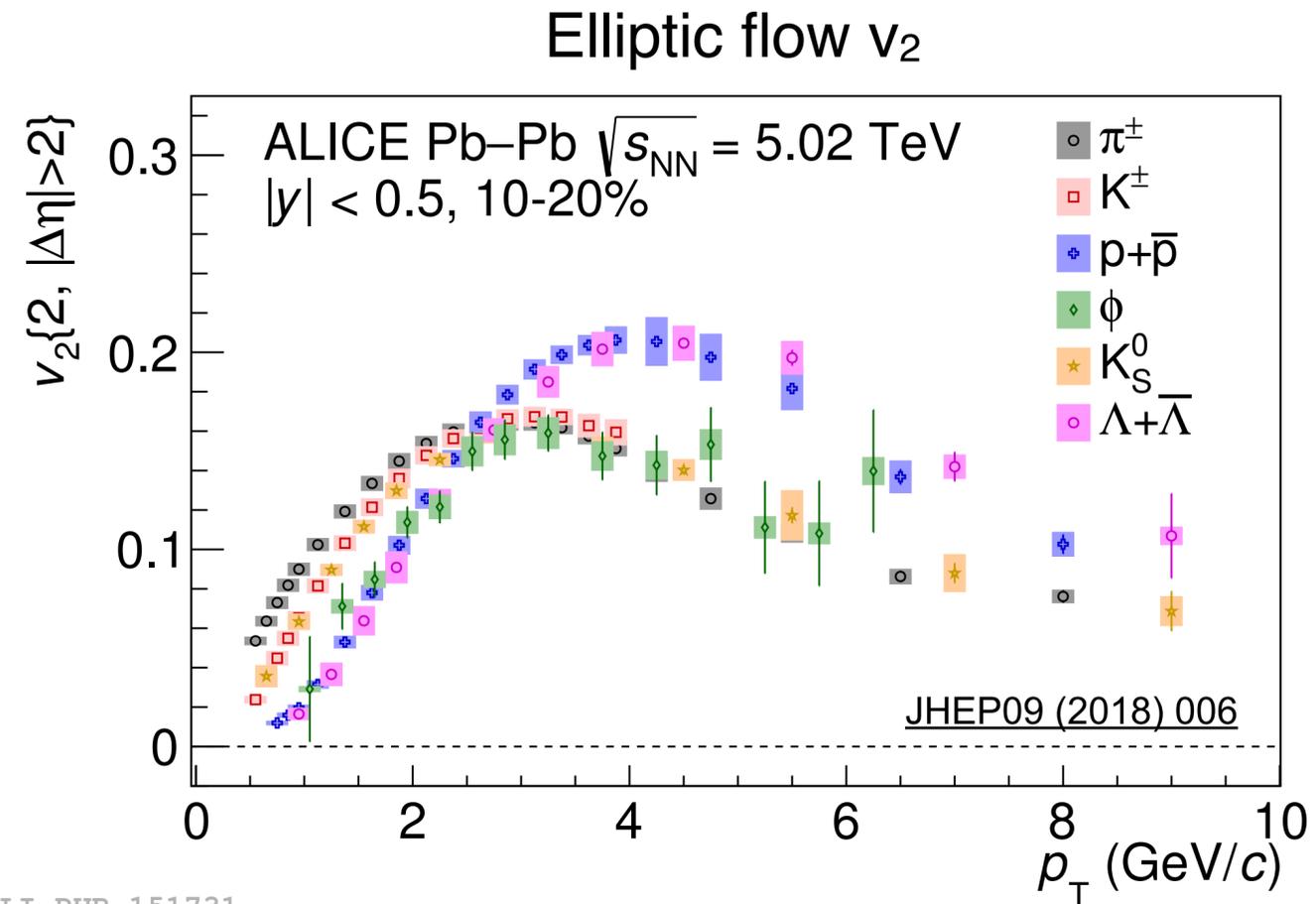
Azimuthal distribution single event



Sum over many events

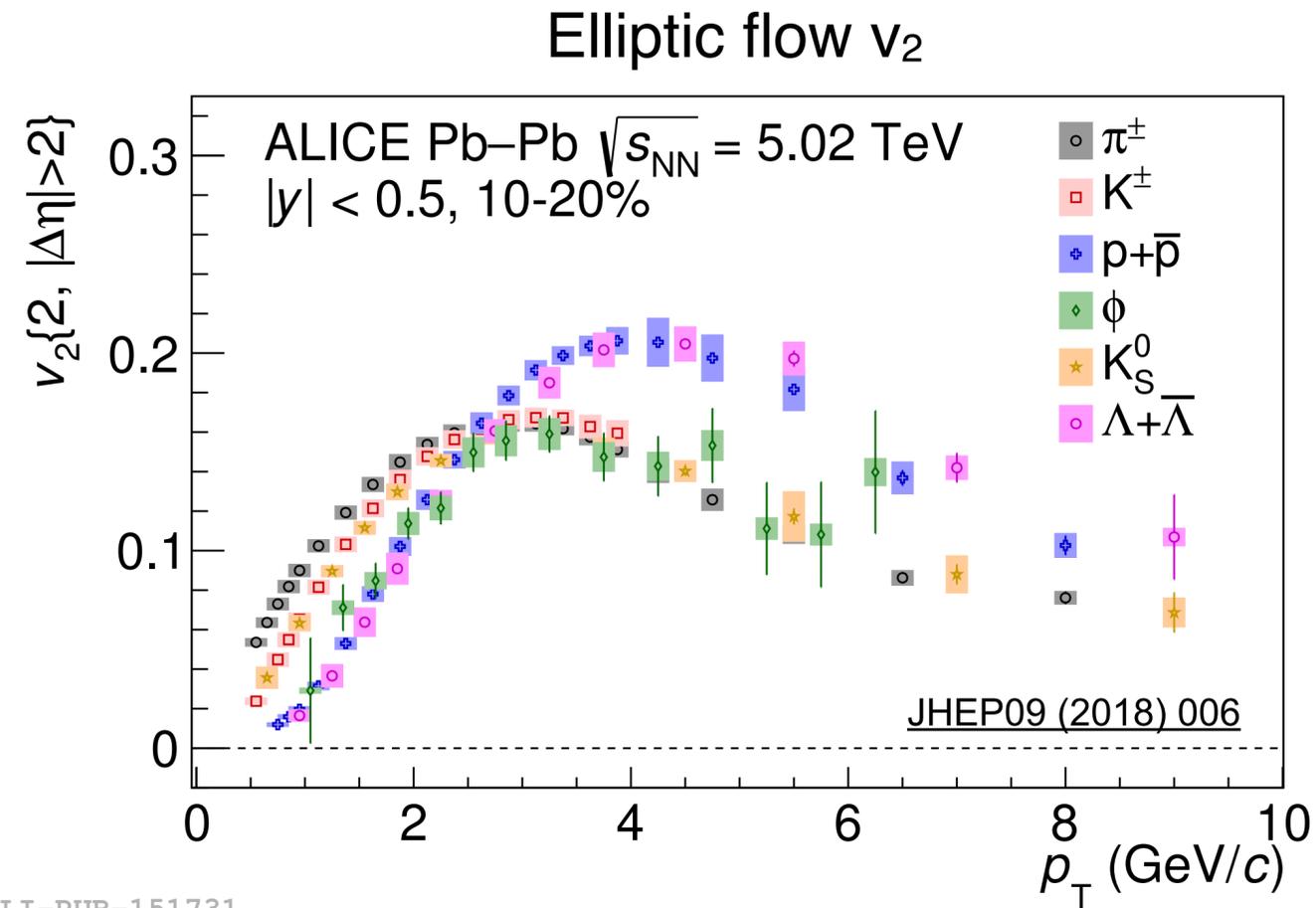


Anisotropic flow: initial state and QGP expansion

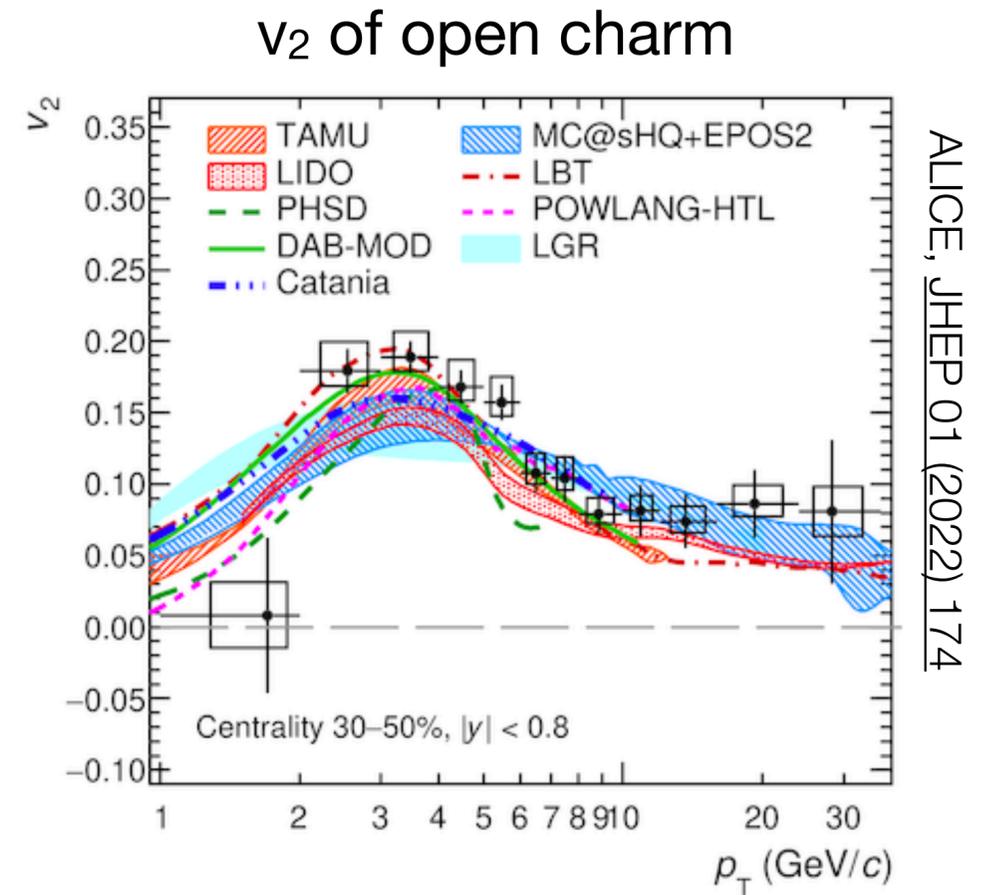


Mass-dependence of v_2 measures flow velocity

Anisotropic flow: initial state and QGP expansion



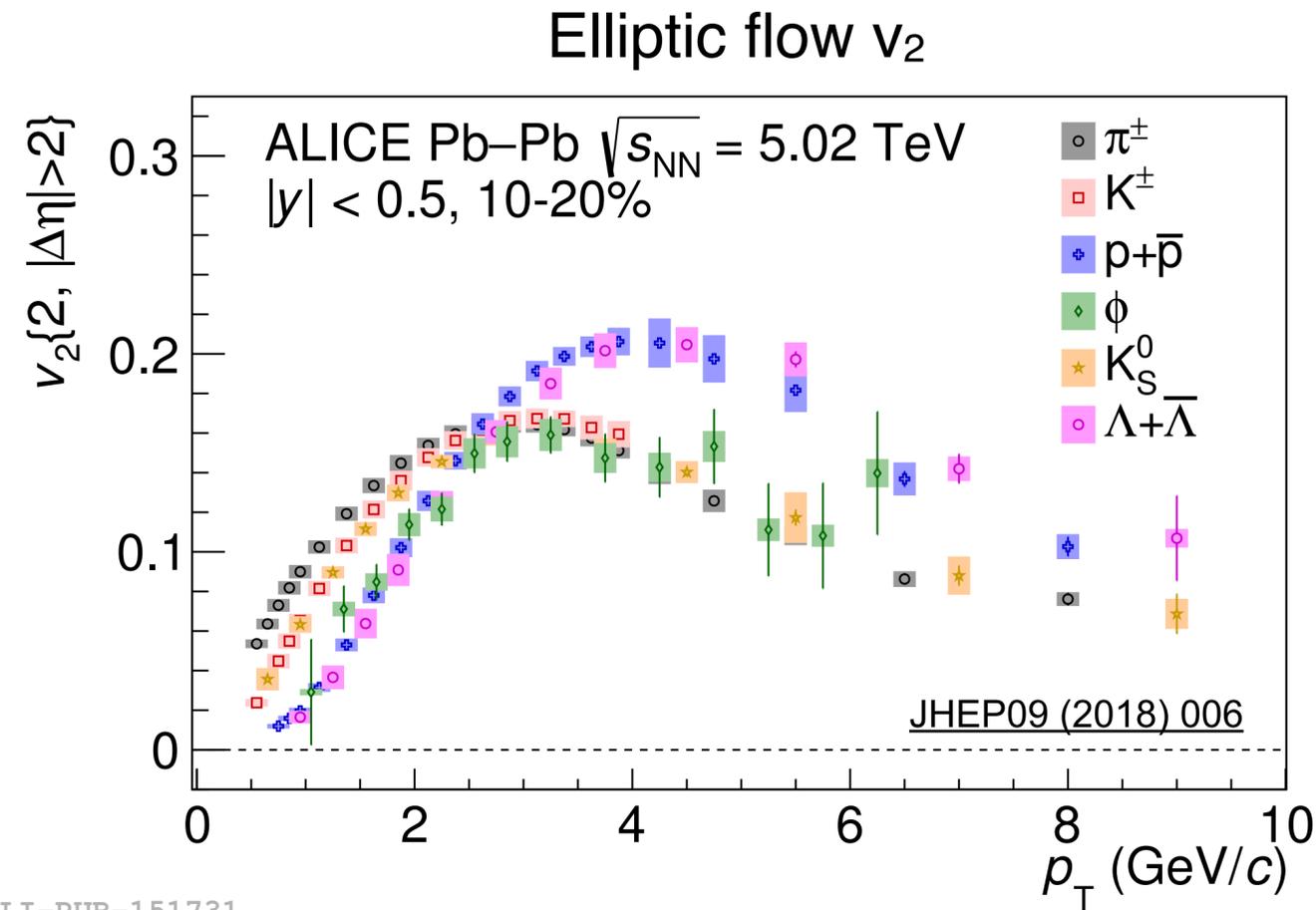
ALI-PUB-151731



Even heavy flavour hadrons flow !

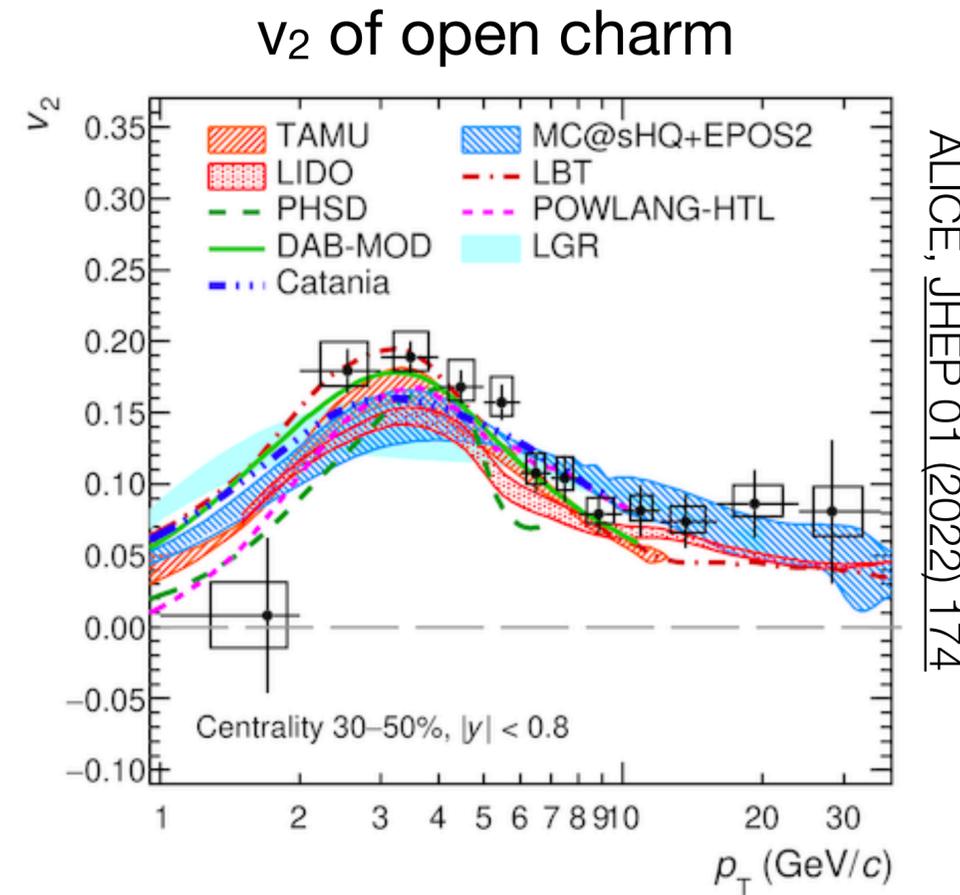
Mass-dependence of v_2 measures flow velocity

Anisotropic flow: initial state and QGP expansion

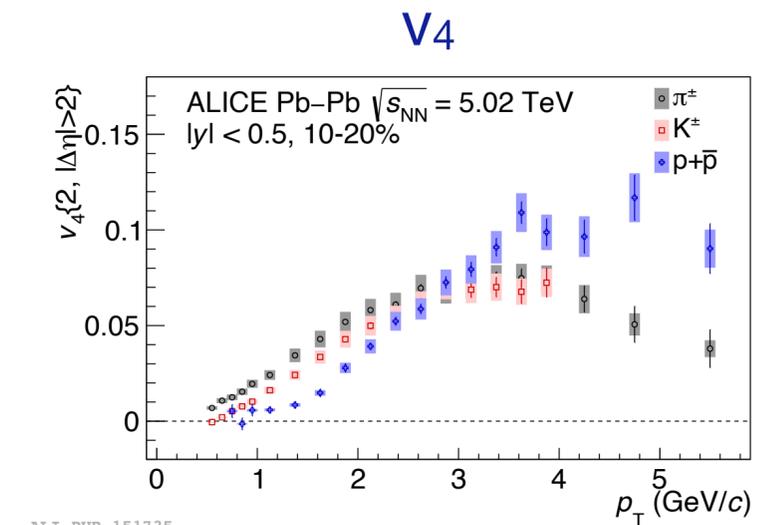
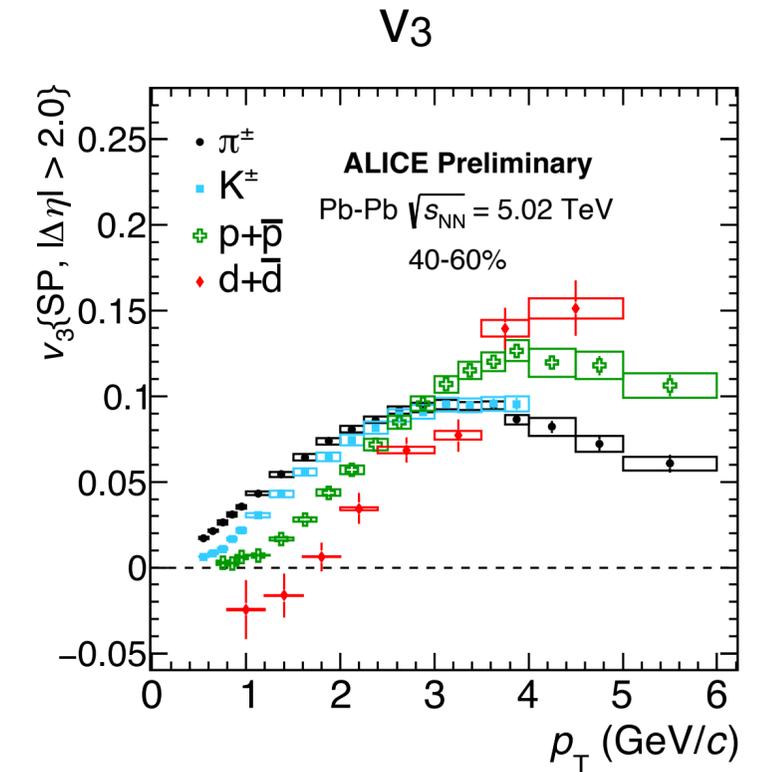


ALI-PUB-151731

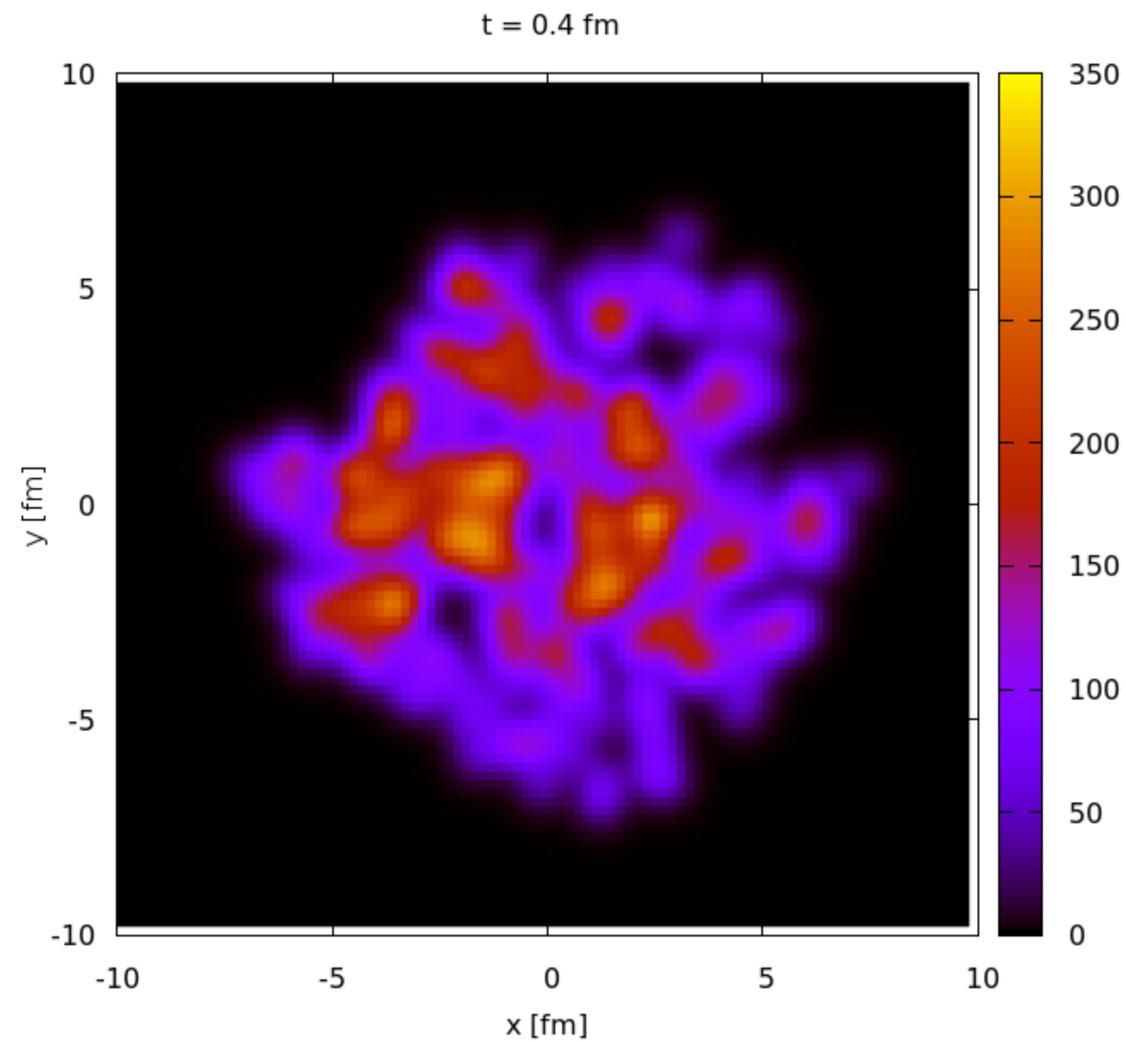
Mass-dependence of v_2 measures flow velocity



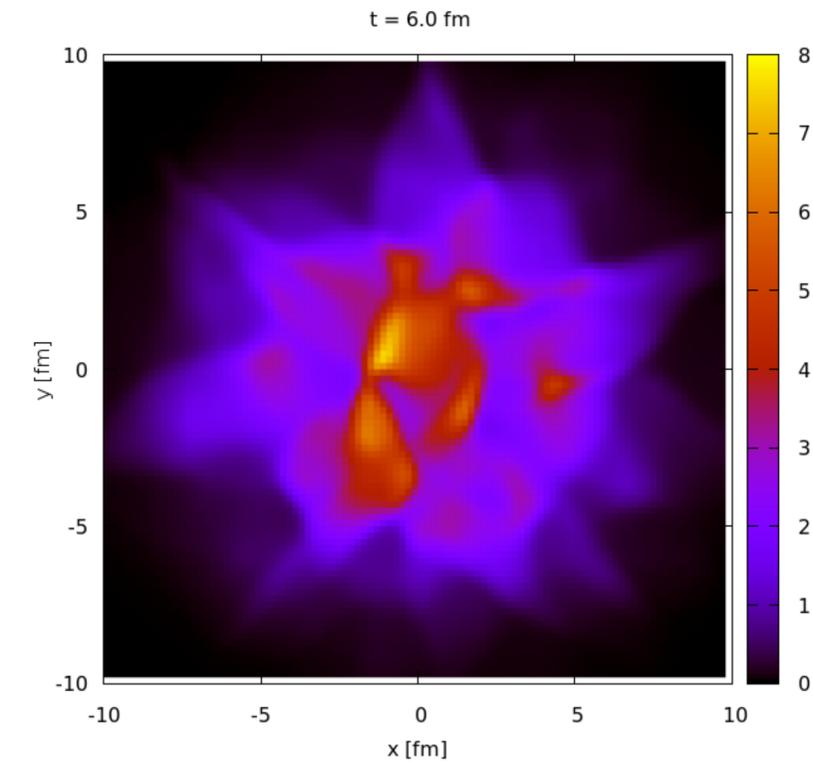
Even heavy flavour hadrons flow !



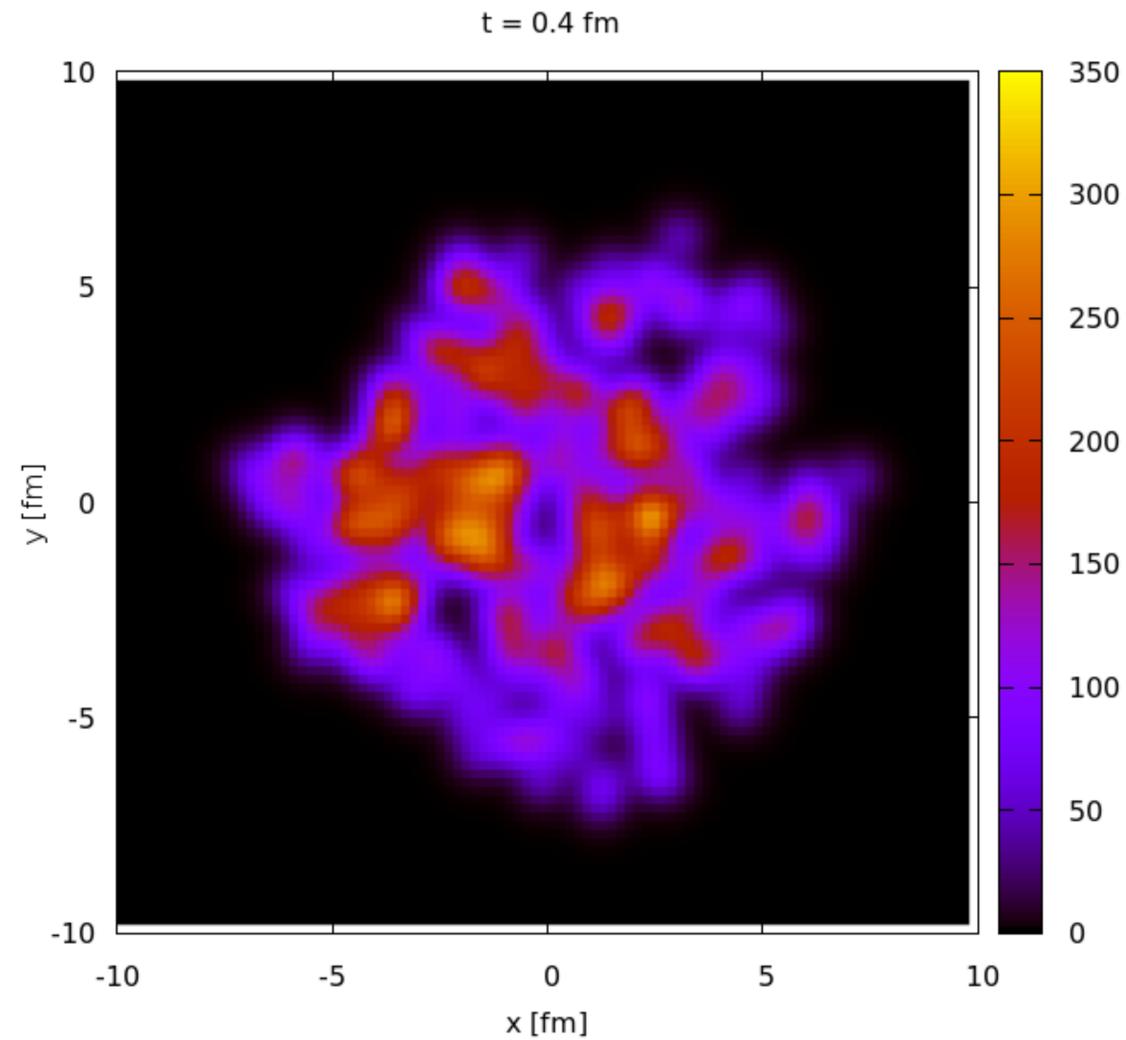
Azimuthal anisotropy: initial and final states



No viscosity
 $\eta/s = 0$



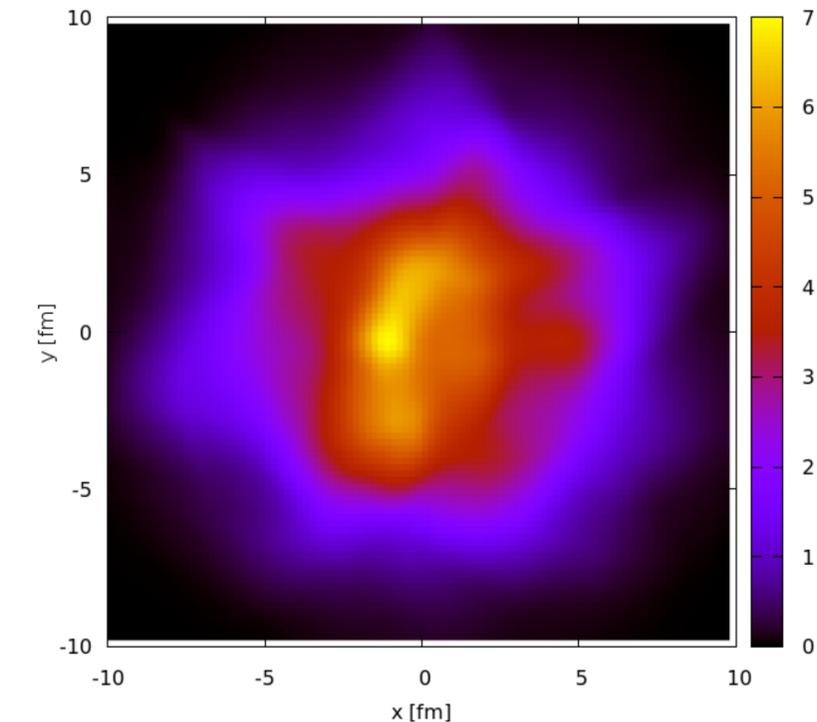
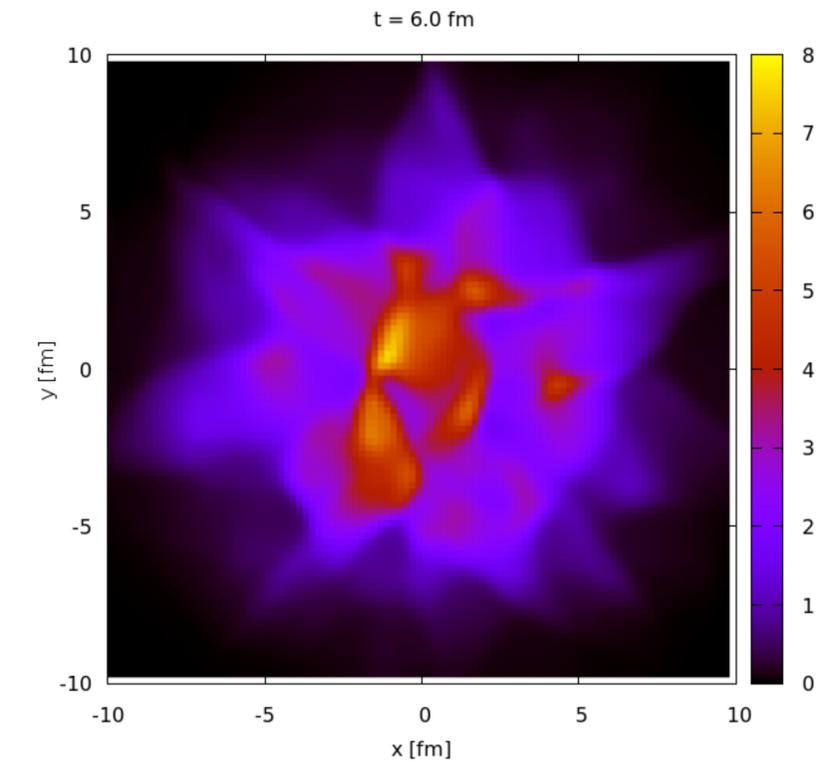
Azimuthal anisotropy: initial and final states



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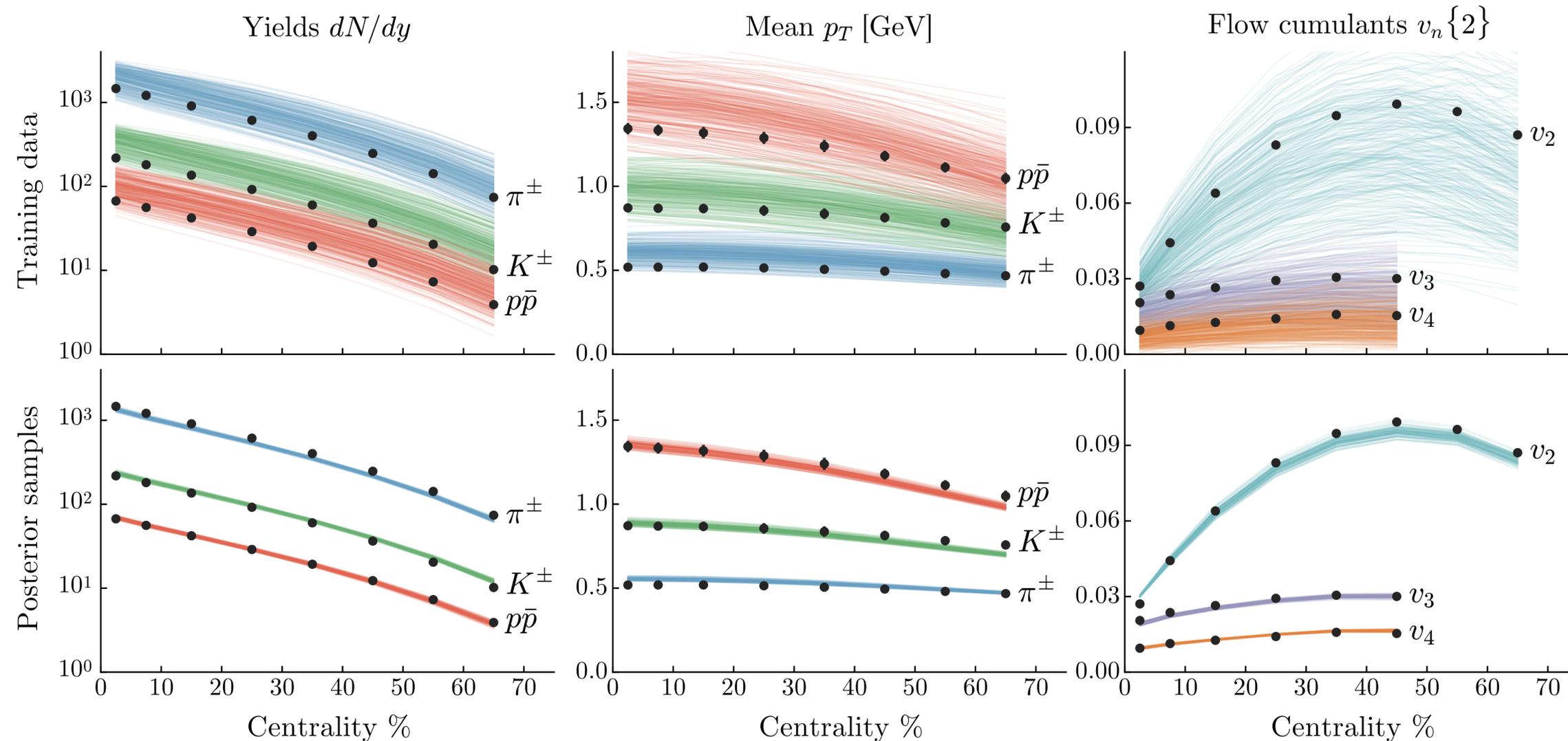
$\eta/s = 0.16$
Low viscosity



Constraining initial state and plasma properties simultaneously: Bayesian inference

J. E. Bernhard et al, arXiv: 1605.03954

Experimental input: yields, mean p_T and harmonic flow vs p_T



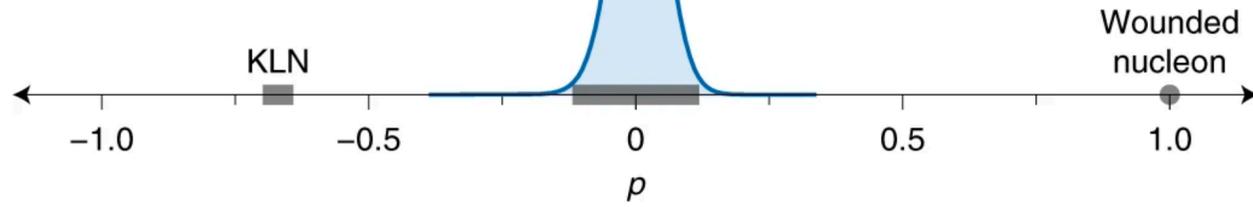
Model: initial anisotropies + medium response

Explores a large parameter space to investigate reliability/robustness of the modeling

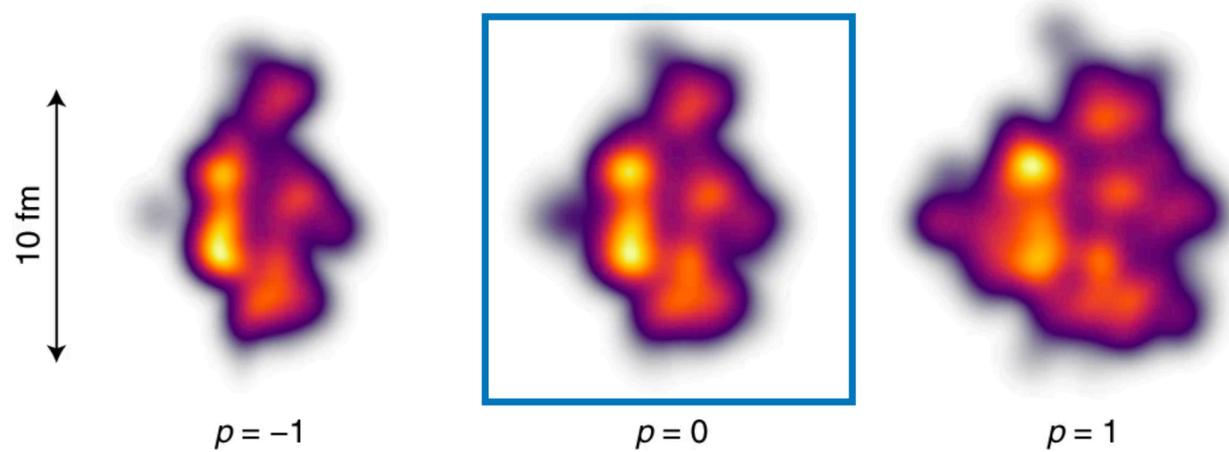
Bayesian analysis of flow: results

Initial state geometry

Calibrated to:
Pb-Pb 2.76 and 5.02 TeV

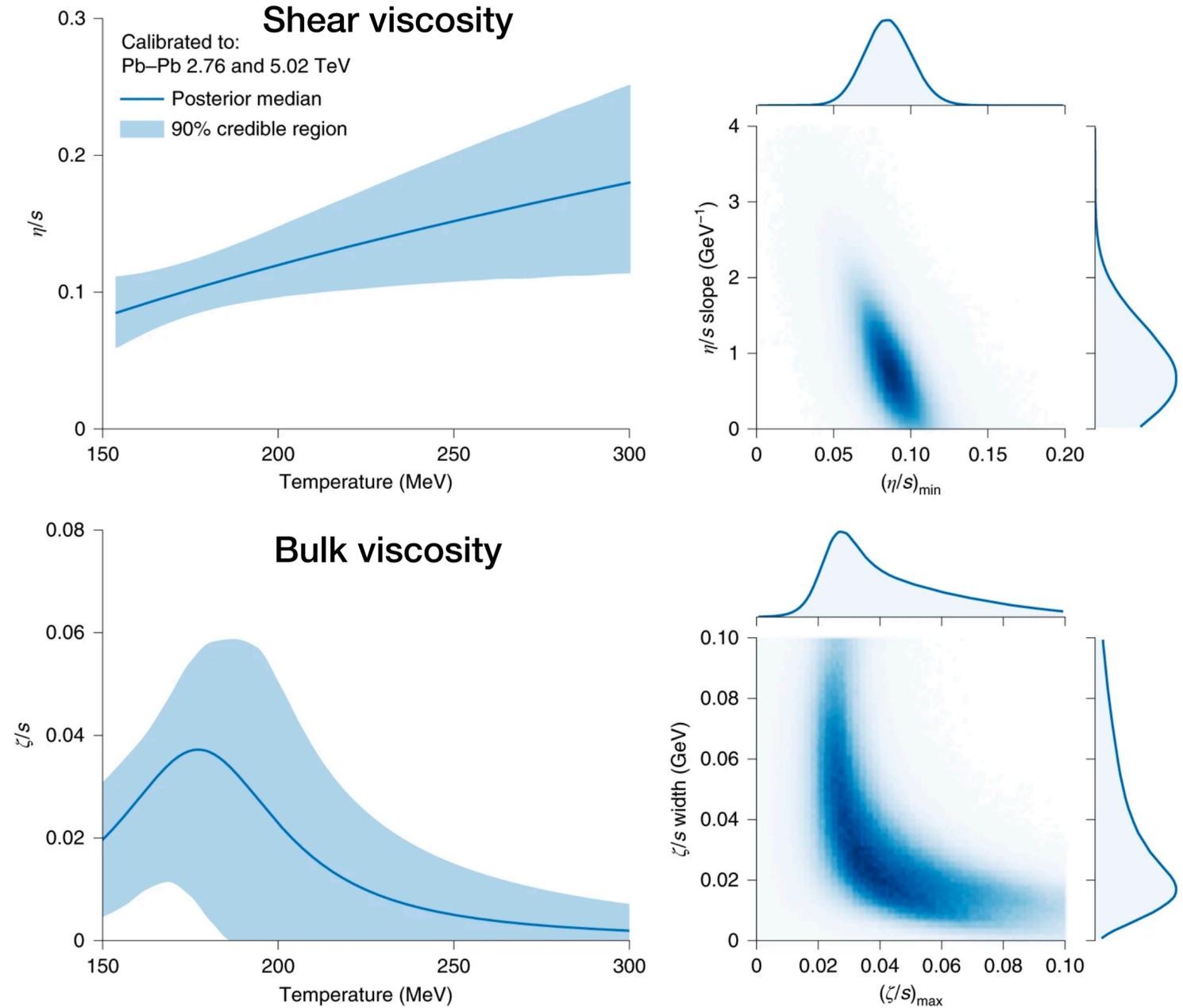


Energy density in transverse plane



Flow data provide information on initial geometry and viscosity of the QGP at the same time

Viscosity vs T



J. E. Bernhard et al, arXiv: 1605.03954

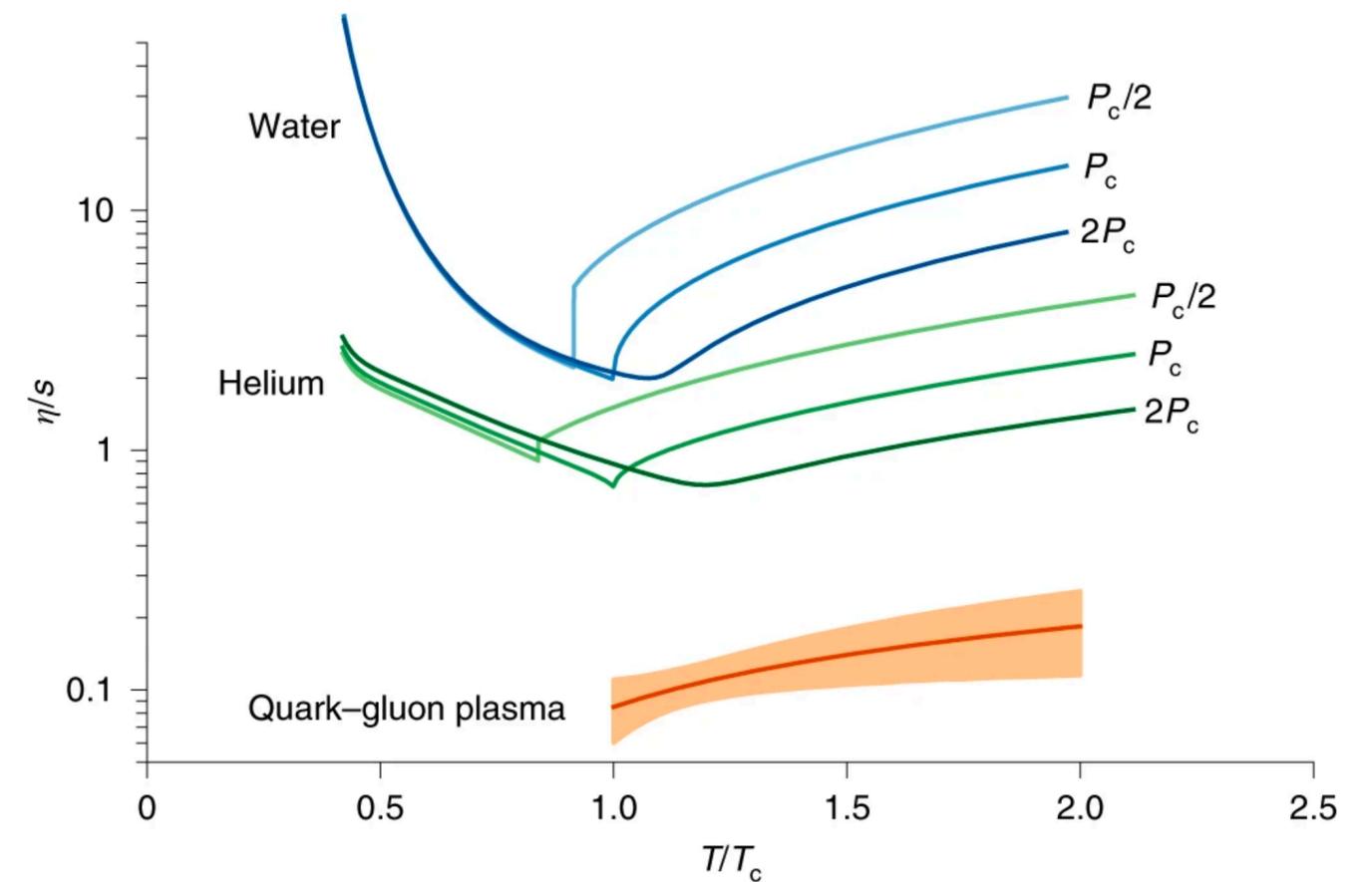
A global fit to anisotropic flow: main result

J. E. Bernhard et al, Nature
Physics 15, 1113–1117,
arXiv: 1605.03954

QGP has a very small ‘specific viscosity’
small mean free path

Viscosity close to fundamental lower bound

Comparison to well-known liquids

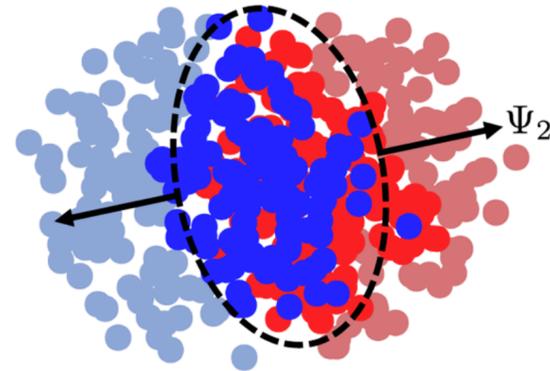


$$\eta = \frac{1}{3} n \bar{p} \lambda$$

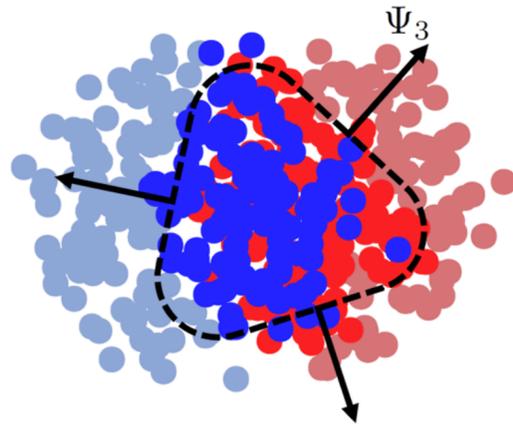
Initial state geometry: event plane correlations

ALICE, [arXiv:2302.01234](https://arxiv.org/abs/2302.01234)

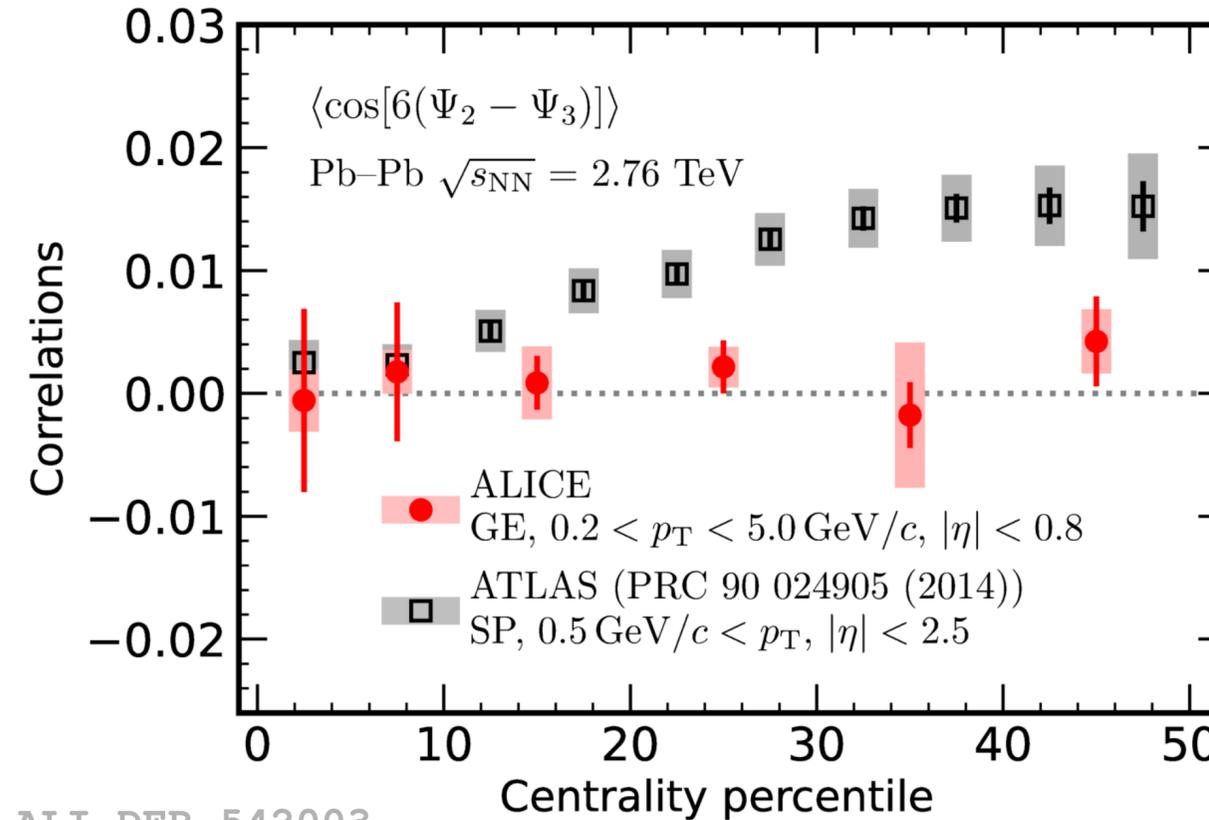
Elliptic deformation



Triangular deformation

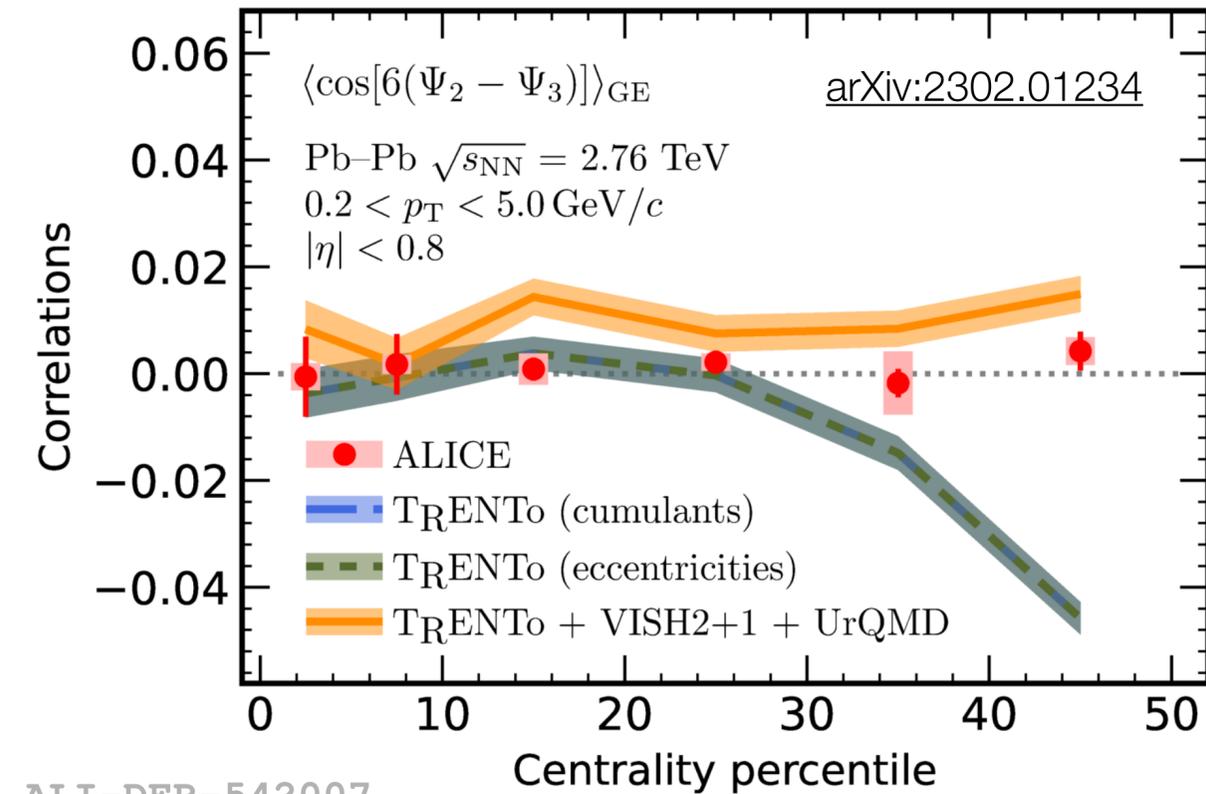


Comparison to previous results



ALI-DER-542003

Comparison to theory



ALI-DER-542007

Correlations between symmetry plane orientations

- New method: reduced sensitivity to numerical fluctuations
- **No significant correlation between Ψ_2 and Ψ_3**
- New results more in line with expectations

Chemical freeze-out

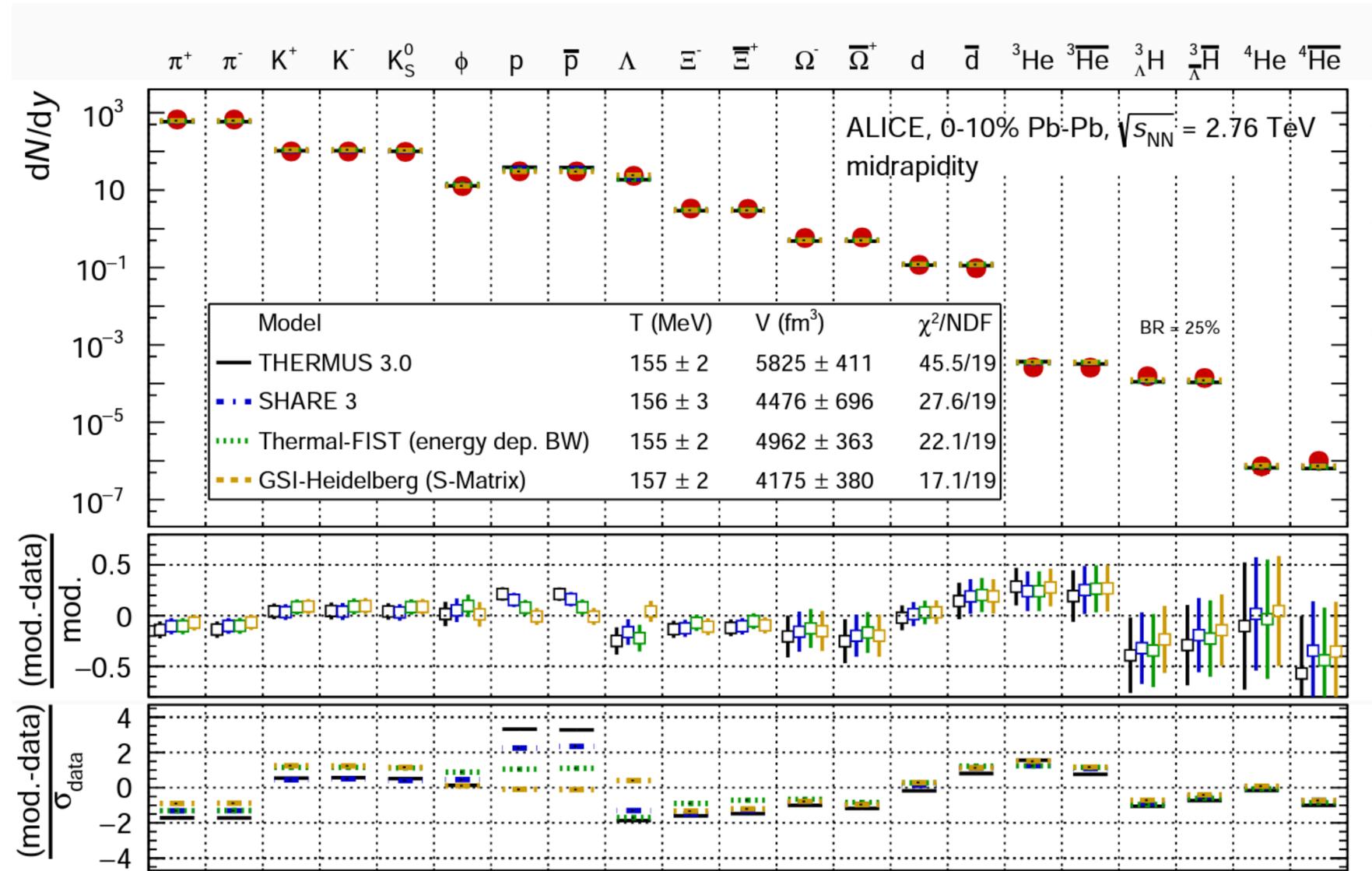
Chemical freeze-out determines hadron yields
 — end of inelastic collisions

- Hadron yields follow thermal distribution with $T = T_c = 155$ MeV

$$N = (2J + 1) e^{-m/T}$$

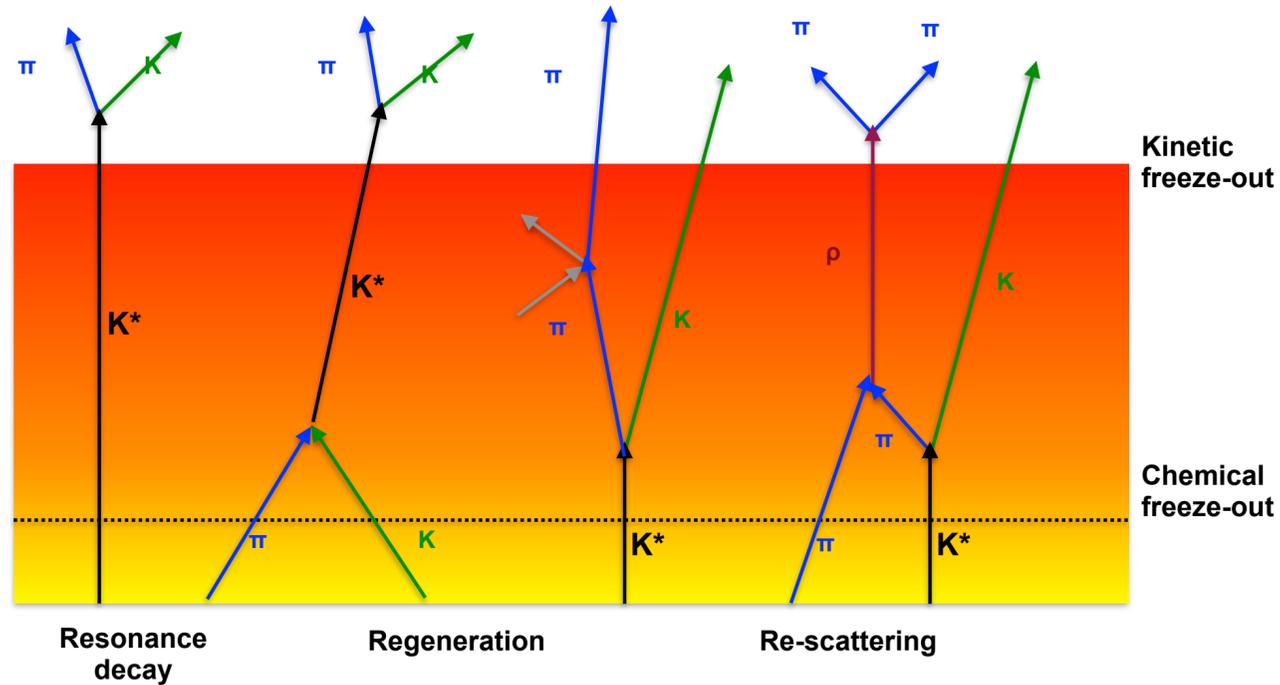
- Chemical freeze-out at phase transition temperature:
no inelastic collisions after phase transition

Hadron yields compared to thermal model calculation



A journey through QCD, ALICE, arXiv:2211.04384

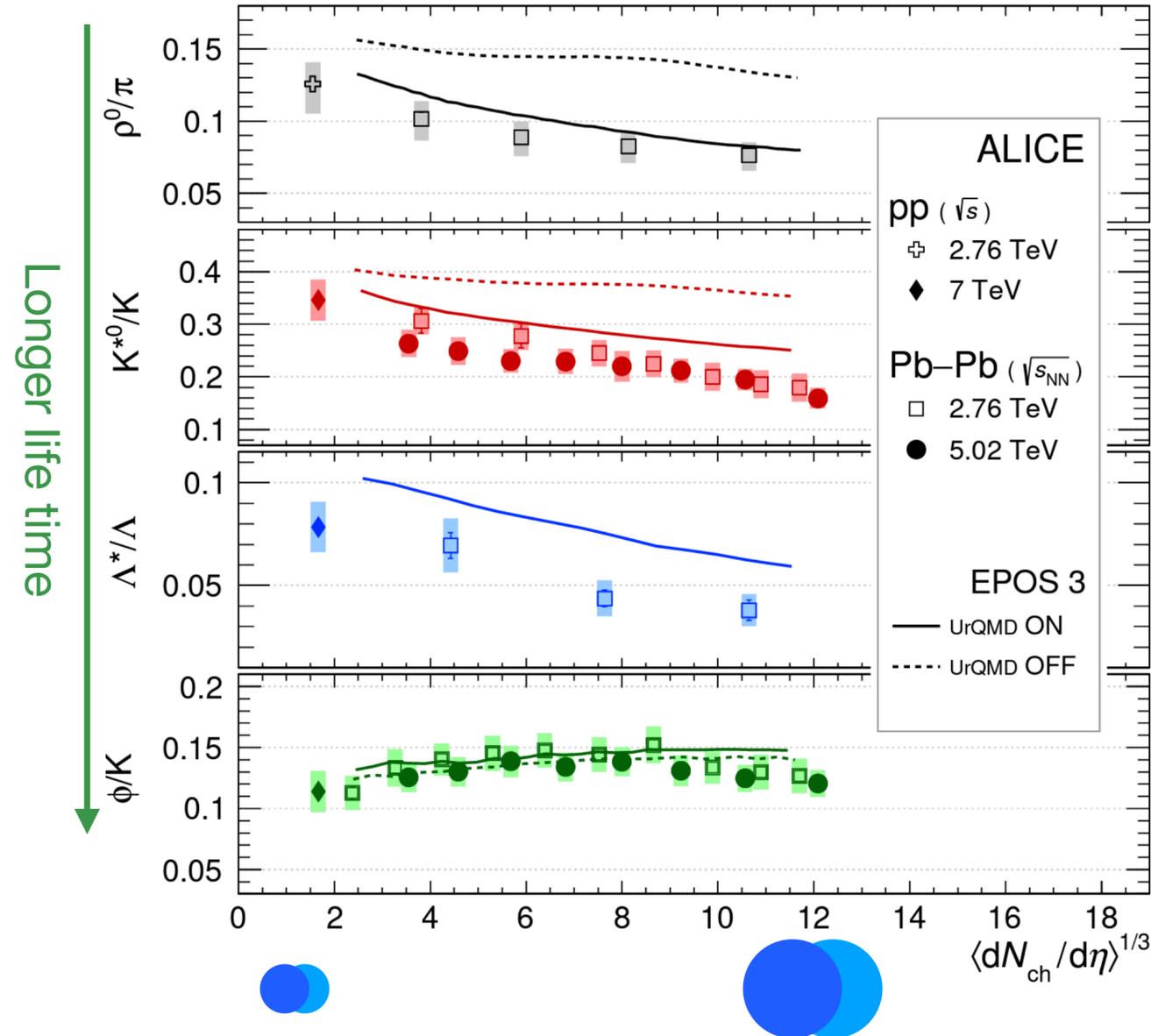
Elastic collisions in the hadronic phase: resonance yields



Re-scattering of decay daughters: signal loss

Regeneration: resonances formed in hadron scattering

Resonance yields vs system size



A journey through QCD, ALICE, arXiv:2211.04384

ρ, K^*, Λ^* reduced yield: final state scattering of decay particles

ϕ : longer life time, yield not affected

Light nuclei: a sensitive probe of the hadronic phase

Hadronic collisions at LHC produce light nuclei
d, t, ^3He , ^4He etc

- Small binding energy $O(\text{MeV}) \ll T$ of hadronic phase
- Expect large break-up probability

Formation by coalescence of protons and neutrons?

- Expect yield $\propto \rho_p \rho_n$
- Model calculations use Wigner function formalism

JUNE 1, 1932

PHYSICAL REVIEW

VOLUME 40

On the Quantum Correction For Thermodynamic Equilibrium

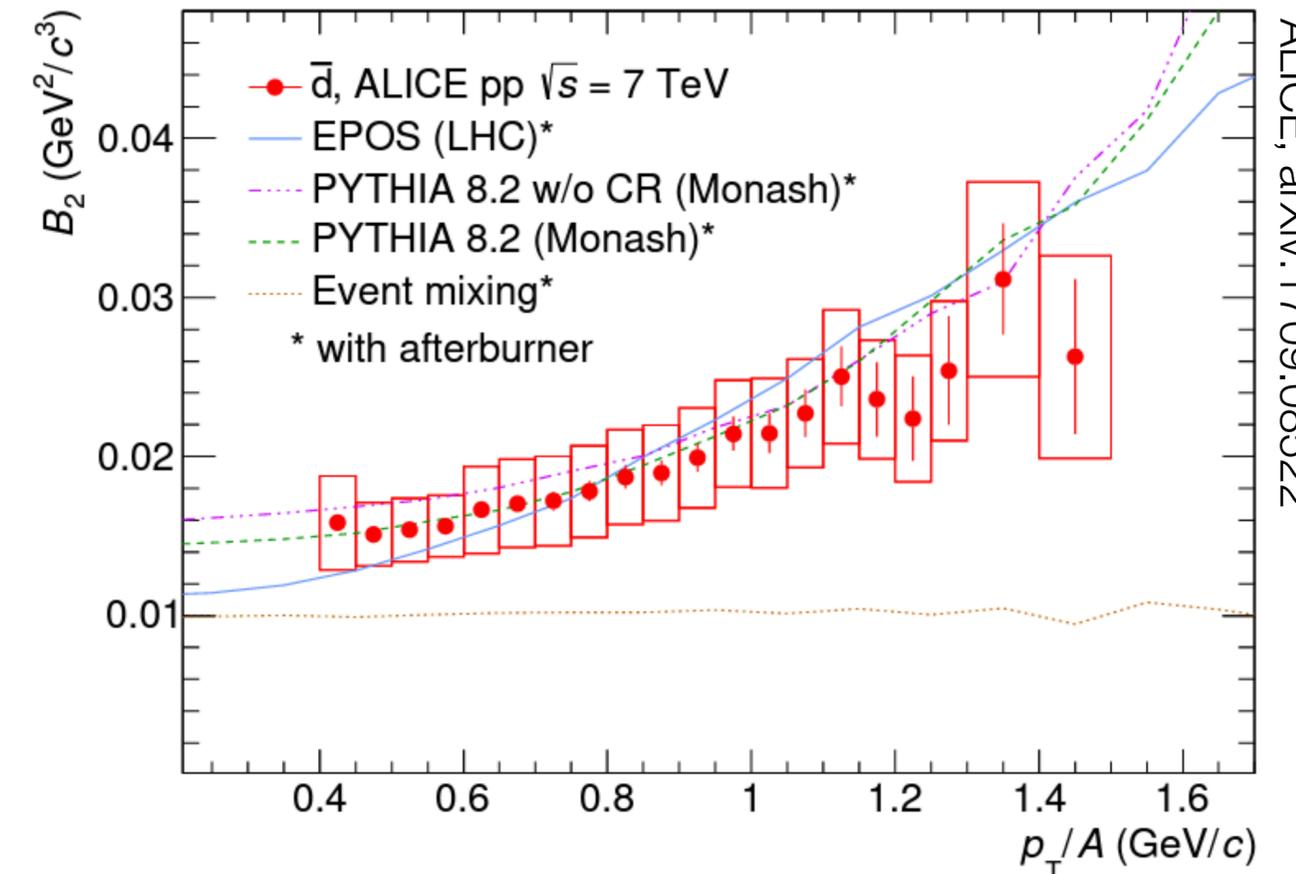
By E. WIGNER

Department of Physics, Princeton University

(Received March 14, 1932)

The probability of a configuration is given in classical theory by the Boltzmann formula $\exp[-V/hT]$ where V is the potential energy of this configuration. For high temperatures this of course also holds in quantum theory. For lower temperatures, however, a correction term has to be introduced, which can be developed into a power series of \hbar . The formula is developed for this correction by means of a probability function and the result discussed.

deuteron coalescence parameter in pp



$$B_2 = E_A \frac{d^3N}{dp_A^3}(p_A) \bigg/ E_p \frac{d^3N}{dp_p^3}(p_p) E_n \frac{d^3N}{dp_n^3}(p_n)$$

$$p_n = p_p \quad p_A = 2p_p$$

The Wigner function (quantum density matrix formalism)

On the Quantum Correction For Thermodynamic Equilibrium

By E. WIGNER

Department of Physics, Princeton University

(Received March 14, 1932)

It does not seem to be easy to make explicit calculations with the form (4) of the mean value. One may resort therefore to the following method.

If a wave function $\psi(x_1 \cdots x_n)$ is given one may build the following expression²

$$\begin{aligned}
 P(x_1, \cdots, x_n; p_1, \cdots, p_n) \\
 = \left(\frac{1}{h\pi}\right)^n \int_{-\infty}^{\infty} \cdots \int dy_1 \cdots dy_n \psi(x_1 + y_1 \cdots x_n + y_n)^* \\
 \psi(x_1 - y_1 \cdots x_n - y_n) e^{2i(p_1 y_1 + \cdots + p_n y_n)/h} \quad (5)
 \end{aligned}$$

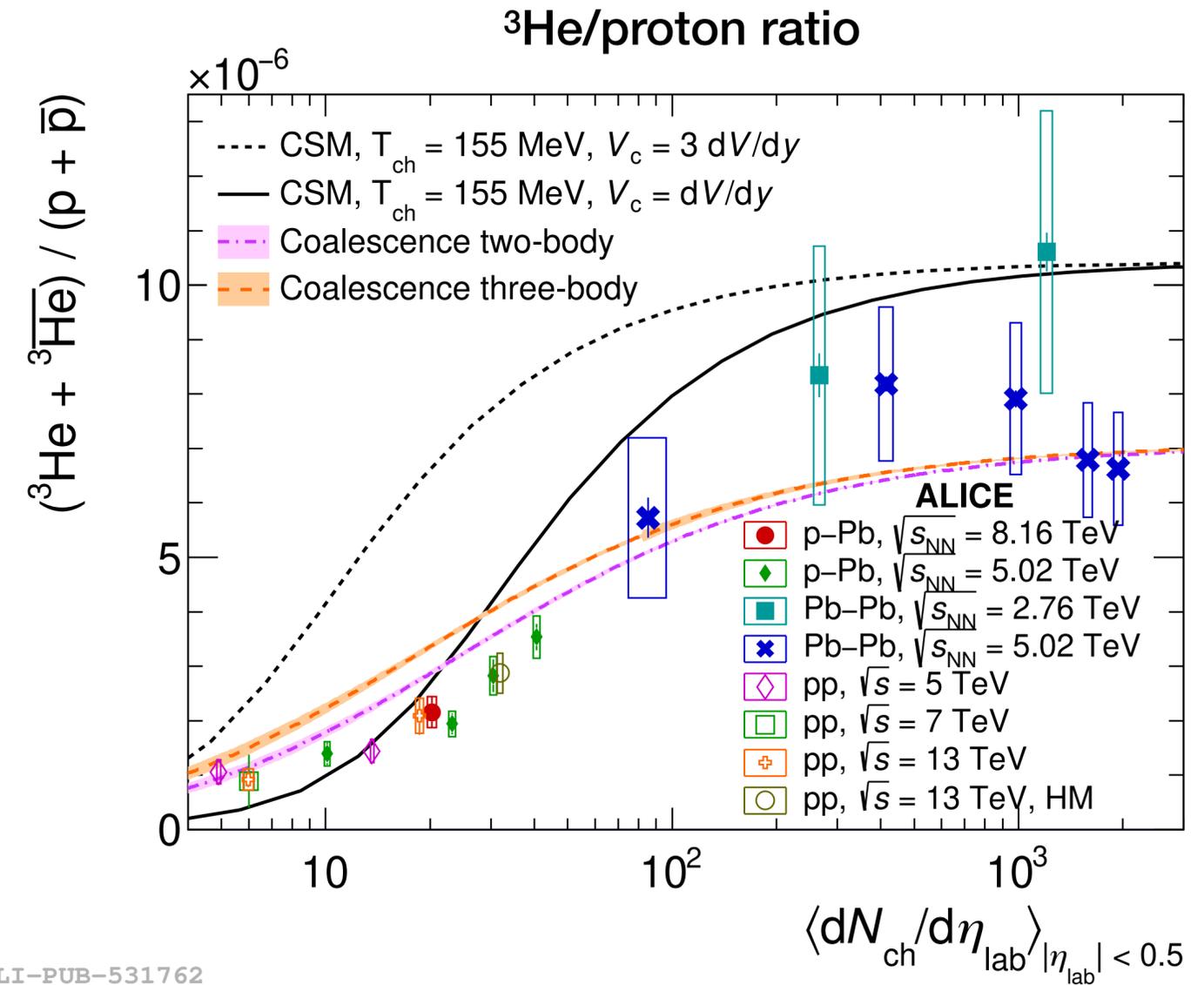
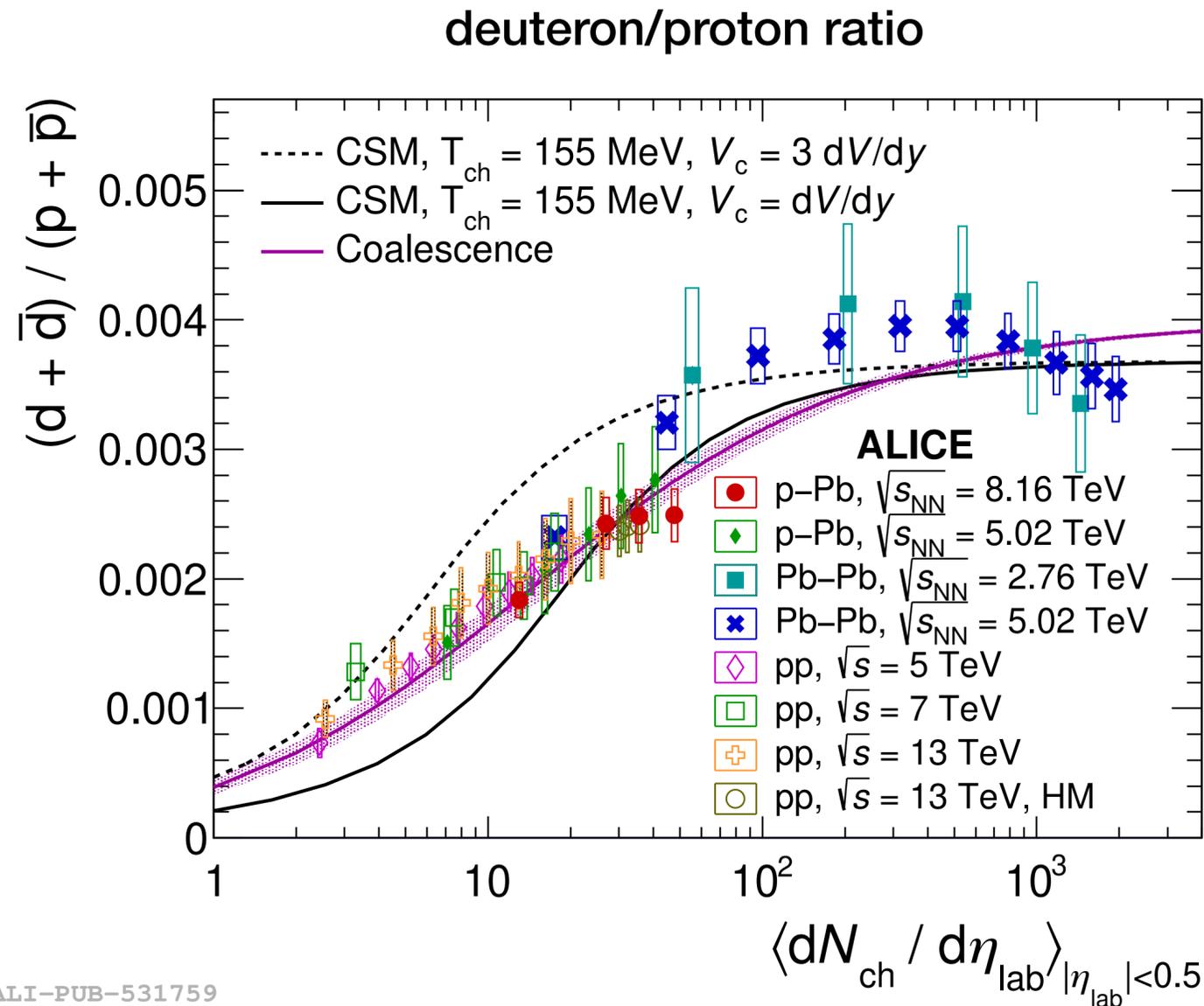
and call it the probability-function of the simultaneous values of $x_1 \cdots x_n$ for the coordinates and $p_1 \cdots p_n$ for the momenta. In (5), as throughout

Of course $P(x_1, \cdots, x_n; p_1, \cdots, p_n)$ cannot be really interpreted as the simultaneous probability for coordinates and momenta, as is clear from the fact, that it may take negative values. But of course this must not hinder the use of it in calculations as an auxiliary function which obeys many relations we would expect from such a probability. It should be noted, furthermore,

Coalescence models for nuclear collisions/QGP combine **proximity conditions in coordinate and momentum space**

- Explored for nuclei by e.g. R Scheibl, U Heinz
- Extended to parton coalescence by J Zimányi, P Lévai, T Csörgö, T S Biró, V Greco, C M Ko, R Fries, R Hwa, and others

Production of light nuclei: thermal model vs coalescence



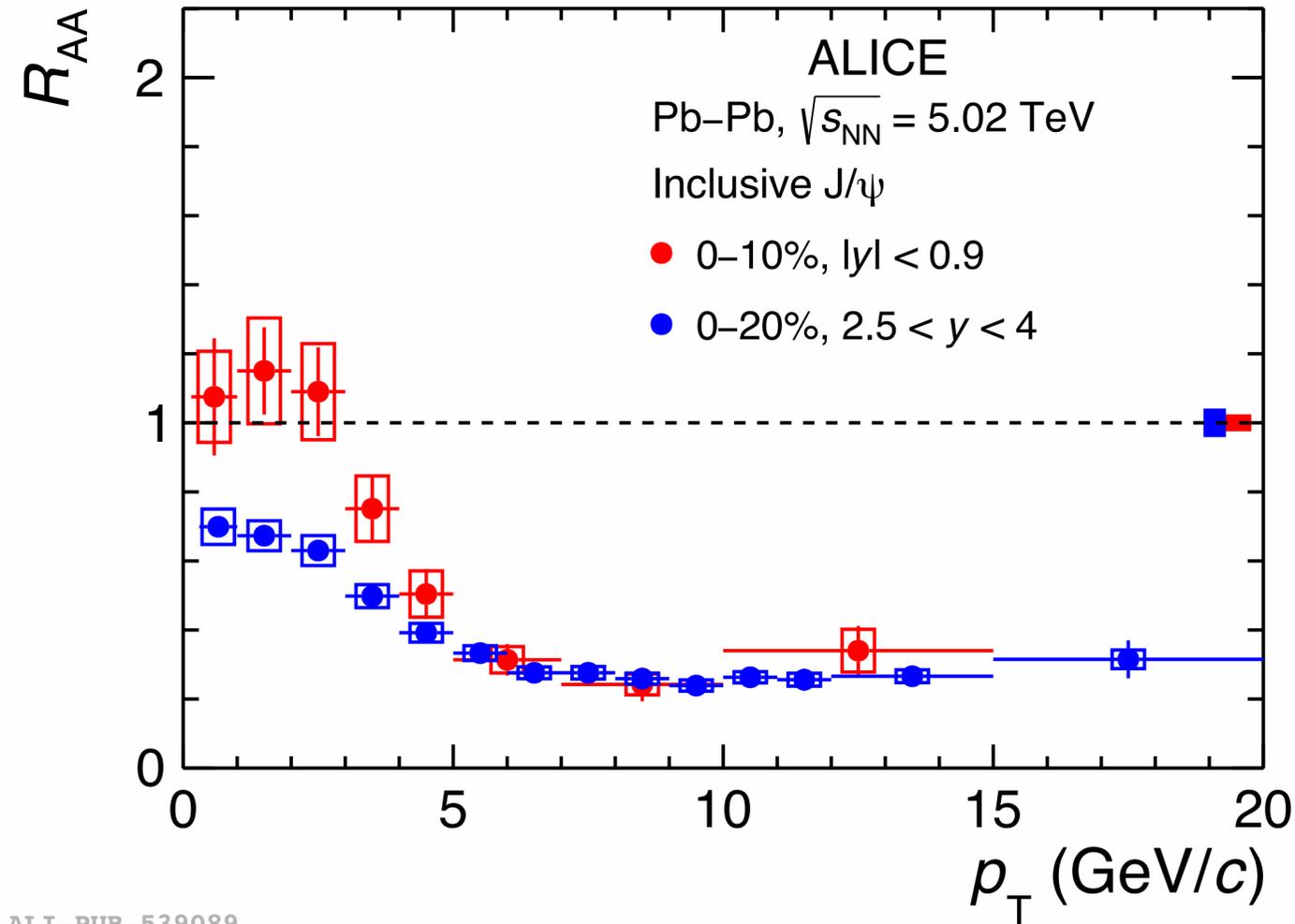
arXiv:2212.04777

Thermal and coalescence calculations give similar result for compact states;
clear differences for larger states

J/ψ: melting and regeneration at the parton level

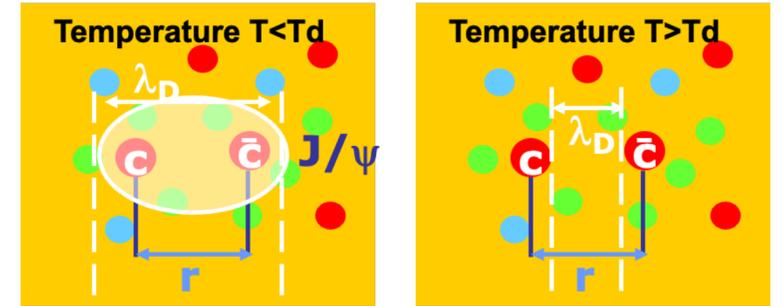
$$R_{AA} = \frac{dN/dp_T|_{AA}}{\langle N_{coll} \rangle dN/dp_T|_{pp}}$$

Nuclear modification factor



ALI-PUB-539089

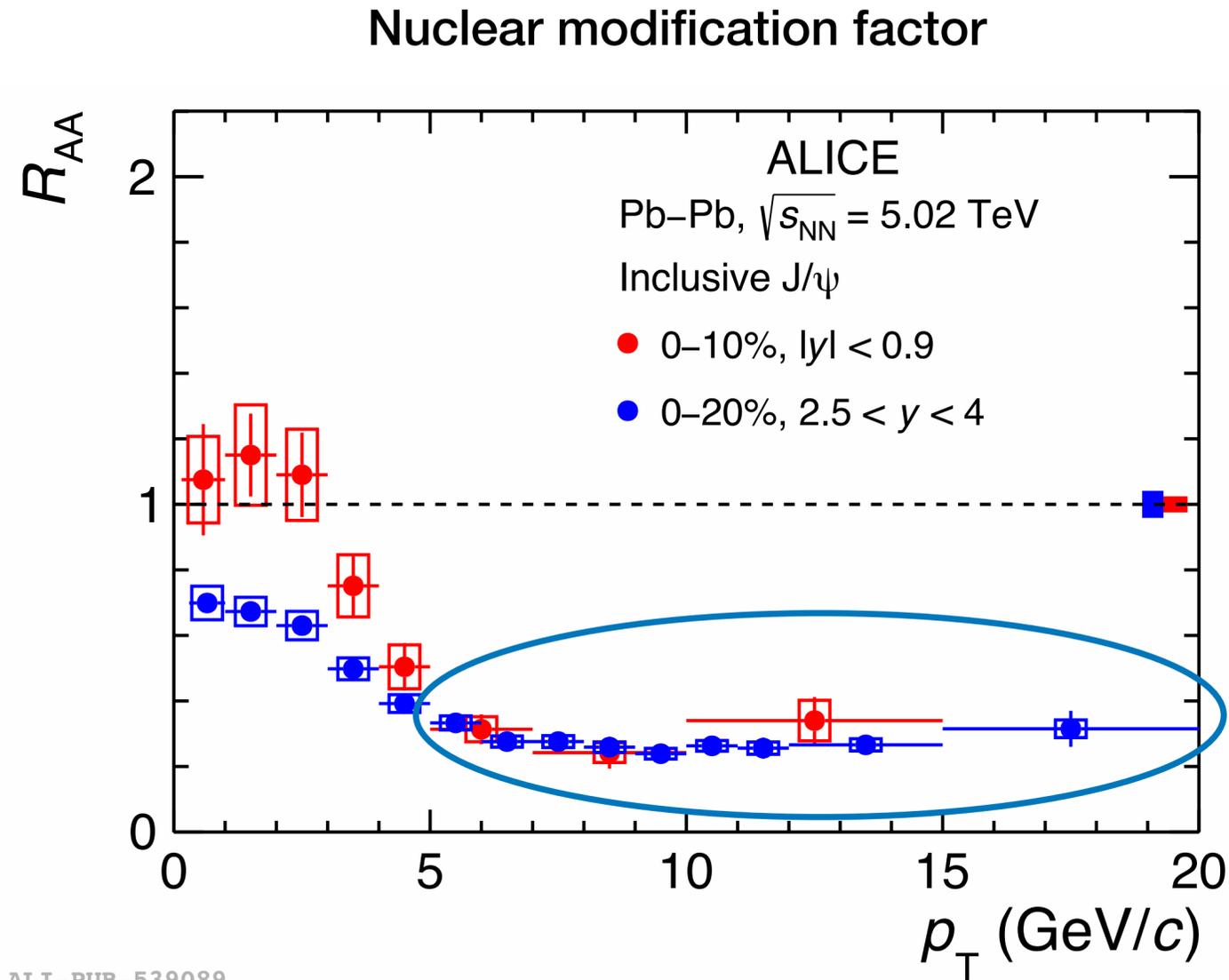
J/ψ: bound state
charm and anti-charm quark



Binding force screened when
 $r > \lambda_d$

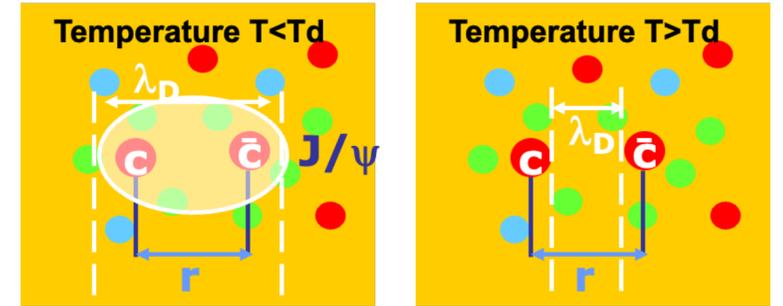
J/ψ: melting and regeneration at the parton level

$$R_{AA} = \frac{dN/dp_T|_{AA}}{\langle N_{coll} \rangle dN/dp_T|_{pp}}$$



ALI-PUB-539089

J/ψ: bound state
charm and anti-charm quark

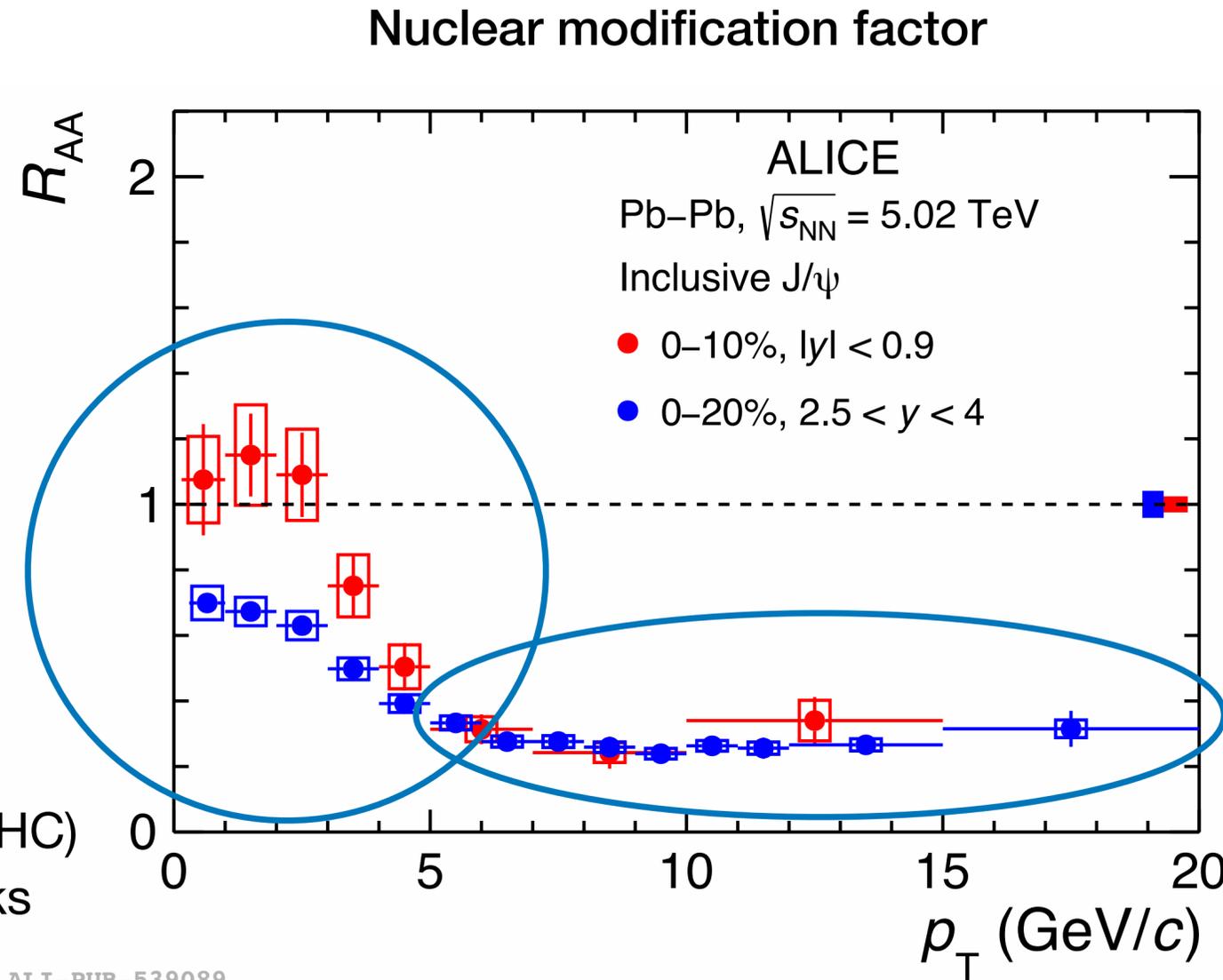


Binding force screened when
 $r > \lambda_d$

High p_T - low density:
quarkonia dissociate in the QGP

J/ψ: melting and regeneration at the parton level

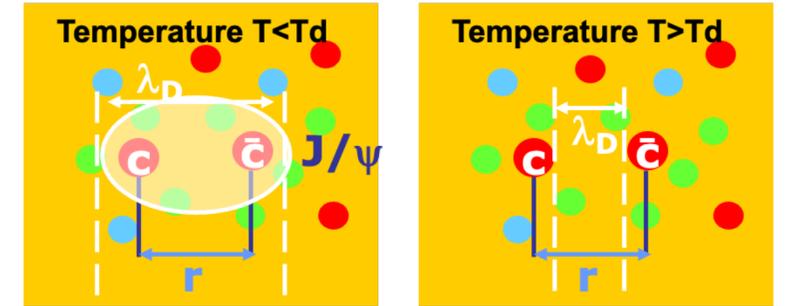
$$R_{AA} = \frac{dN/dp_T|_{AA}}{\langle N_{coll} \rangle dN/dp_T|_{pp}}$$



Low p_T :
large density of charm quarks (at LHC)
regeneration: coalescence of quarks

ALI-PUB-539089

J/ψ: bound state
charm and anti-charm quark

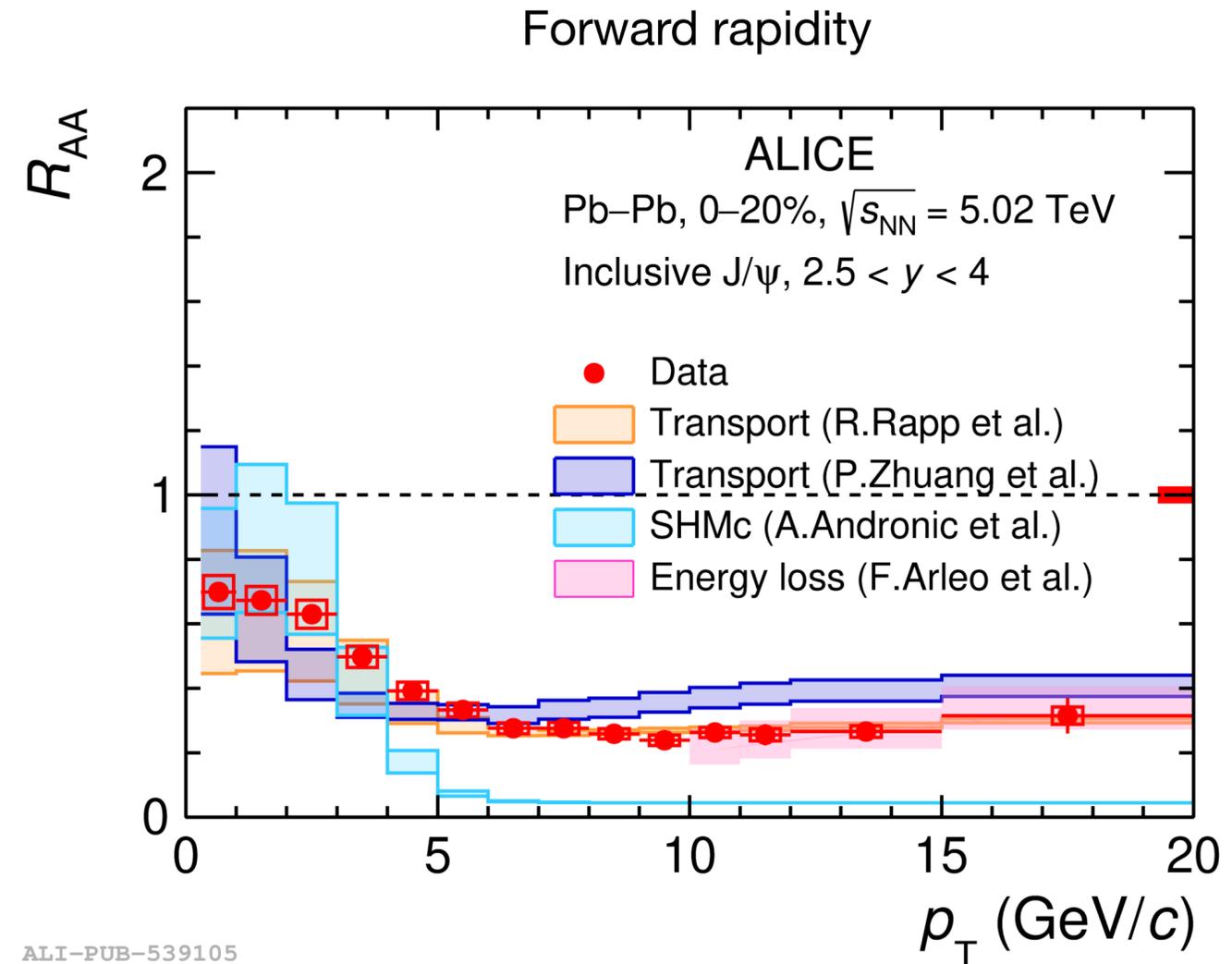
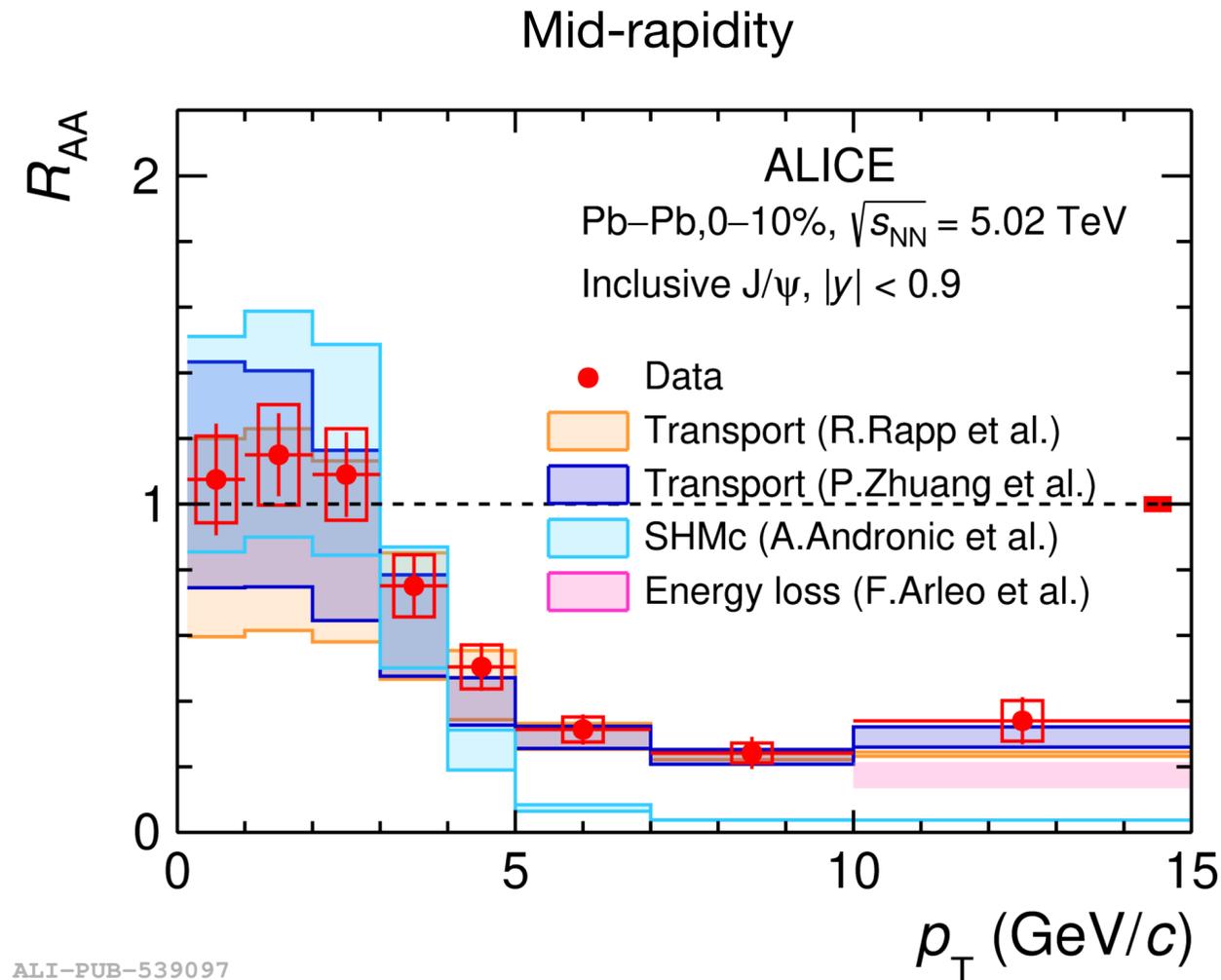


Binding force screened when
 $r > \lambda_d$

High p_T - low density:
quarkonia dissociate in the QGP

J/ψ production in Pb-Pb collisions: melting and recombination

arXiv:2303.13361



$$R_{AA} = \frac{dN/dp_T|_{AA}}{\langle N_{coll} \rangle dN/dp_T|_{pp}}$$

- Balance between melting and recombination at low p_T
- Rapidity dependence exposes density dependence

Back to the earliest stages: direct photon production

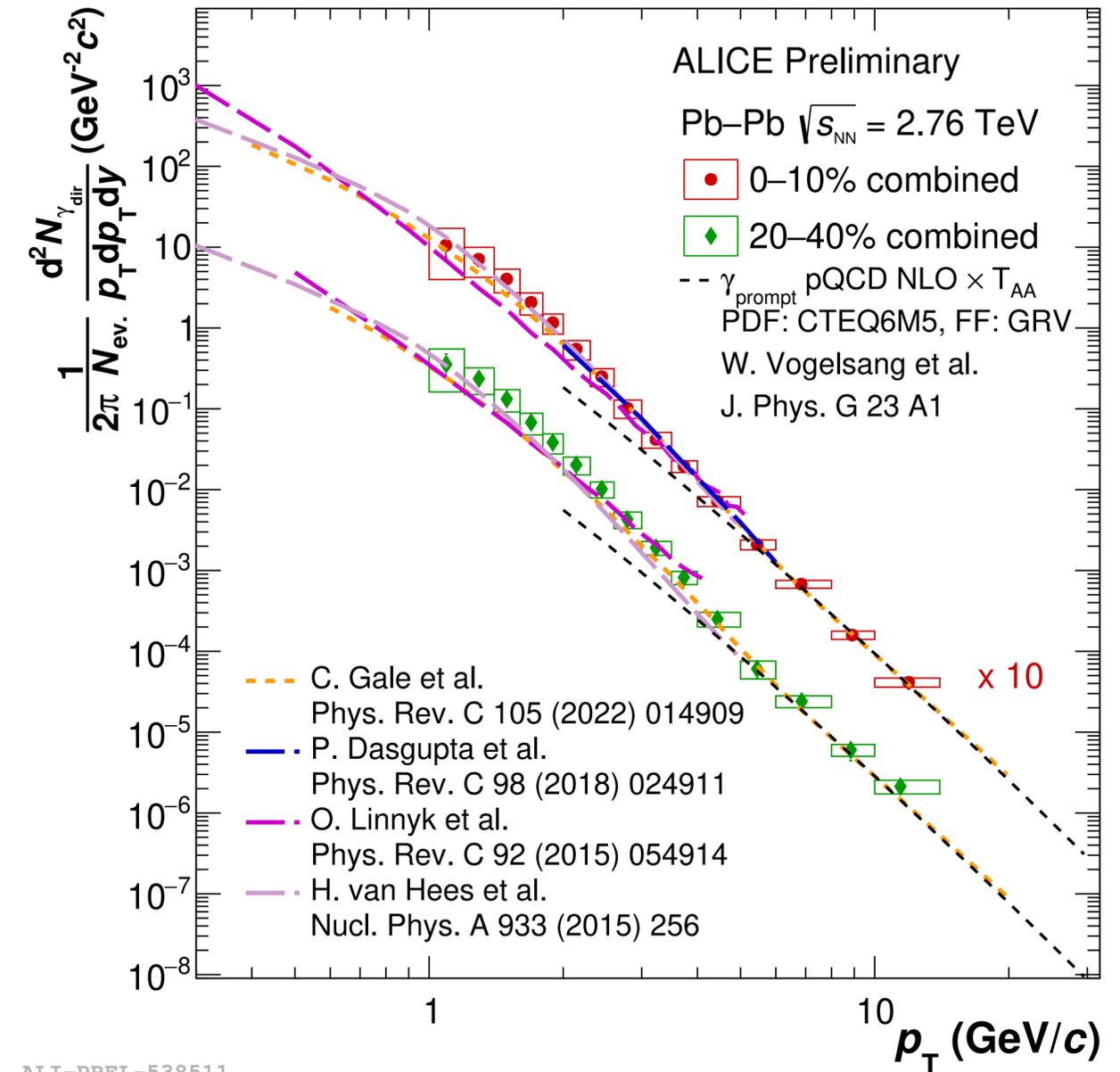
Large background: decay photons from π^0 , η , ...

⇒ Challenging measurement

Main sources:

- High p_T : hard scattering; quark-gluon Compton process
- Low p_T : thermal radiation

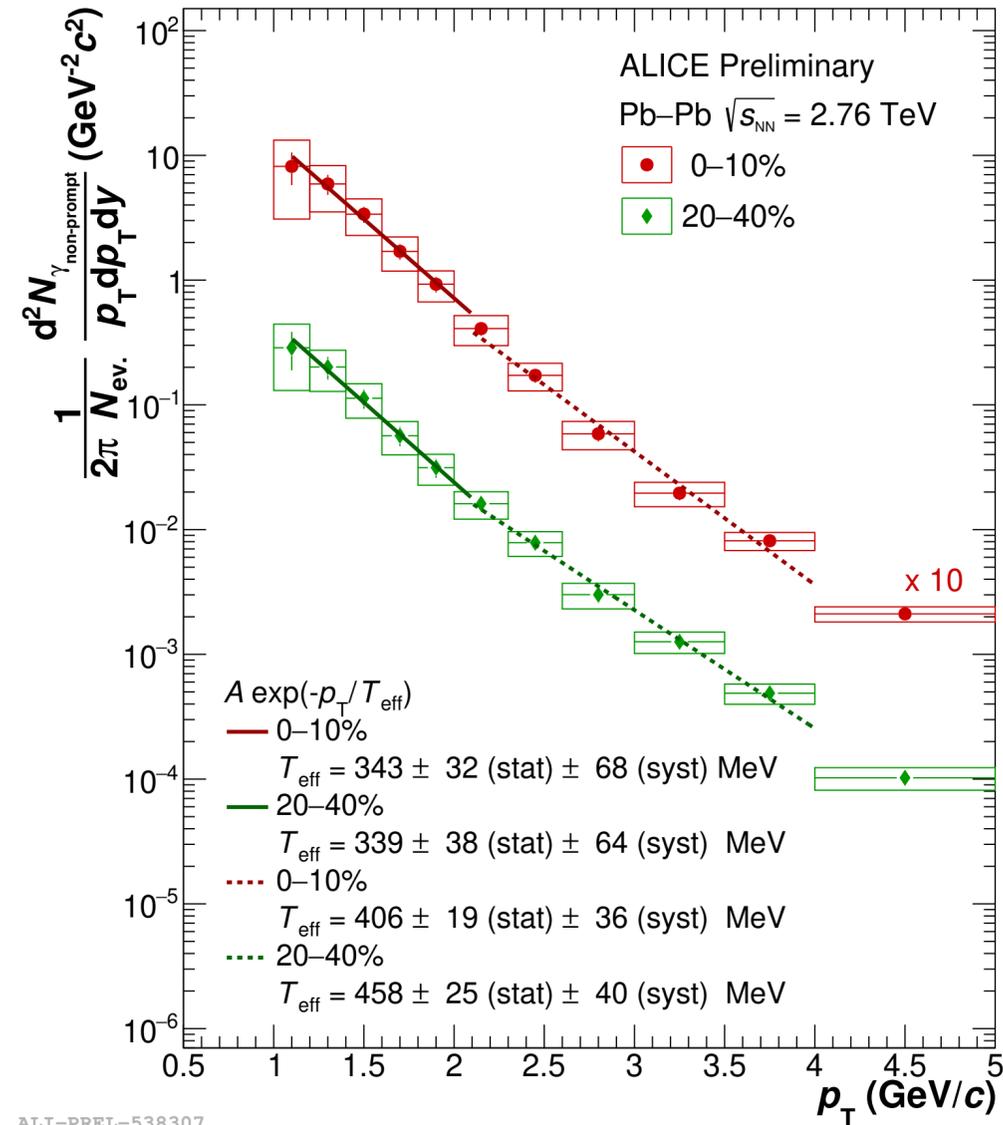
Excess at low p_T : thermal photons



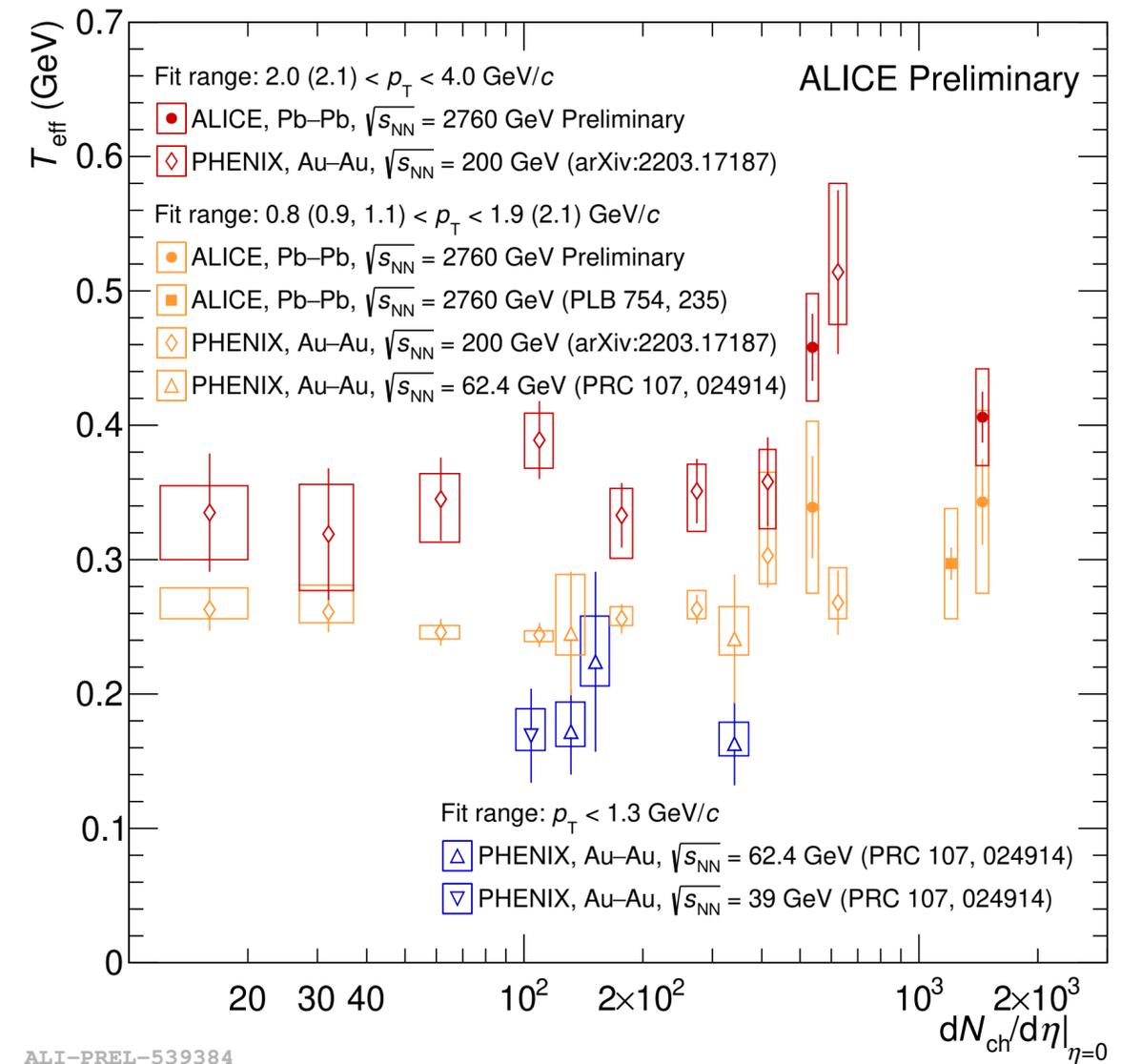
ALICE, PLB 754, 235

Direct photon excess: thermal production

Direct photon excess spectrum



Spectral slope: apparent temperature



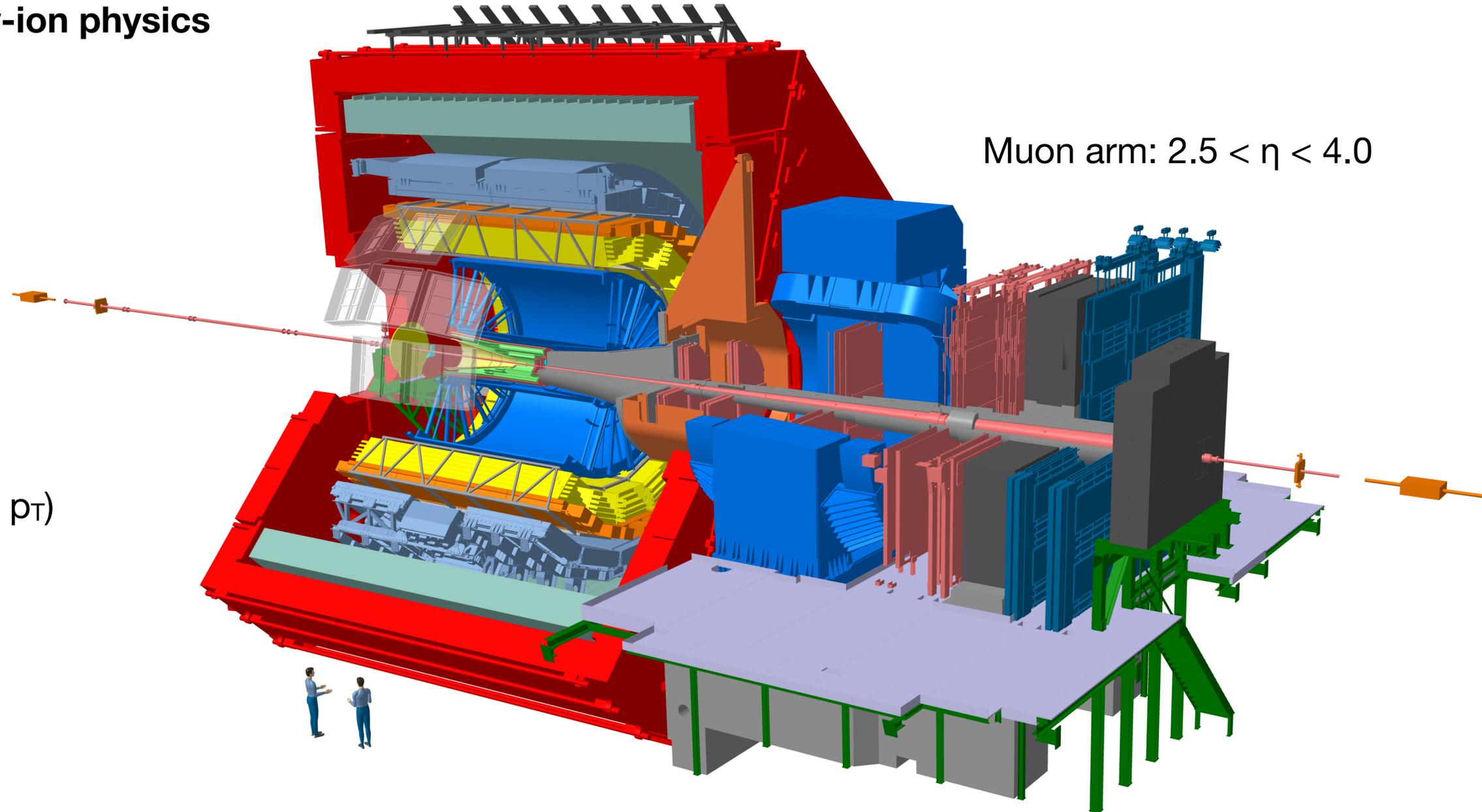
Thermal emission visible for mid-central and central events

Apparent temperature larger at LHC than RHIC
Absolute temperature depends on blue shift

The ALICE detector

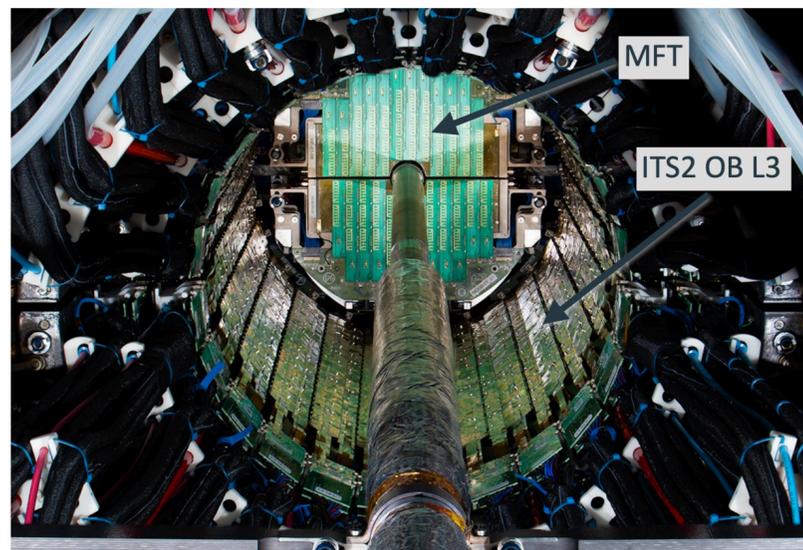
A general purpose detector for heavy-ion physics

- Low material budget
- High-resolution tracking
 - Silicon tracker
 - Time projection chamber
- Particle identification
 - TPC dE/dx
 - Time of Flight
 - Ring-imaging Cherenkov (high p_T)
 - Muon ID (forward)
 - Transition radiation detector
- EM calorimeters

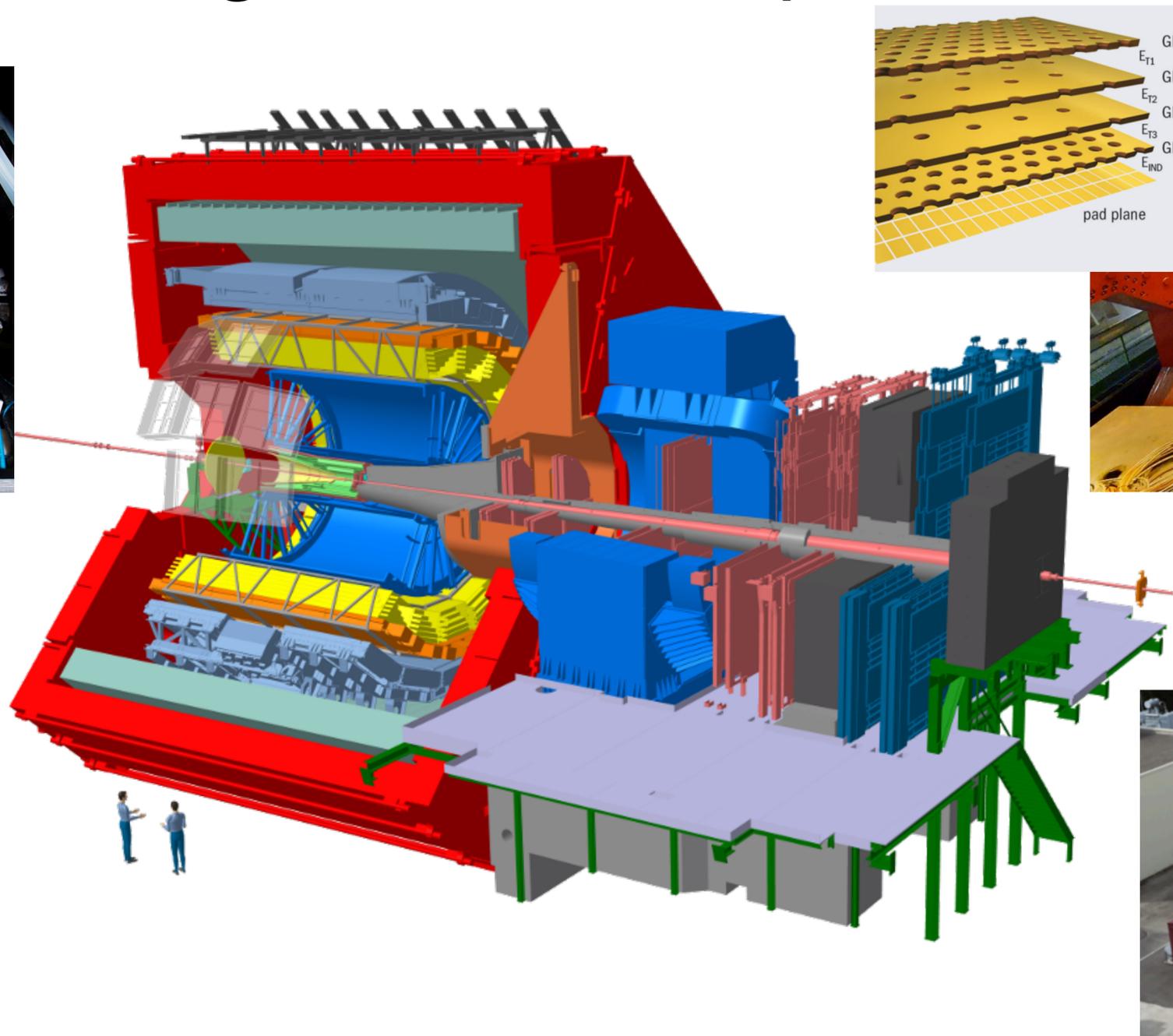
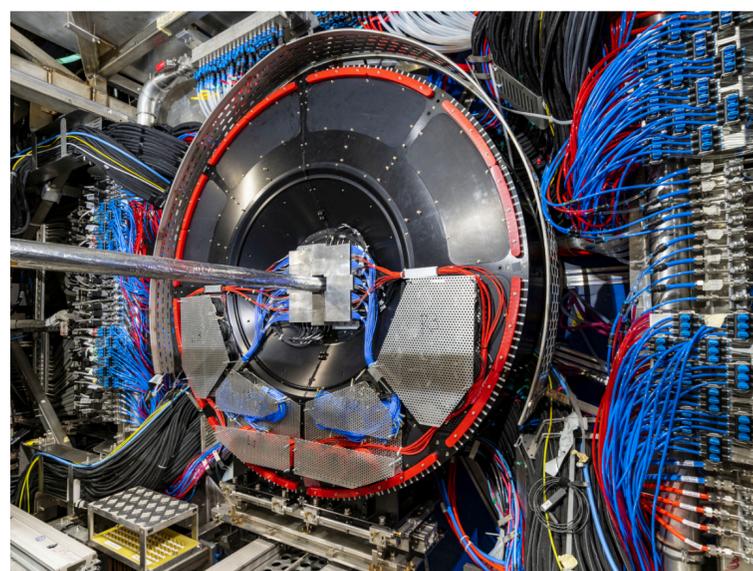


ALICE upgrades in Long Shutdown 2 (2019-2021)

New ITS and MFT

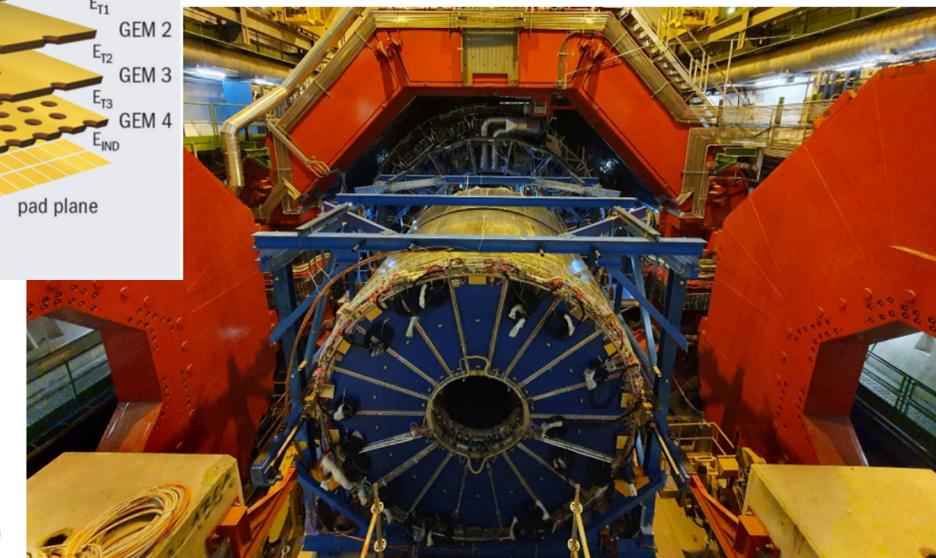
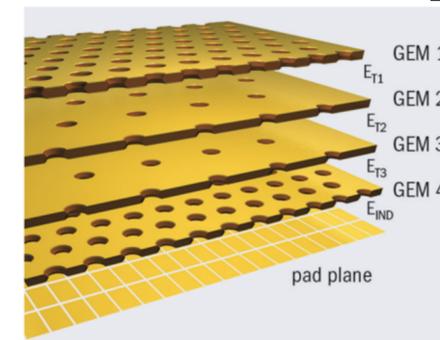


Full pixel detector
Improved spatial resolution
Fast Interaction Trigger



ALICE LS2 upgrade paper: [arXiv:2302.01238](https://arxiv.org/abs/2302.01238)

TPC: GEM readout



Continuous readout

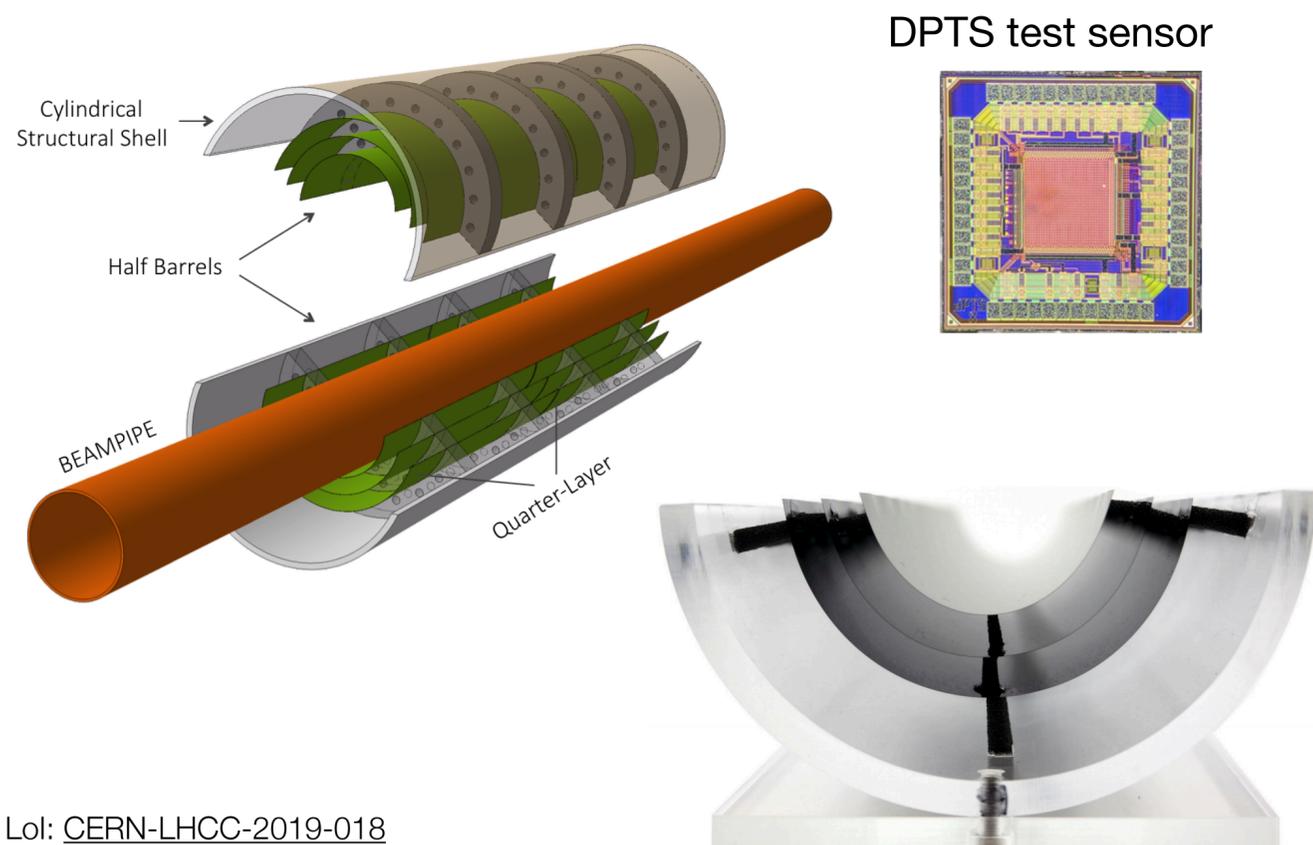
Online event processing



Run 3, 4: collect 13 nb⁻¹ Pb-Pb: 50x more minimum bias data; 10x more triggered data

Future upgrades: ITS 3 and FoCal

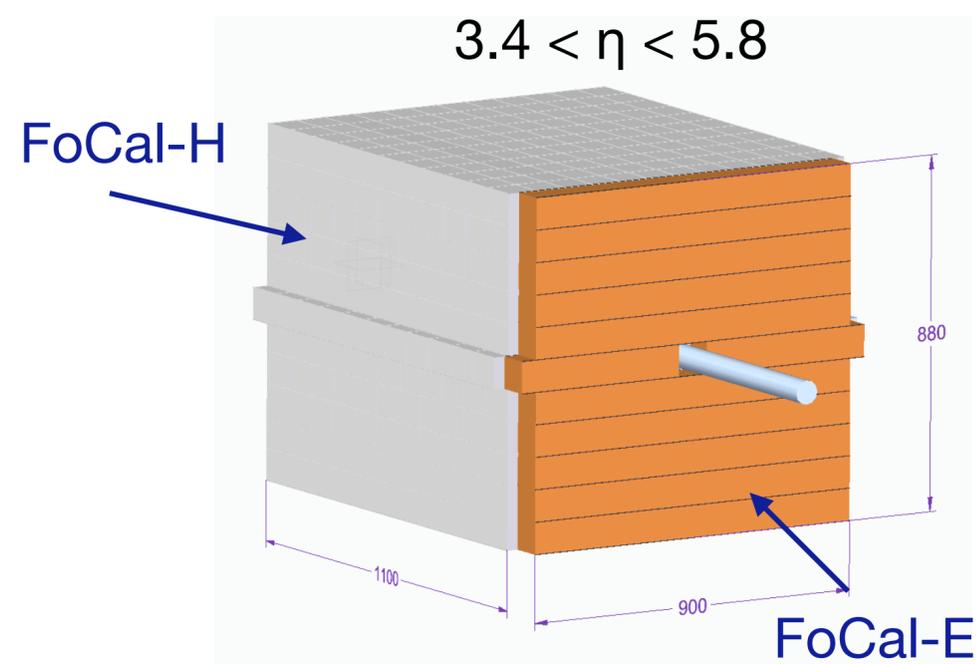
ITS 3: ultra-light, fully cylindrical tracking layers



LoI: CERN-LHCC-2019-018
 DPTS test paper arXiv:2212.08621

- Improved performance for
- Heavy flavour reconstruction
 - **Di-lepton measurements**

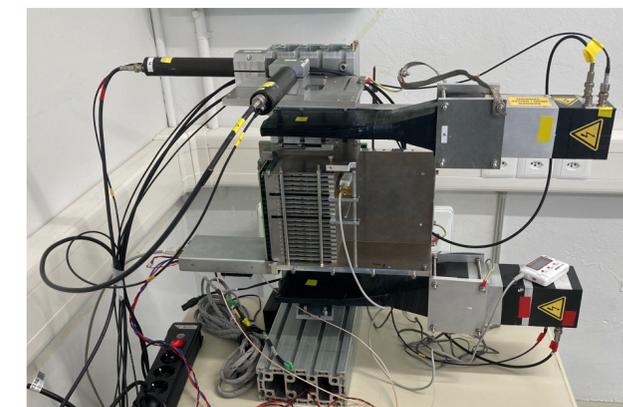
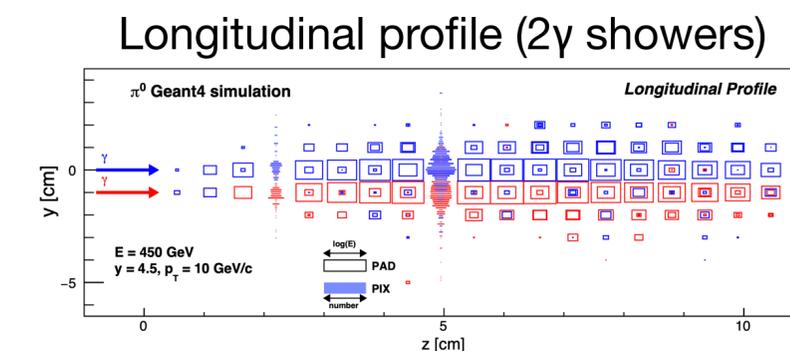
FoCal: high-granularity forward calorimeter



LoI: CERN-LHCC-2020-009

High-granularity Si-W EM calorimeter for **photons and π^0**

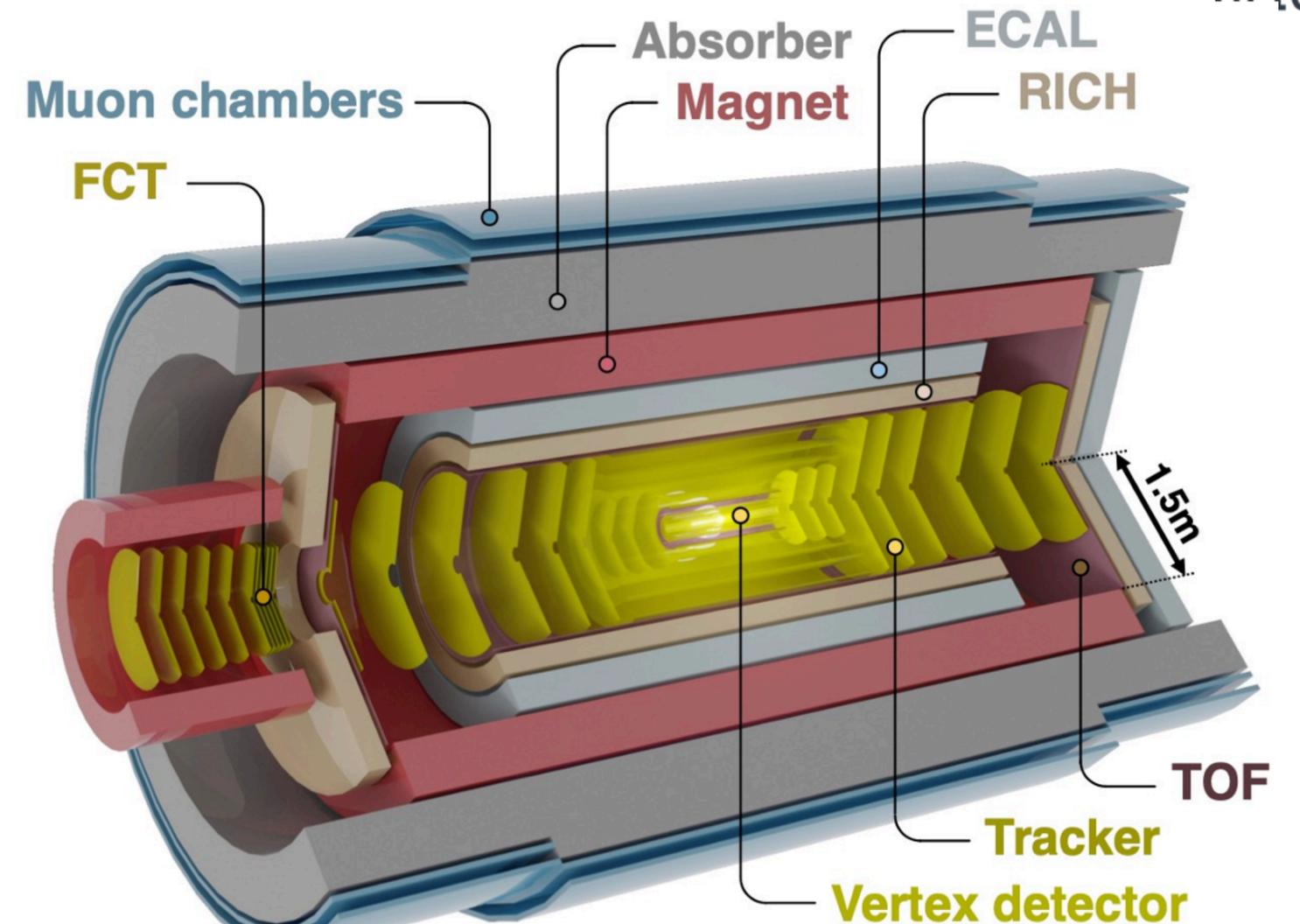
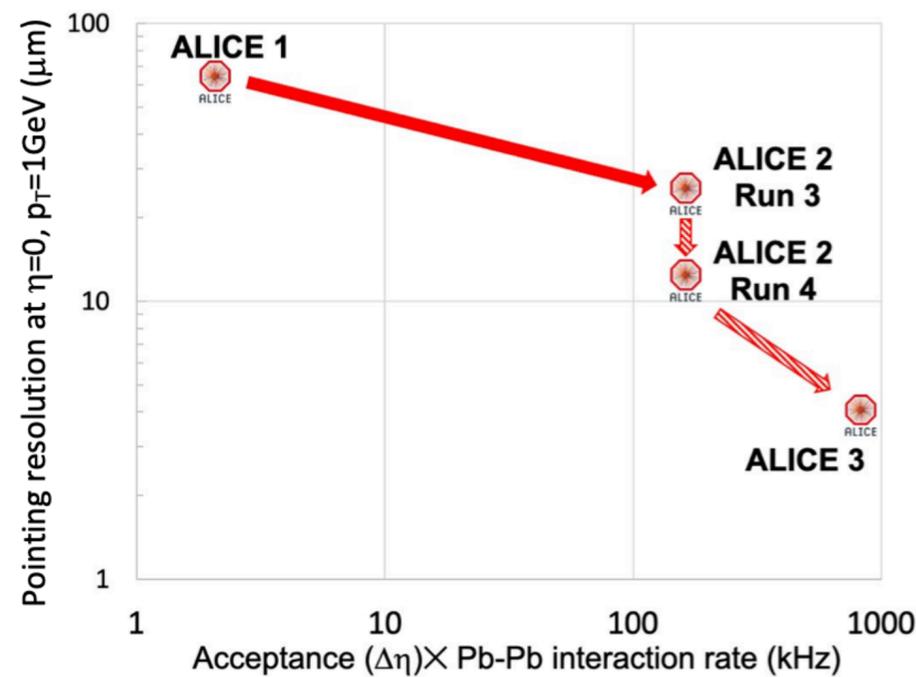
- **Small-x physics** in pp and p-Pb
- Forward π^0 in Pb-Pb



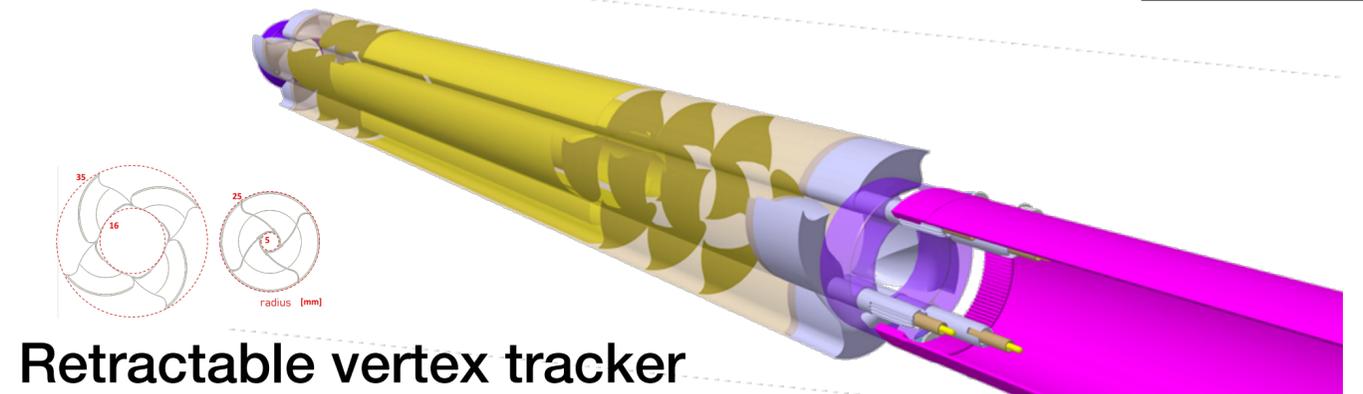
LHC Run 5 and 6: ALICE 3

- Compact all-silicon tracker with high-resolution vertex detector
- Particle Identification over large acceptance: muons, electrons, hadrons, photons
- Fast read-out and online processing

Improvement of detector performance for the upgrades

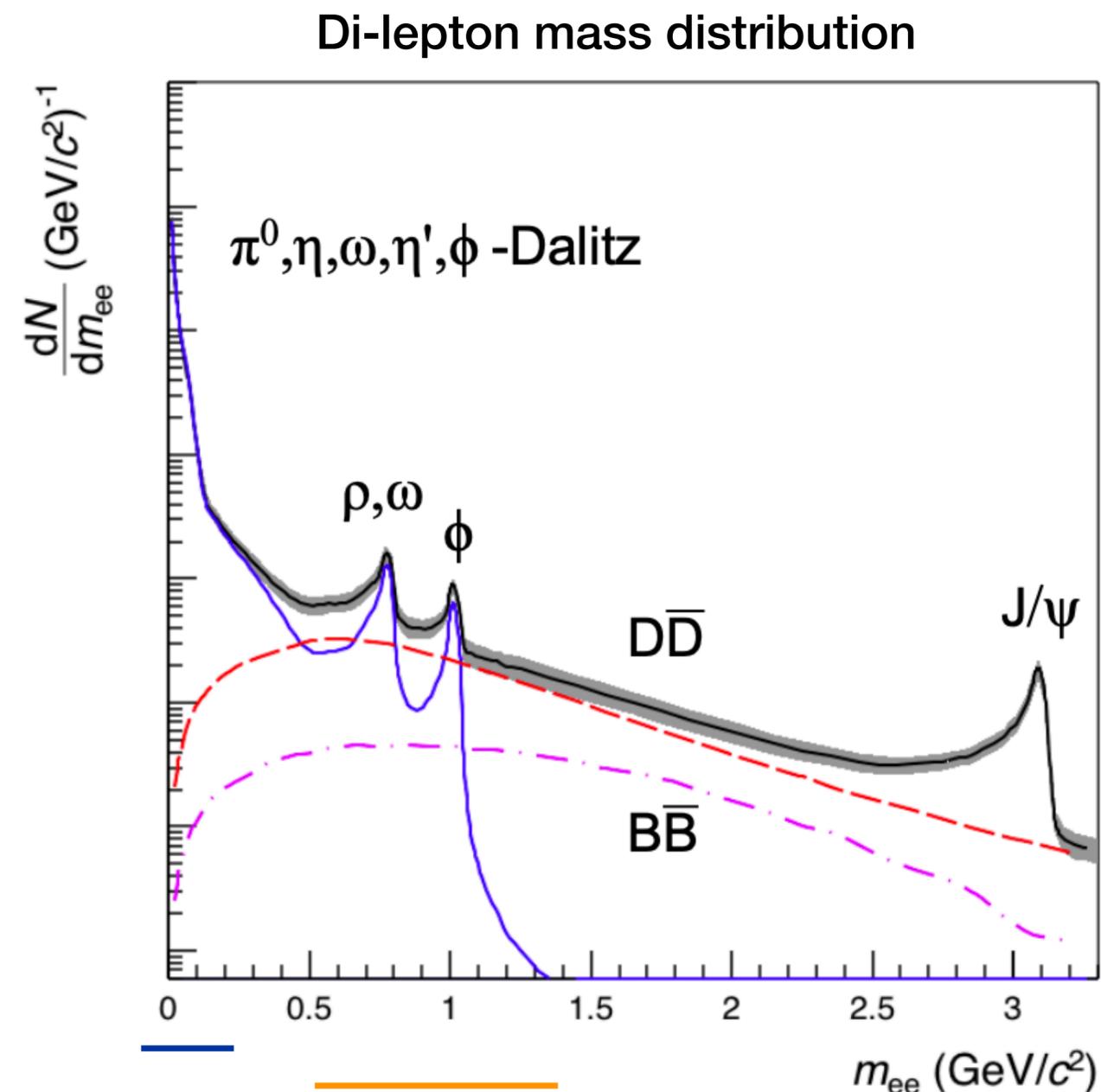


Letter of Intent: [LHCC-2022-009](#)



Di-lepton emission: virtual photons

- Virtual photons: e^+e^- pairs
- Cleaner signal:
 - $m_{ee} > 1 \text{ GeV}/c^2$ removes light flavour decay background
 - Remaining background: heavy flavour pairs
- Slope of mass spectrum not blue-shifted
- Vector meson spectral functions sensitive to chiral symmetry restoration



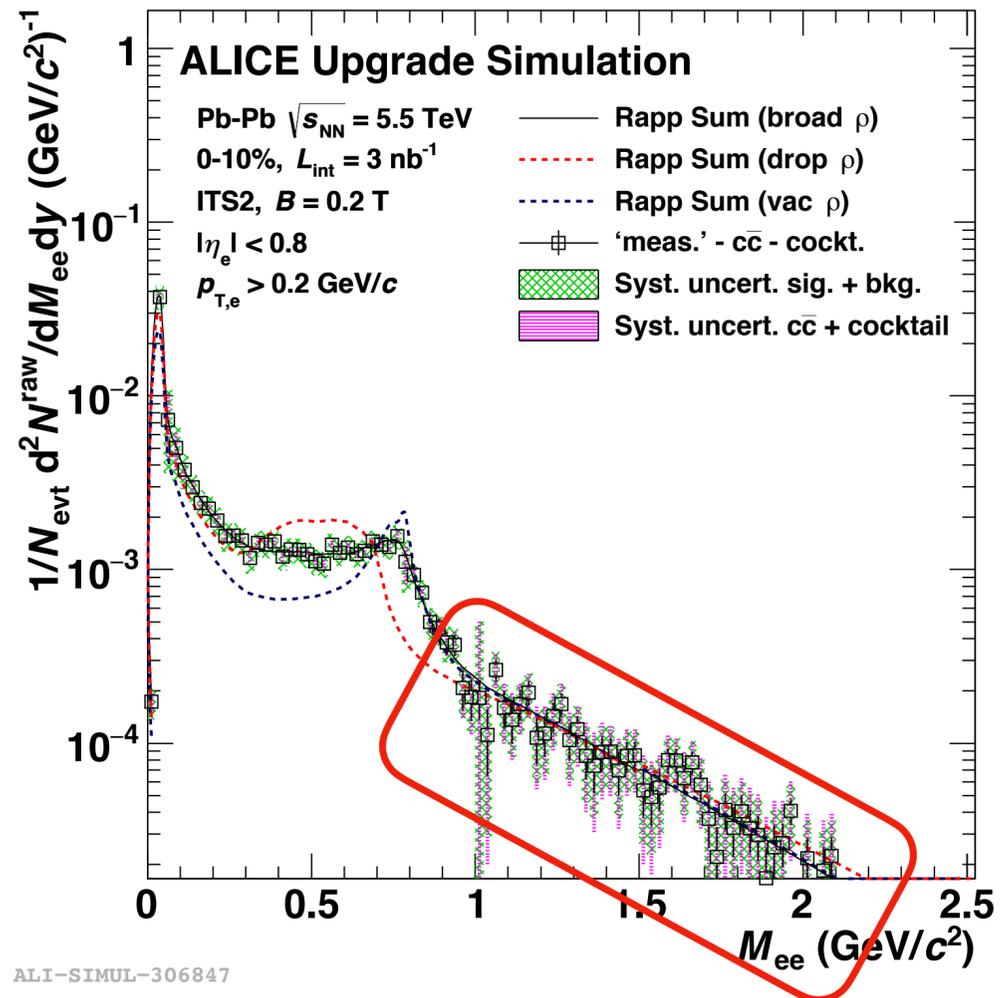
Very low mass:
 π_0 decay background
 conductivity

ω/ϕ region:
 chiral symmetry and
 ρ - a_1 mixing

Large mass:
 thermal emission,
 early times

Dielectrons: chiral symmetry and thermal emission

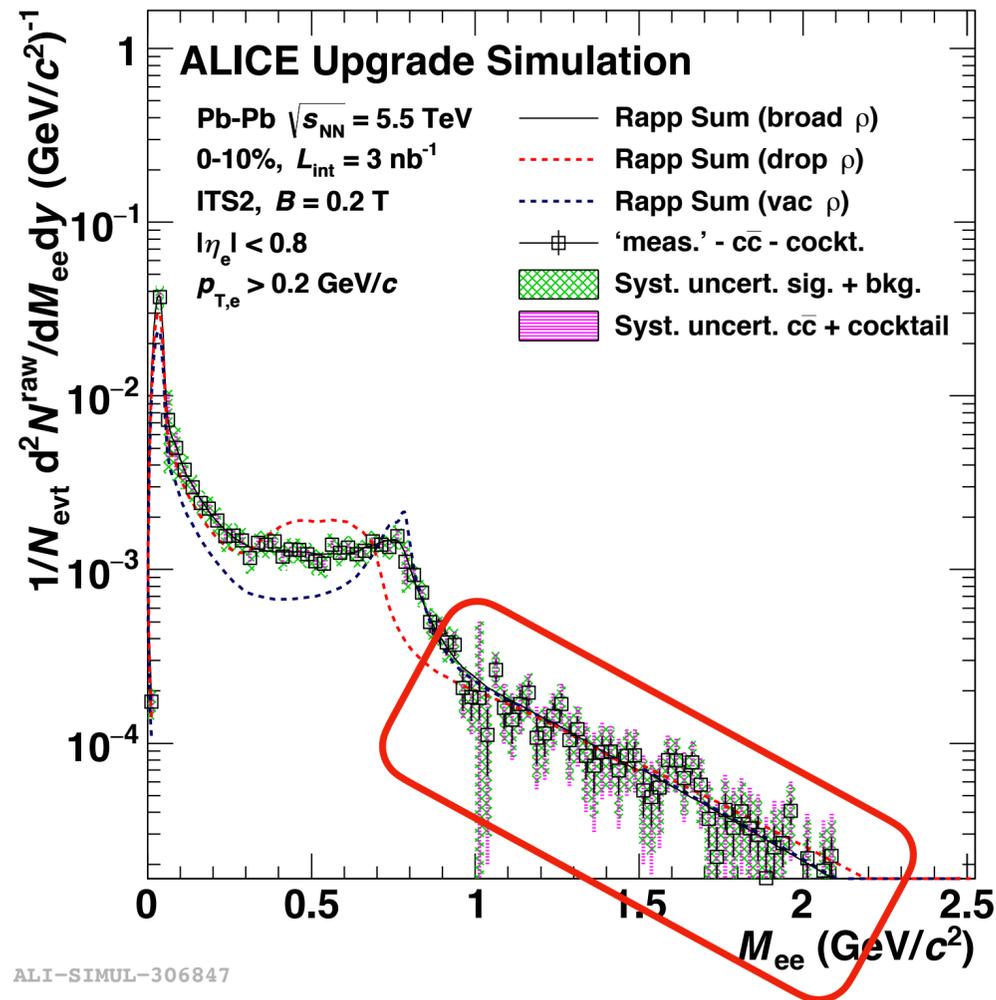
Dileptons in run 3 and 4



Run 3 and 4: first measurements
of thermal dilepton emission at LHC

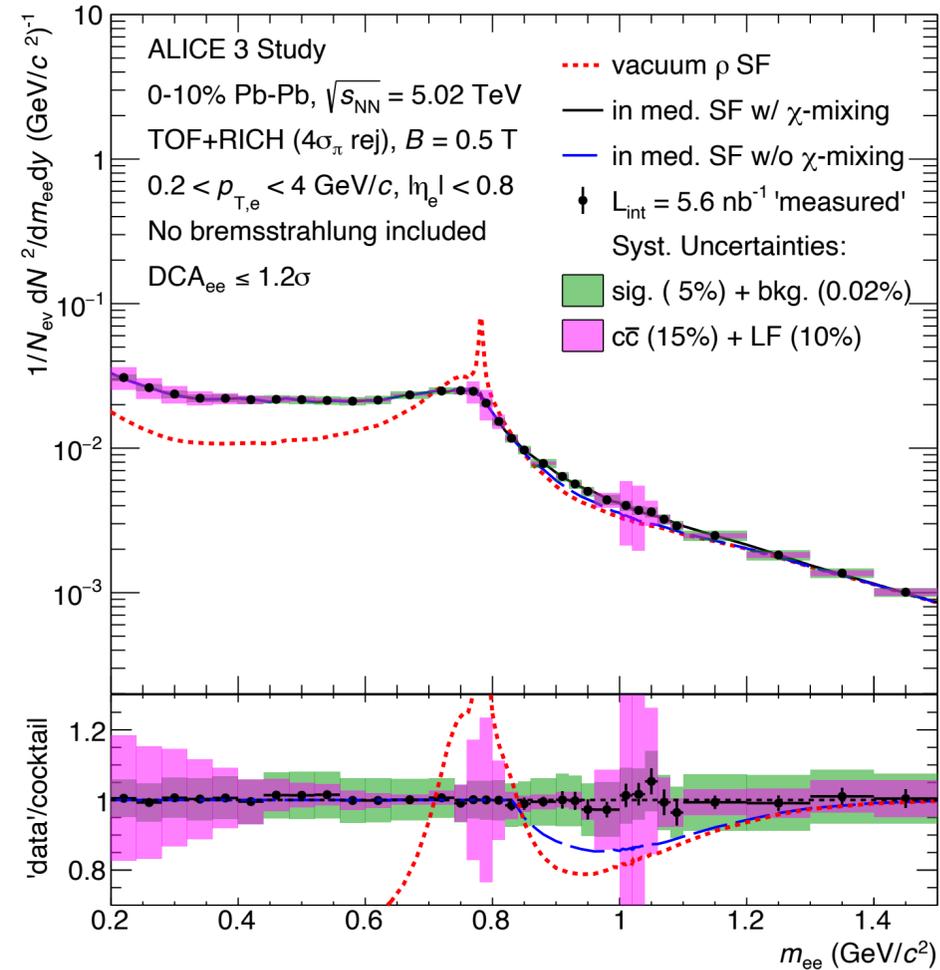
Dielectrons: chiral symmetry and thermal emission

Dileptons in run 3 and 4



Run 3 and 4: first measurements of thermal dilepton emission at LHC

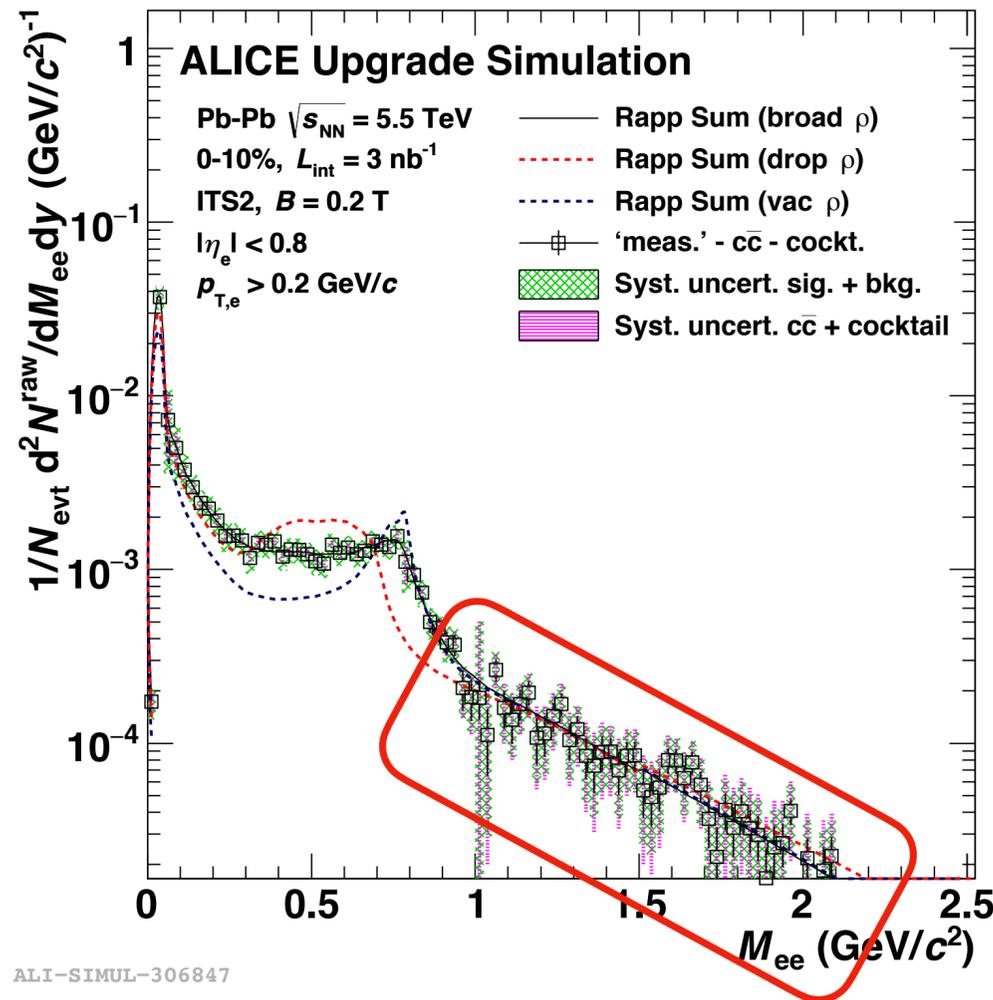
ALICE 3 mass spectrum



High precision:
access $\rho - a_1$ mixing

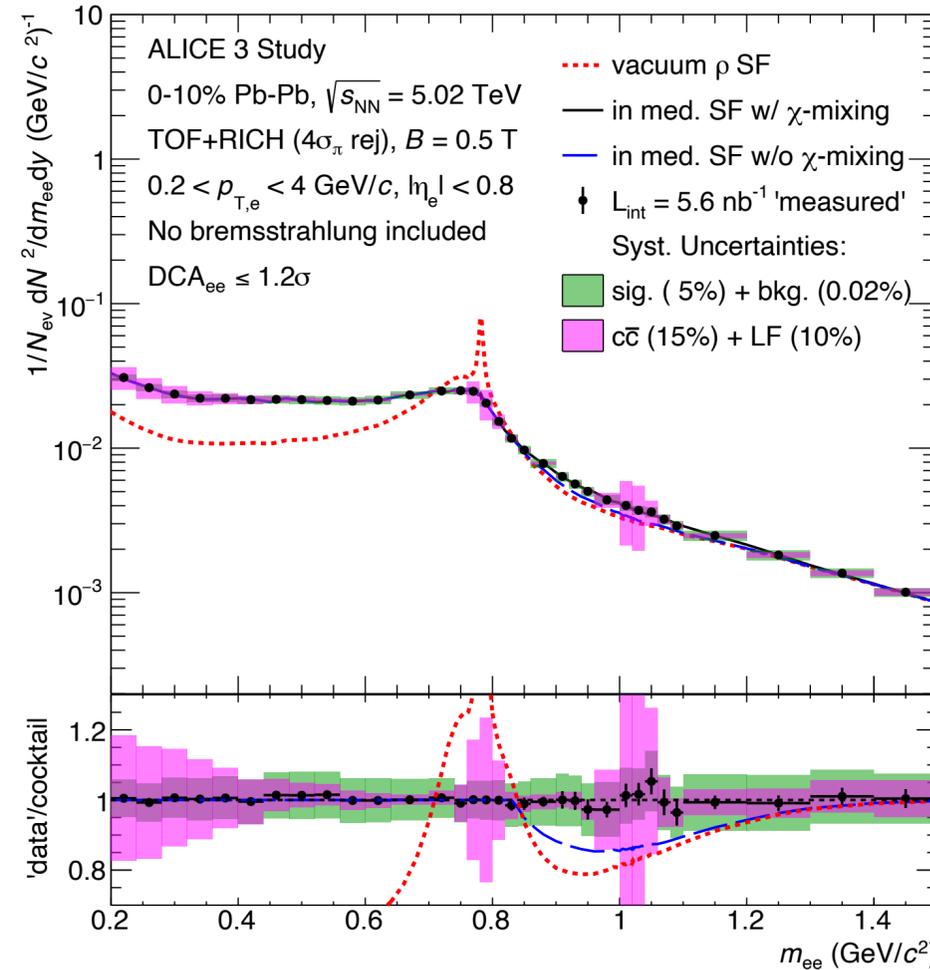
Dielectrons: chiral symmetry and thermal emission

Dileptons in run 3 and 4



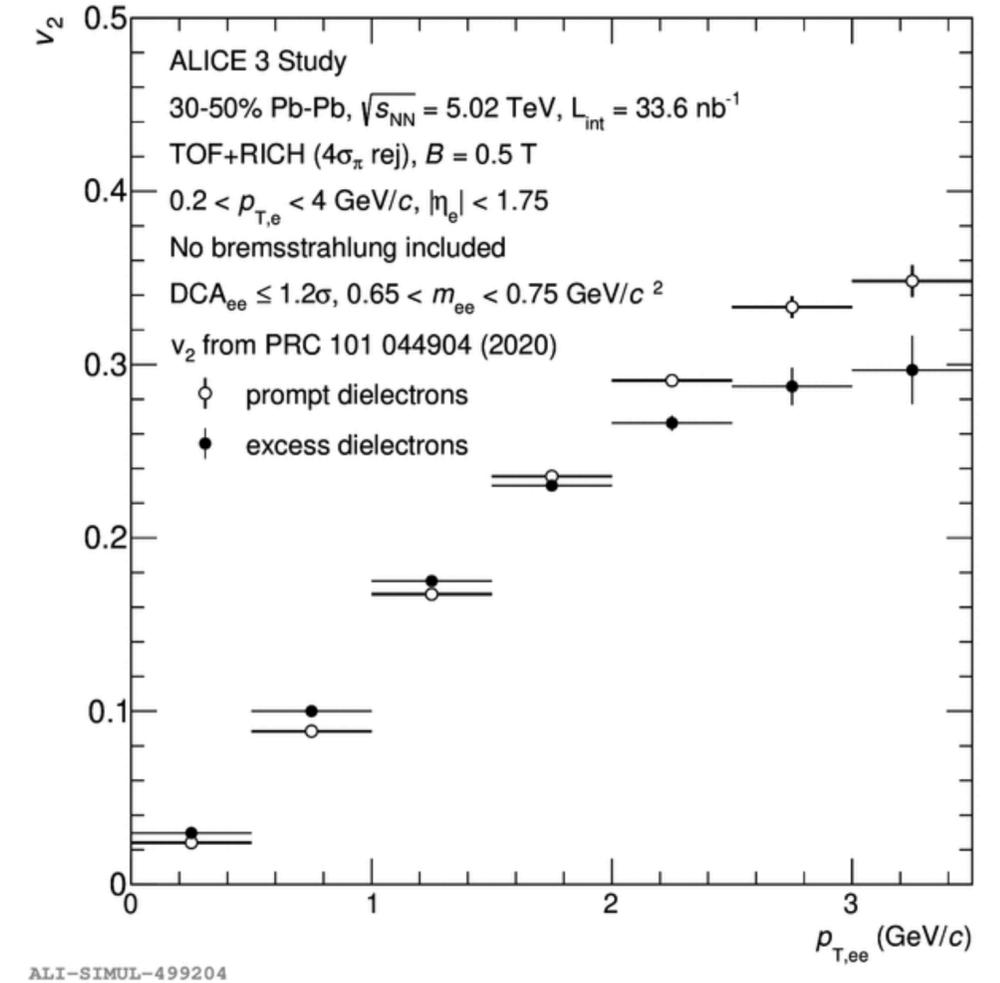
Run 3 and 4: first measurements of thermal dilepton emission at LHC

ALICE 3 mass spectrum



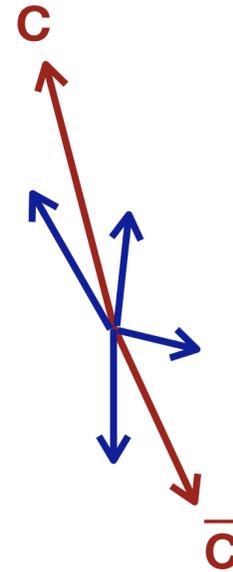
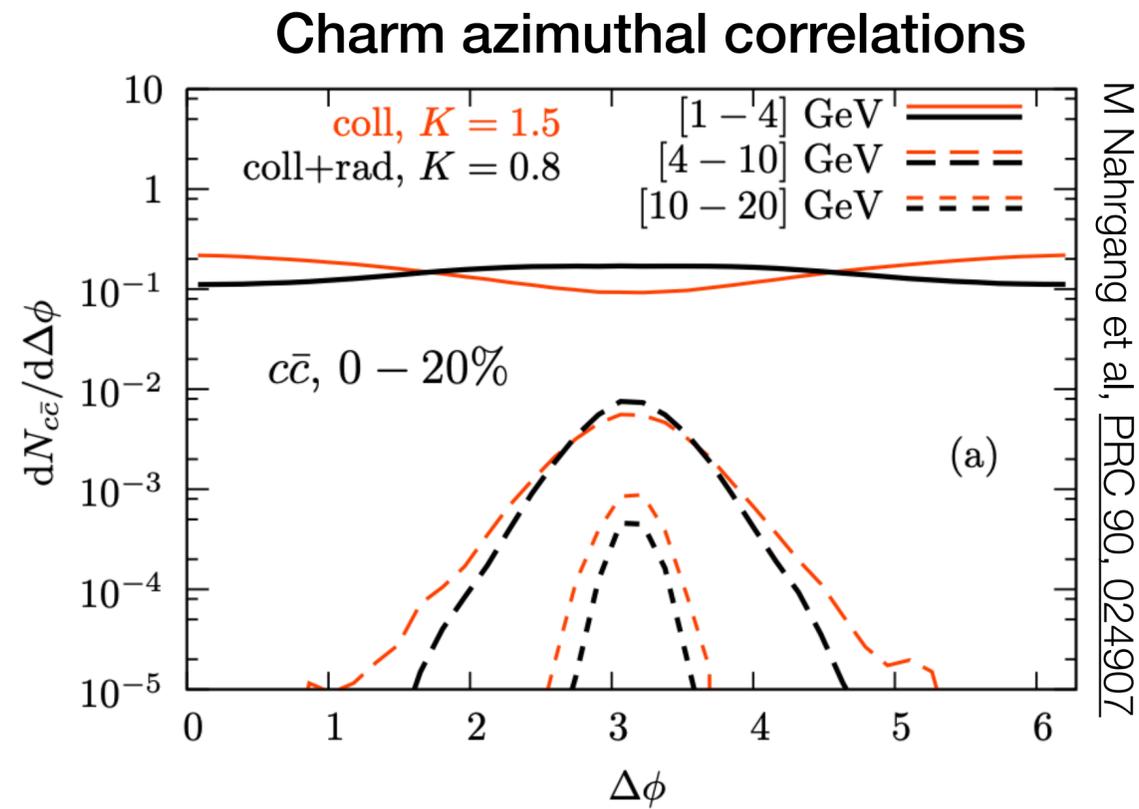
High precision:
 access $\rho - a_1$ mixing

Dielectron v_2



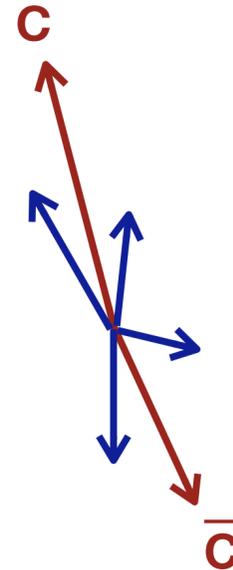
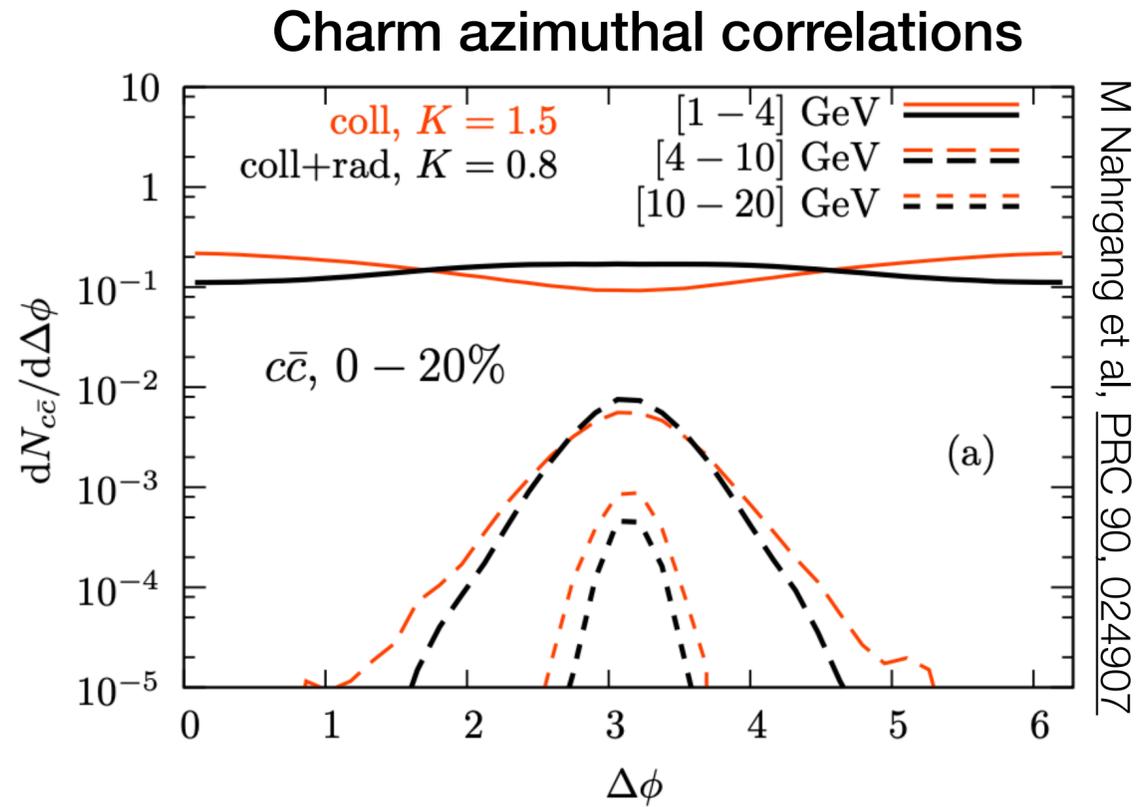
Excellent precision for dilepton v_2 vs p_T in different mass ranges
 → time evolution of emission

$D\bar{D}$ azimuthal correlations

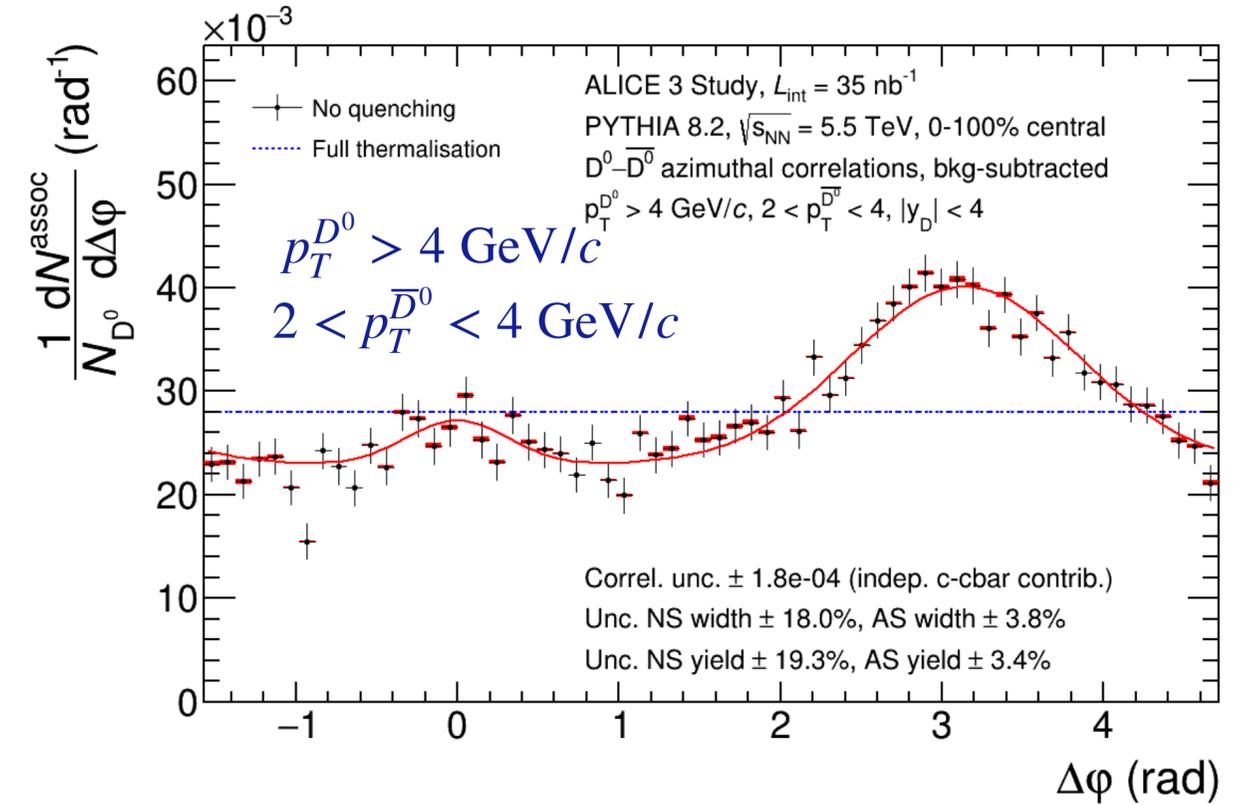


- Angular decorrelation **directly probes QGP scattering**
 - Signal strongest at low p_T
- Very challenging measurement:
 - need good purity, efficiency and η coverage
 - **heavy-ion measurement only possible with ALICE 3**

D \bar{D} azimuthal correlations

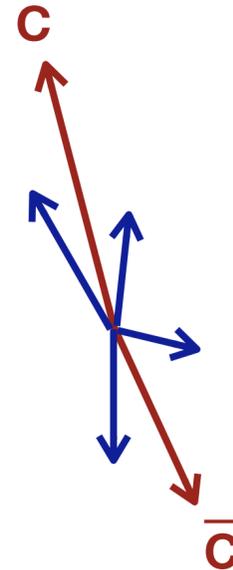
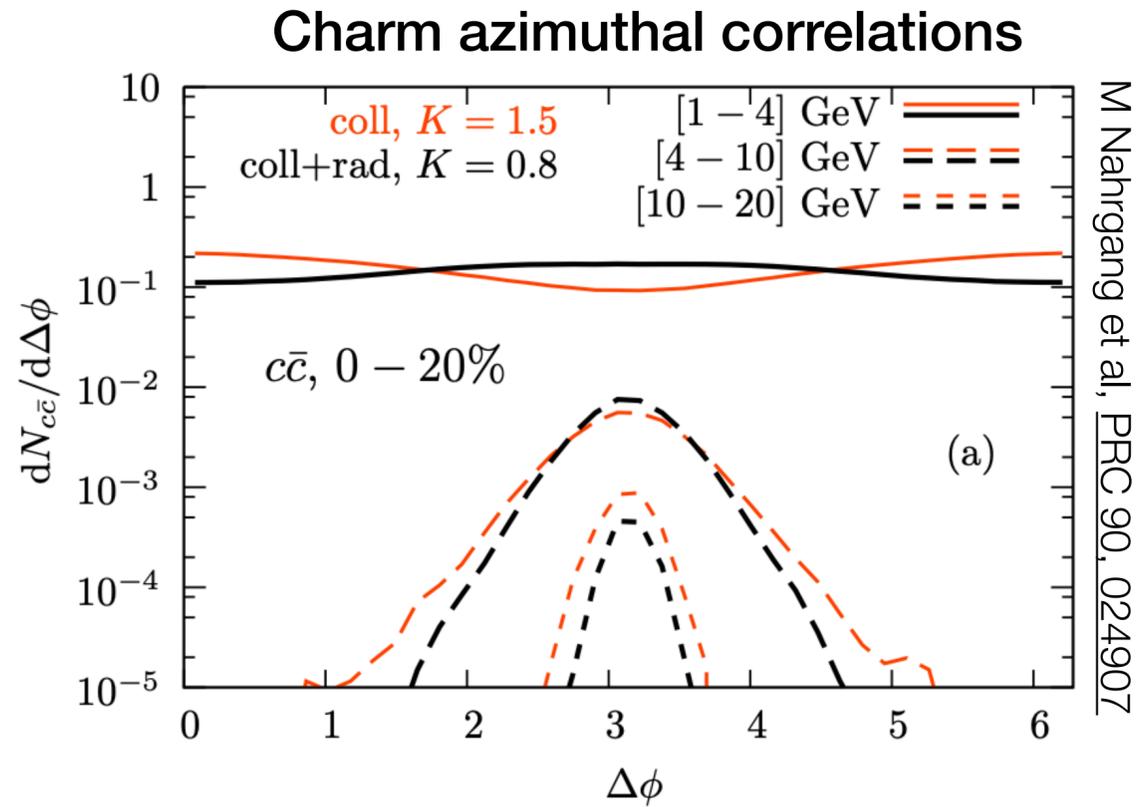


ALICE 3 projection: D \bar{D} correlations

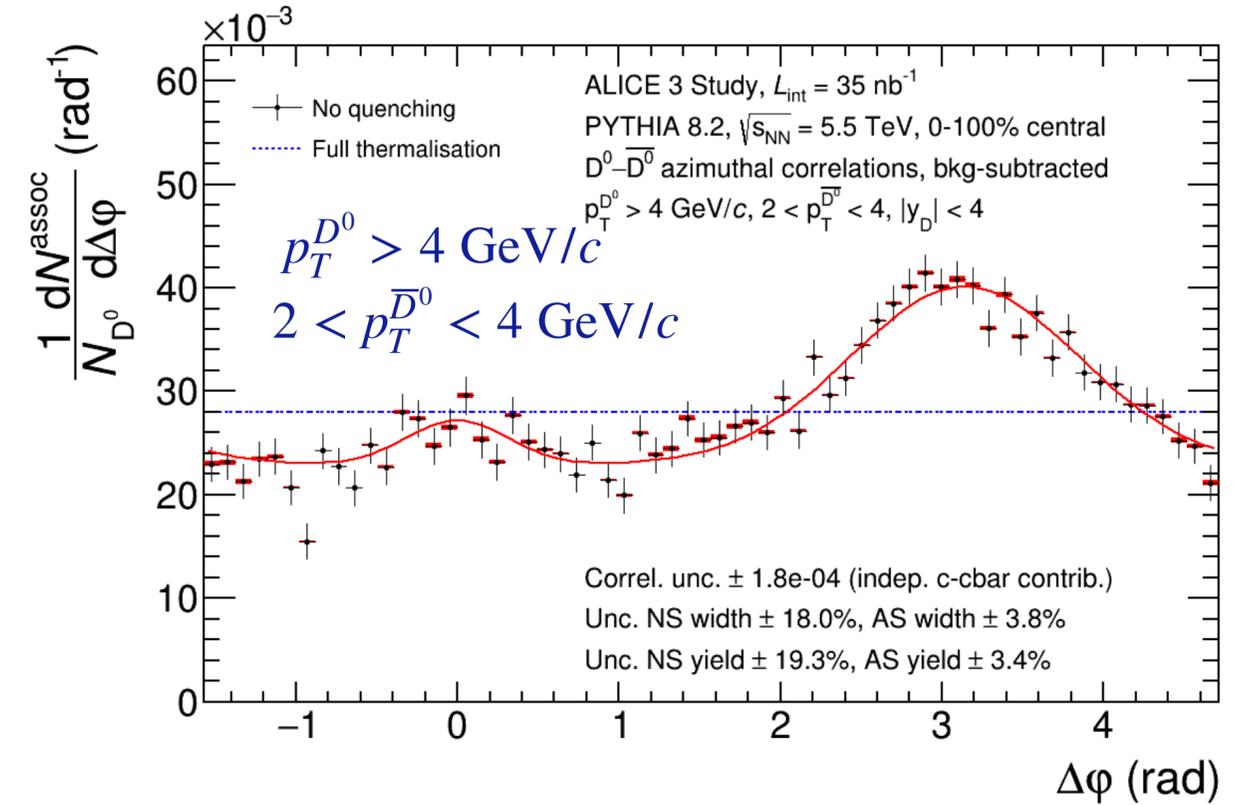


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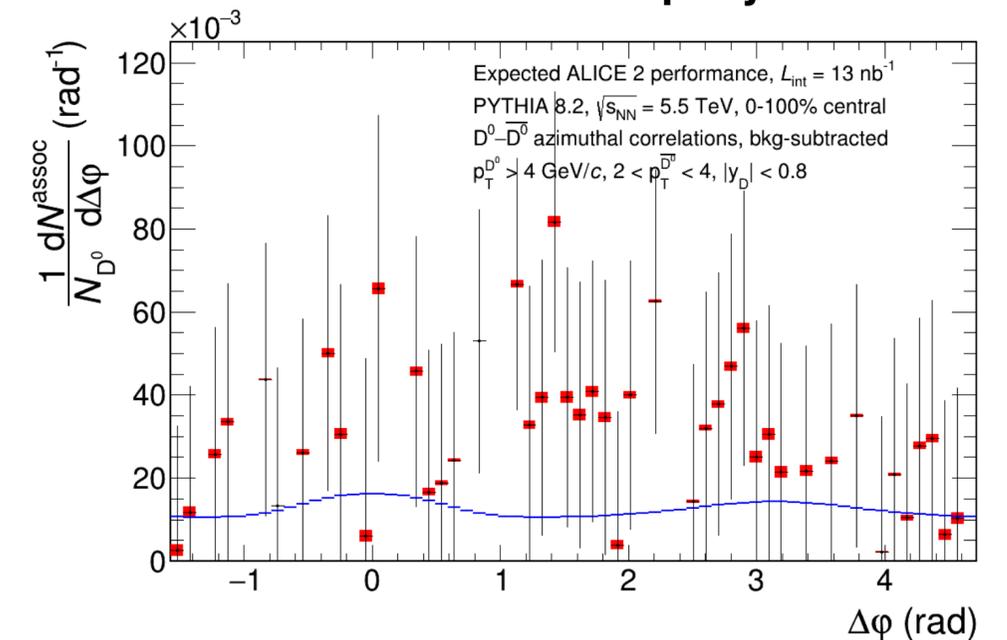
D \bar{D} azimuthal correlations



ALICE 3 projection: D \bar{D} correlations



ALICE Run 3 + 4 projection

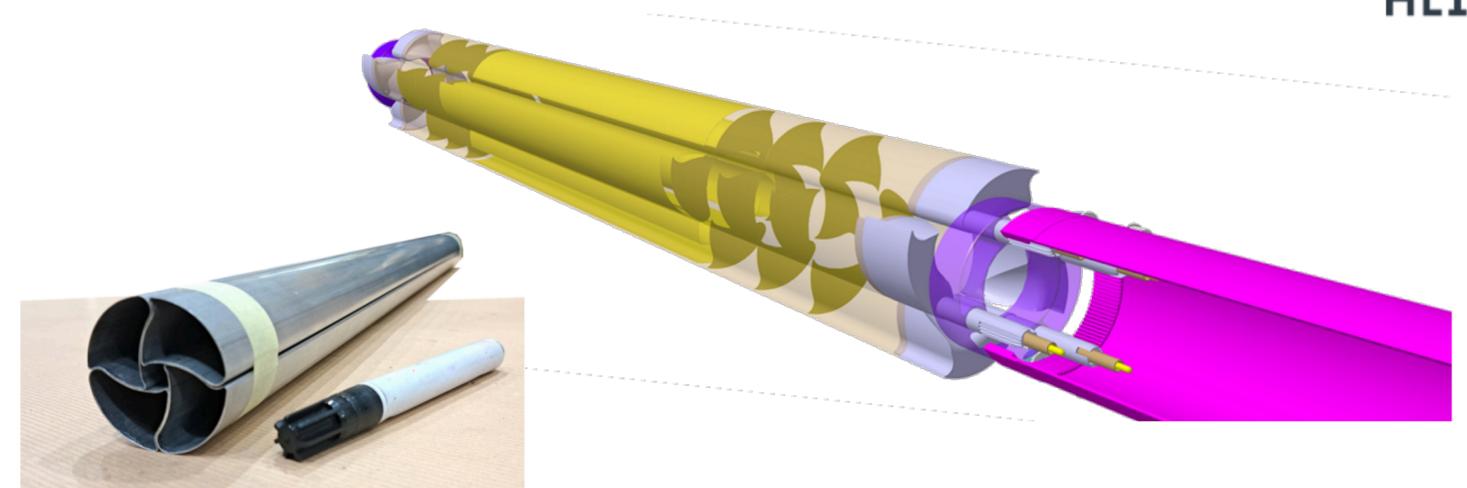


- Angular decorrelation **directly probes QGP scattering**
 - Signal strongest at low p_T
- Very challenging measurement: need good purity, efficiency and η coverage
 → **heavy-ion measurement only possible with ALICE 3**

Strategic detector R&D for ALICE 3

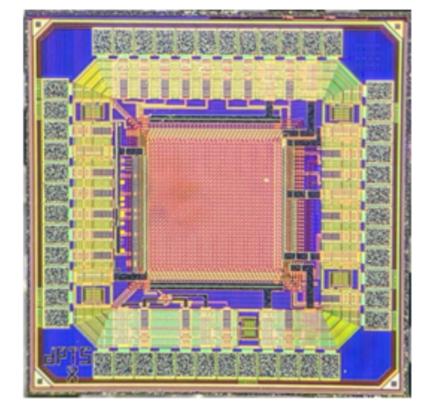
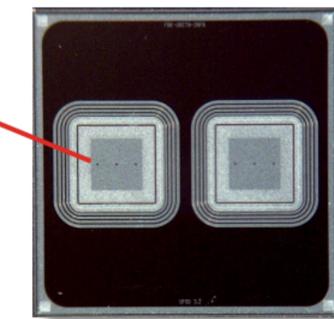
- **Vertex tracker mechanics**
- **Silicon tracking sensor development**
 - Improve read-out speed, radiation hardness
 - Low-power for cooling, large scale application
 - Module integration
- **Integrated timing sensor development**
 - Goal: achieve monolithic timing sensors
 - Improve time resolution to 20 ps
 - Combine photon sensors and timing sensors?
- **Muon detector development**
 - Large area, low noise detectors

R&D for ALICE 3 is starting up
Many exciting opportunities

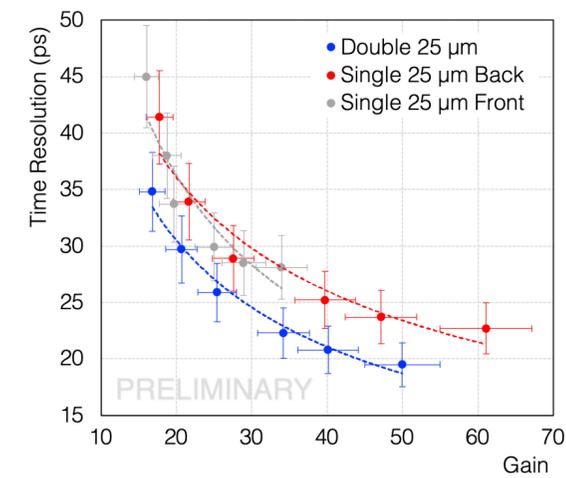


First very thin LGAD prototypes produced by FBK

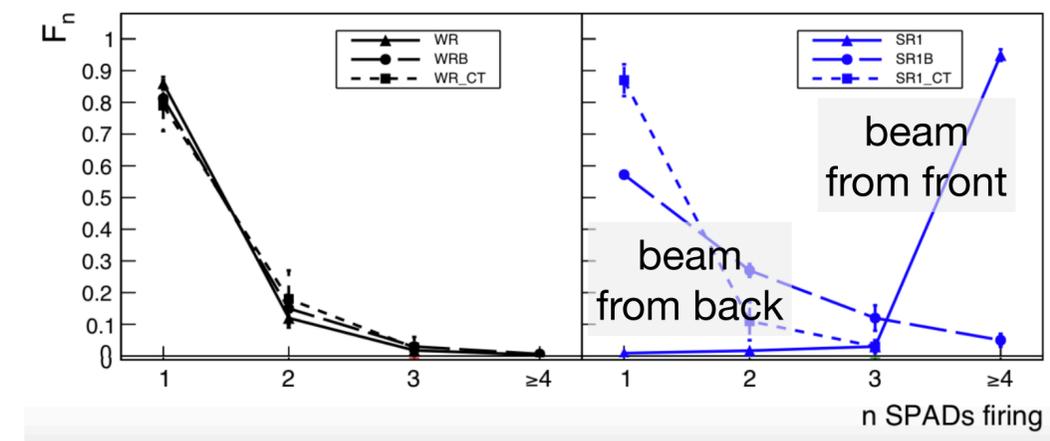
25 μm and **35 μm**-thick FBK single channel
 Area = 1x1 mm²



LGAD time resolution



SiPM hadron response tests



F. Carnesecchi et al, EPJ Plus 138 1, 99

F. Carnesecchi et al, arXiv:2202.04169

Summary

- ALICE studies the *condensed matter* of the strong interaction at high temperature: **the Quark Gluon Plasma**
- Key properties of the plasma are being determined from data:
 - Shear viscosity: close to lower bound
 - First constraints on bulk viscosity
- Charm quarks move with the fluid expansion: **rapid thermalisation**
 - Beauty less so? \Rightarrow Run 3 and 4 to improve precision
- **Indicate short mean free path strong (residual) interactions**
- Future directions
 - Measurement of **initial temperature in reach for upcoming runs**
 - Future upgrades to understand interactions, thermalisation

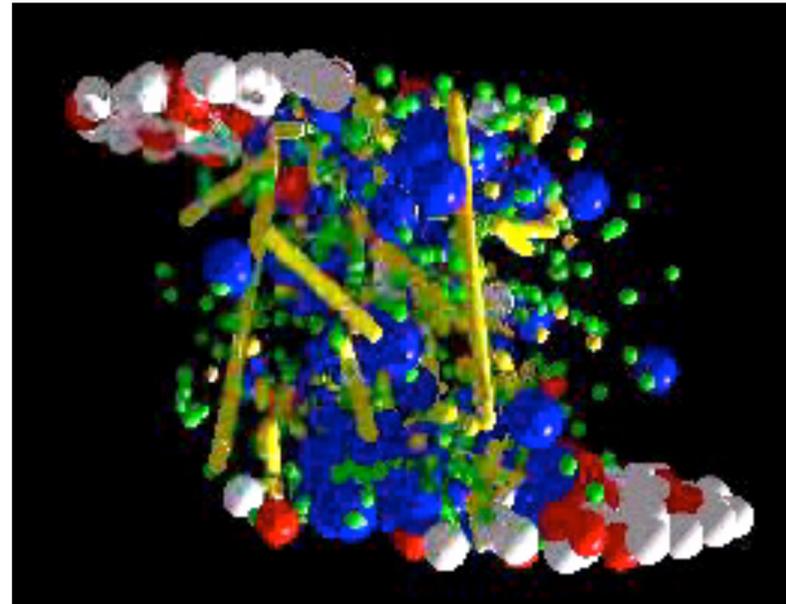
Thank you for your attention!

Probing the Quark-Gluon Plasma

Probe beam



particles



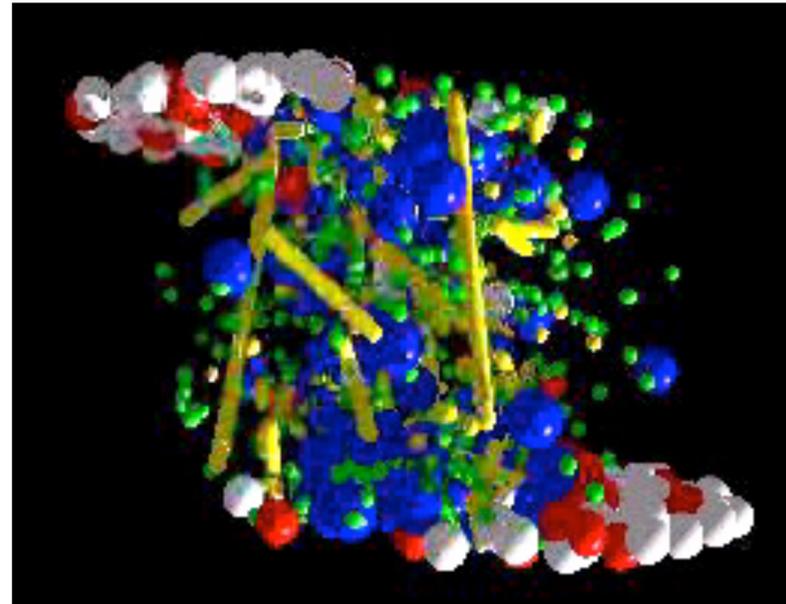
Detector

Probing the Quark-Gluon Plasma

Probe beam



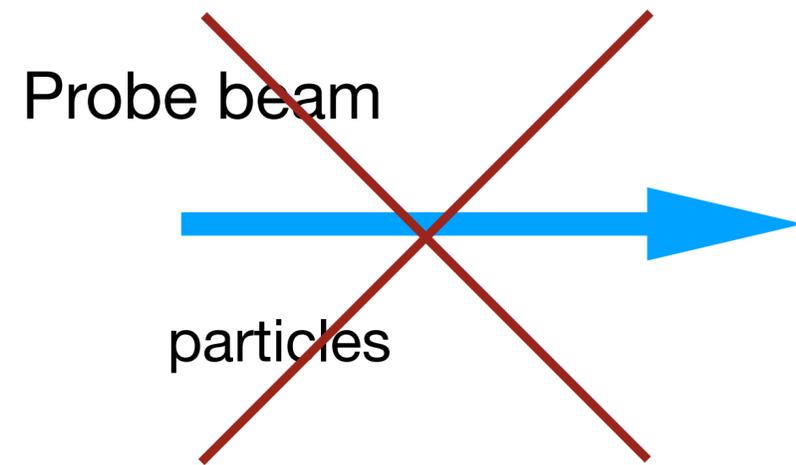
particles



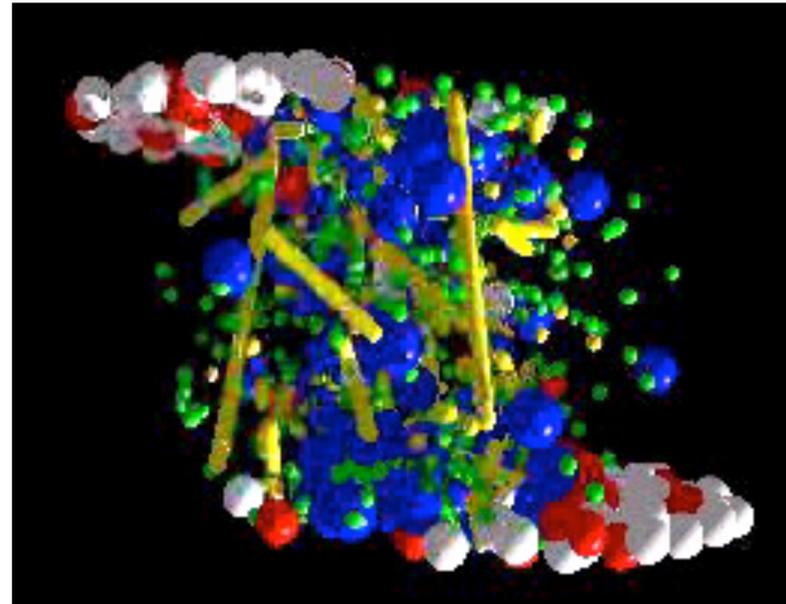
Not feasible:
Short life time
Small size (~ 10 fm)

Detector

Probing the Quark-Gluon Plasma

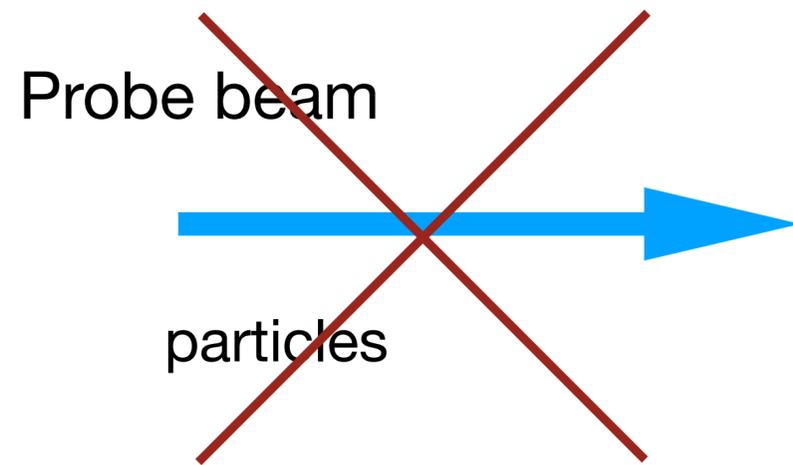


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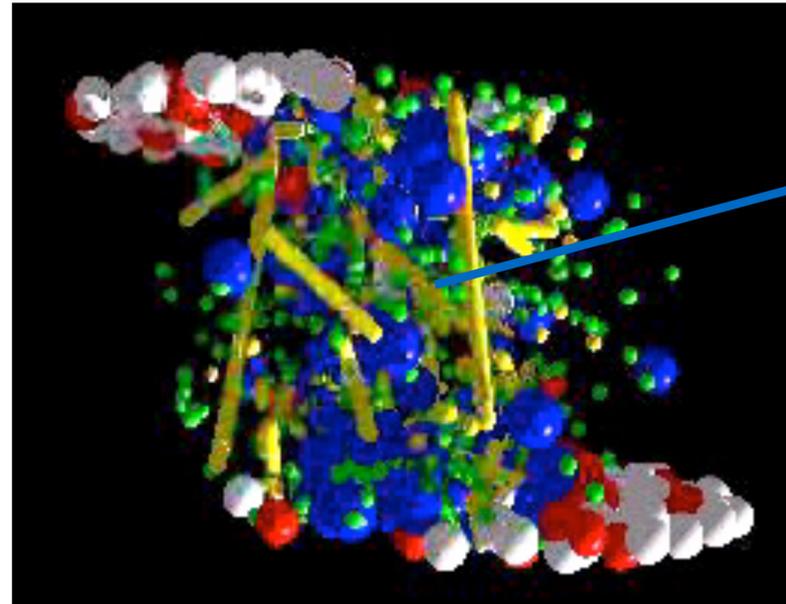


Detector

Probing the Quark-Gluon Plasma



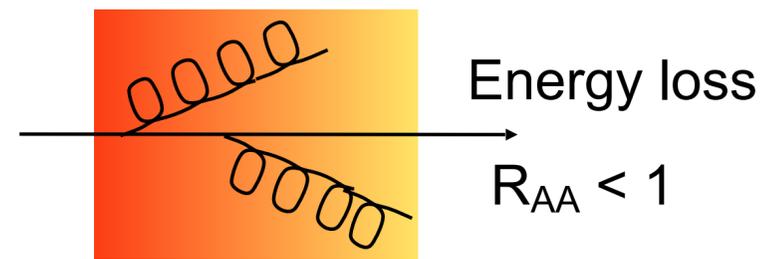
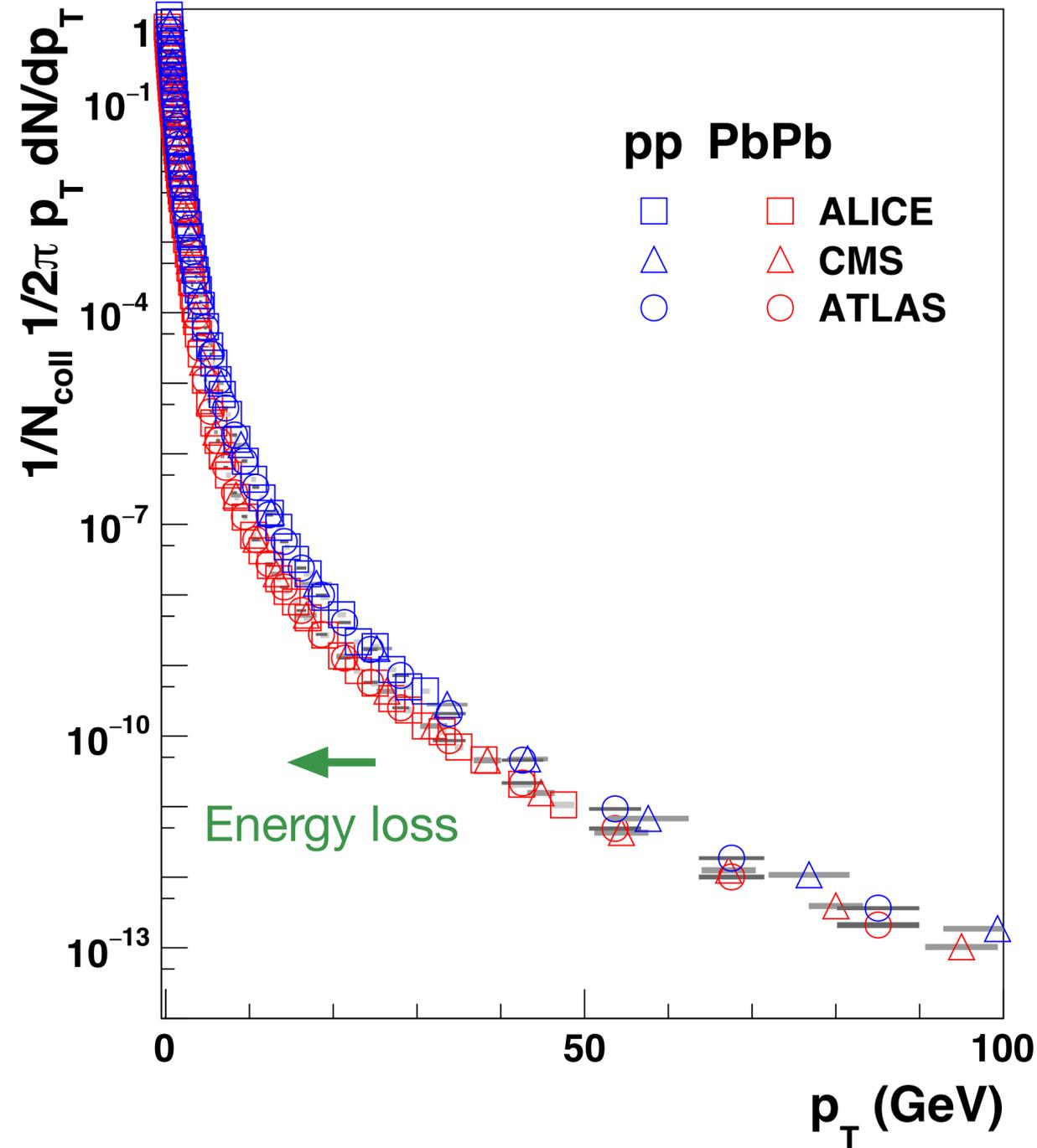
Not feasible:
Short life time
Small size (~ 10 fm)



Use self-generated probe:
quarks, gluons from hard scattering
large transverse momentum

Nuclear modification of p_T spectra

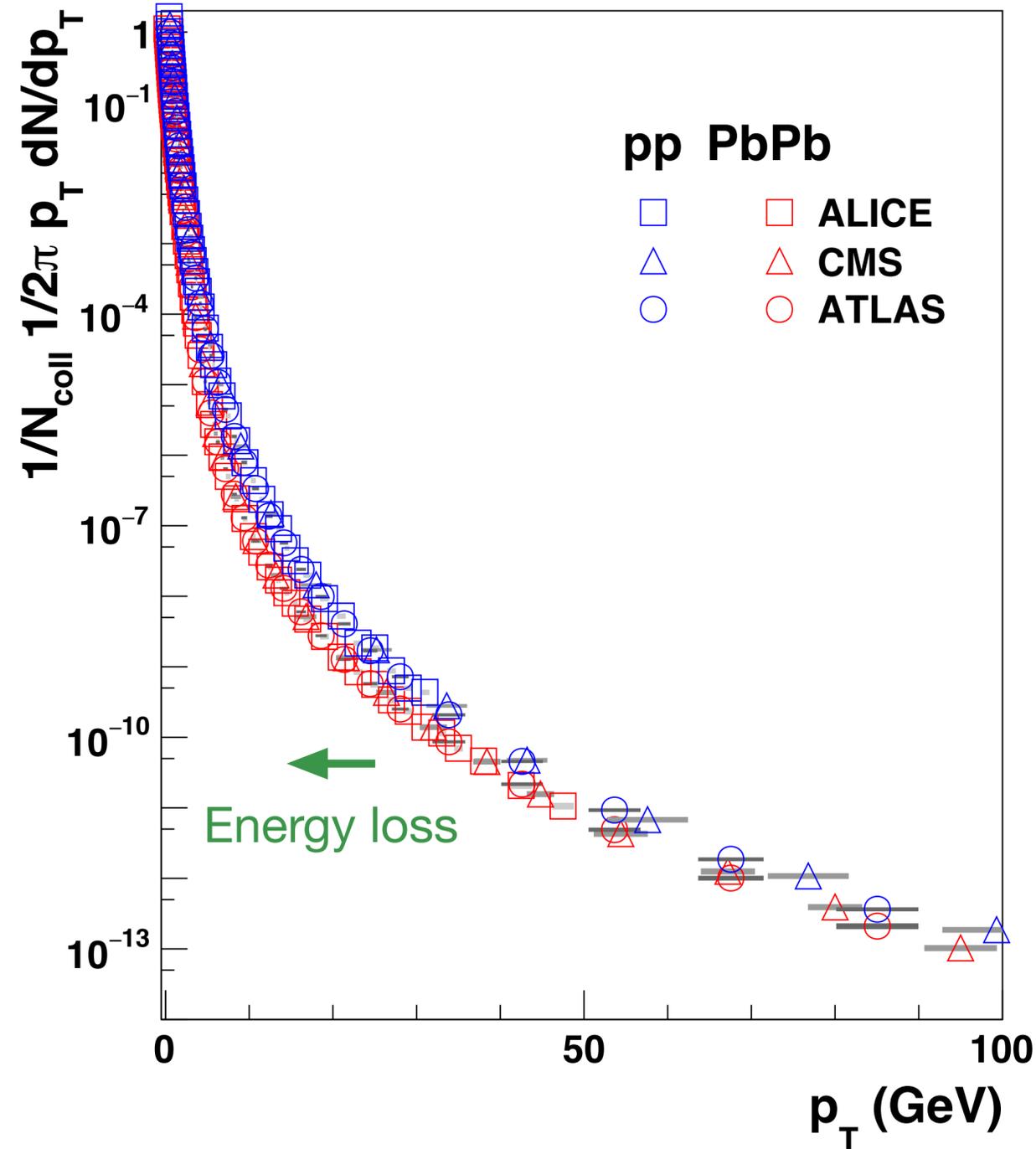
Charged particle p_T spectra



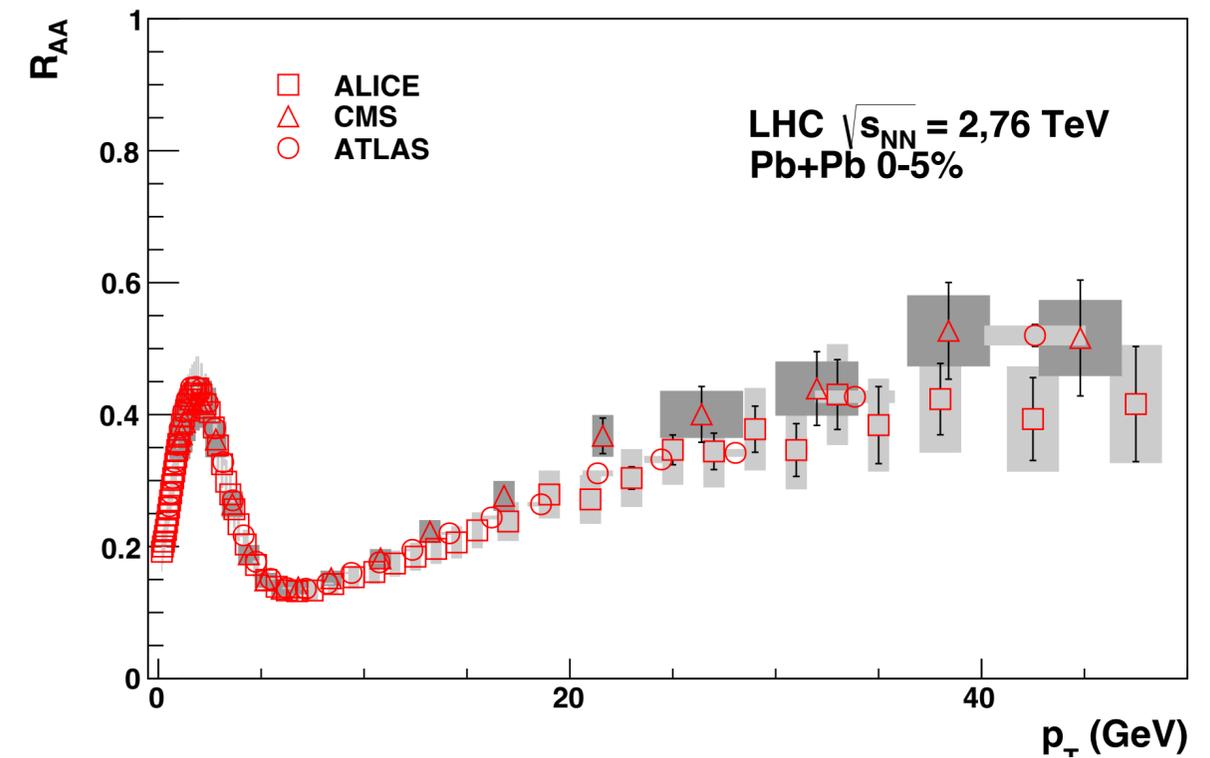
Nuclear modification of p_T spectra

ALICE, PLB720, 52
 CMS, EPJC, 72, 1945
 ATLAS, arXiv:1504.04337

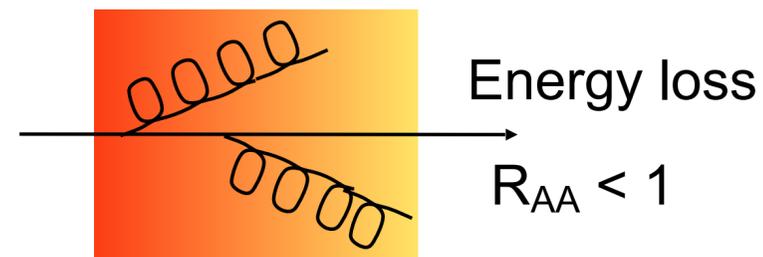
Charged particle p_T spectra



Nuclear modification factor



$$R_{AA} = \frac{dN/dp_T|_{A+A}}{N_{coll} dN/dp_T|_{p+p}}$$

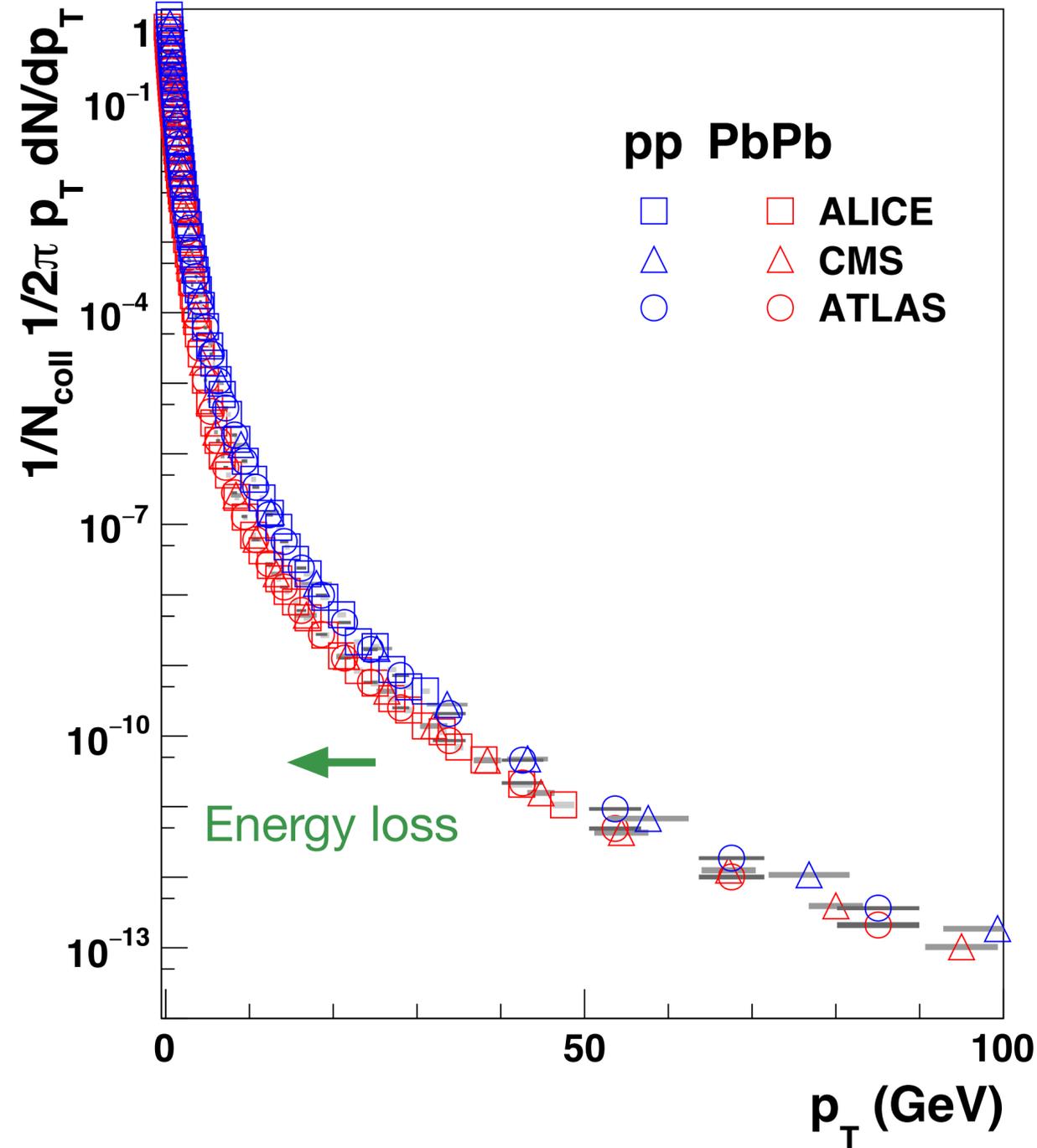


Pb+Pb: clear suppression ($R_{AA} < 1$): parton energy loss

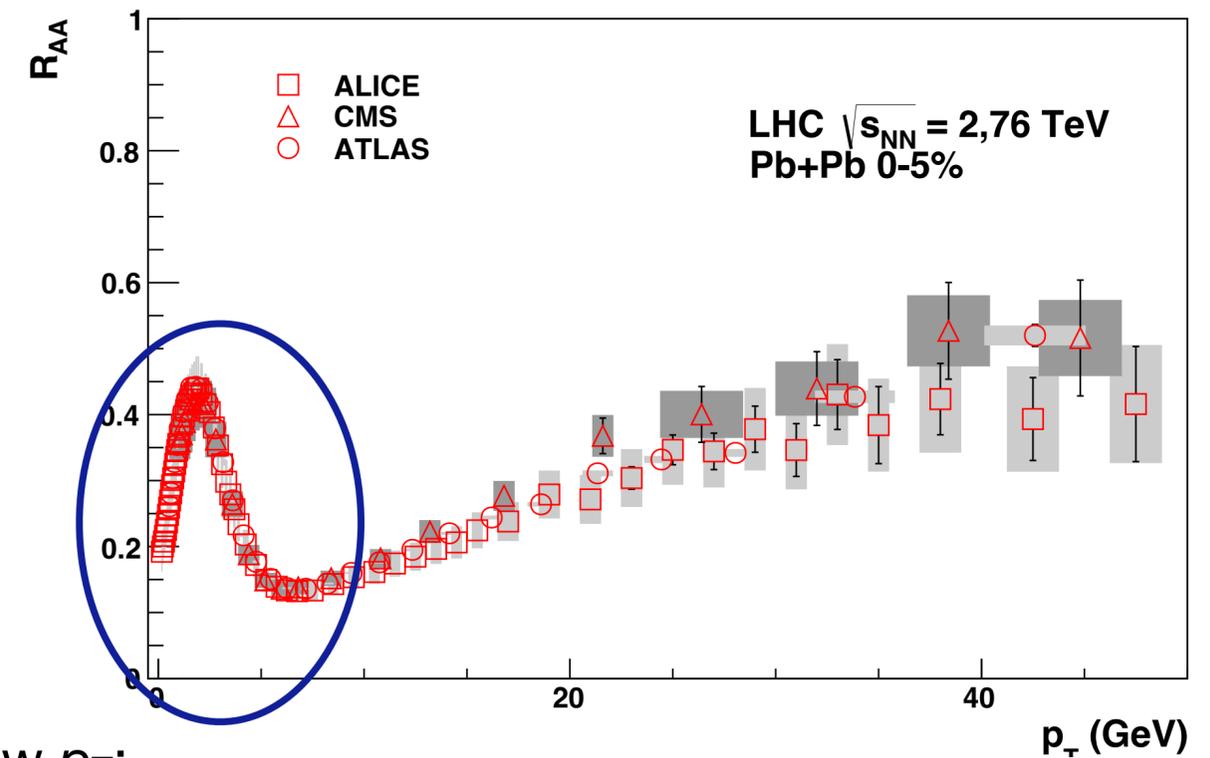
Nuclear modification of p_T spectra

ALICE, PLB720, 52
 CMS, EPJC, 72, 1945
 ATLAS, arXiv:1504.04337

Charged particle p_T spectra

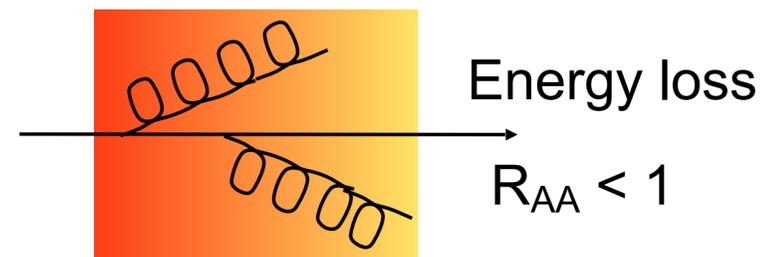


Nuclear modification factor



Low p_T :
 soft production,
 N_{part} scaling

$$R_{AA} = \frac{dN/dp_T|_{A+A}}{N_{coll} dN/dp_T|_{p+p}}$$

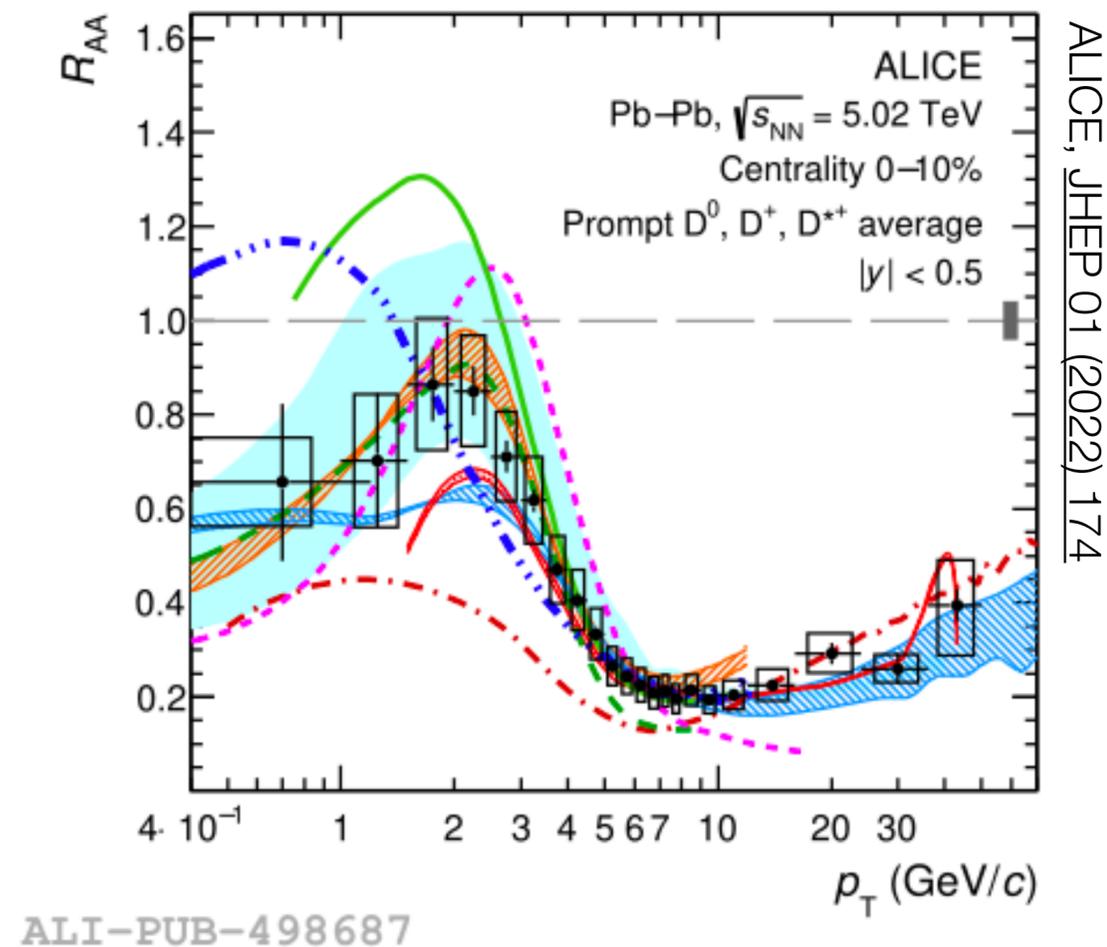


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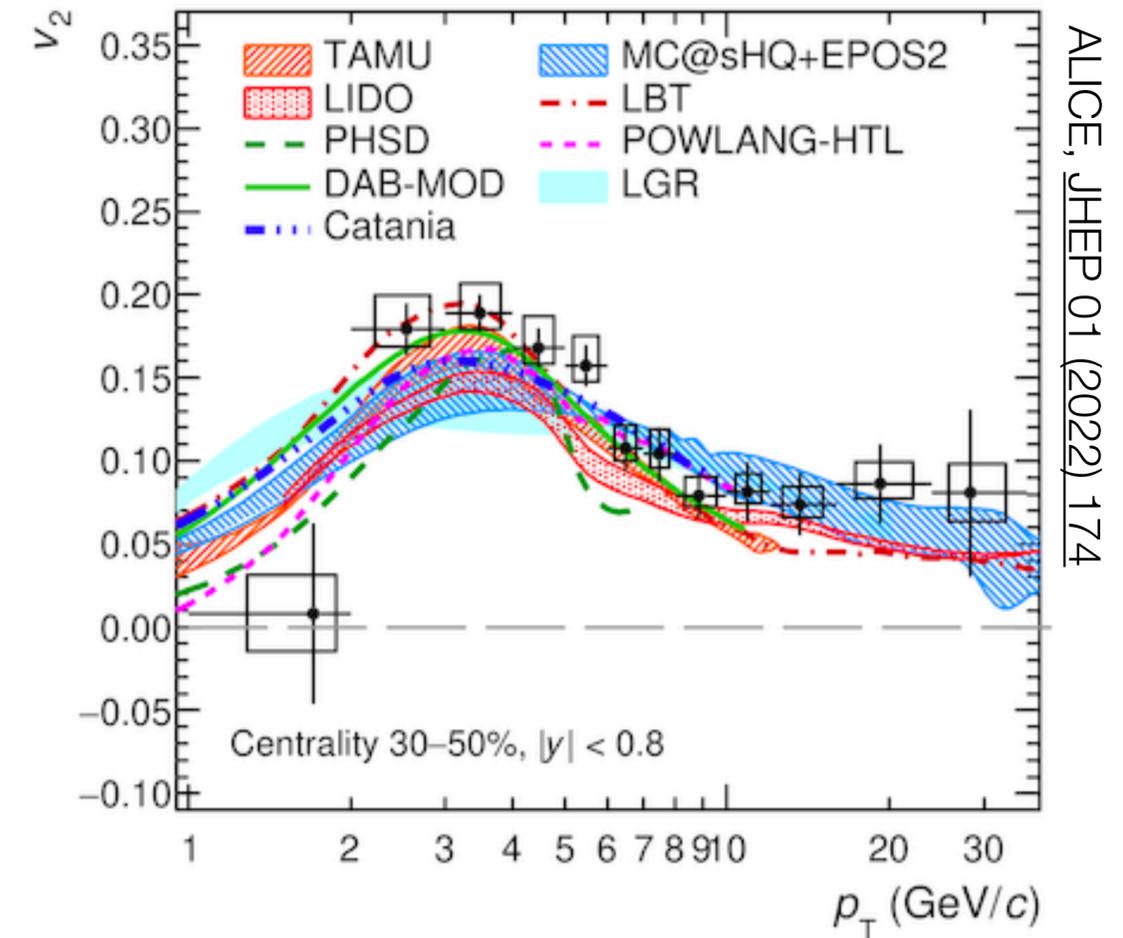
Nuclear modification and elliptic flow of D mesons

D mesons contain a charm quark $m \gg T$, that is produced in an initial hard scattering

Nuclear modification factor



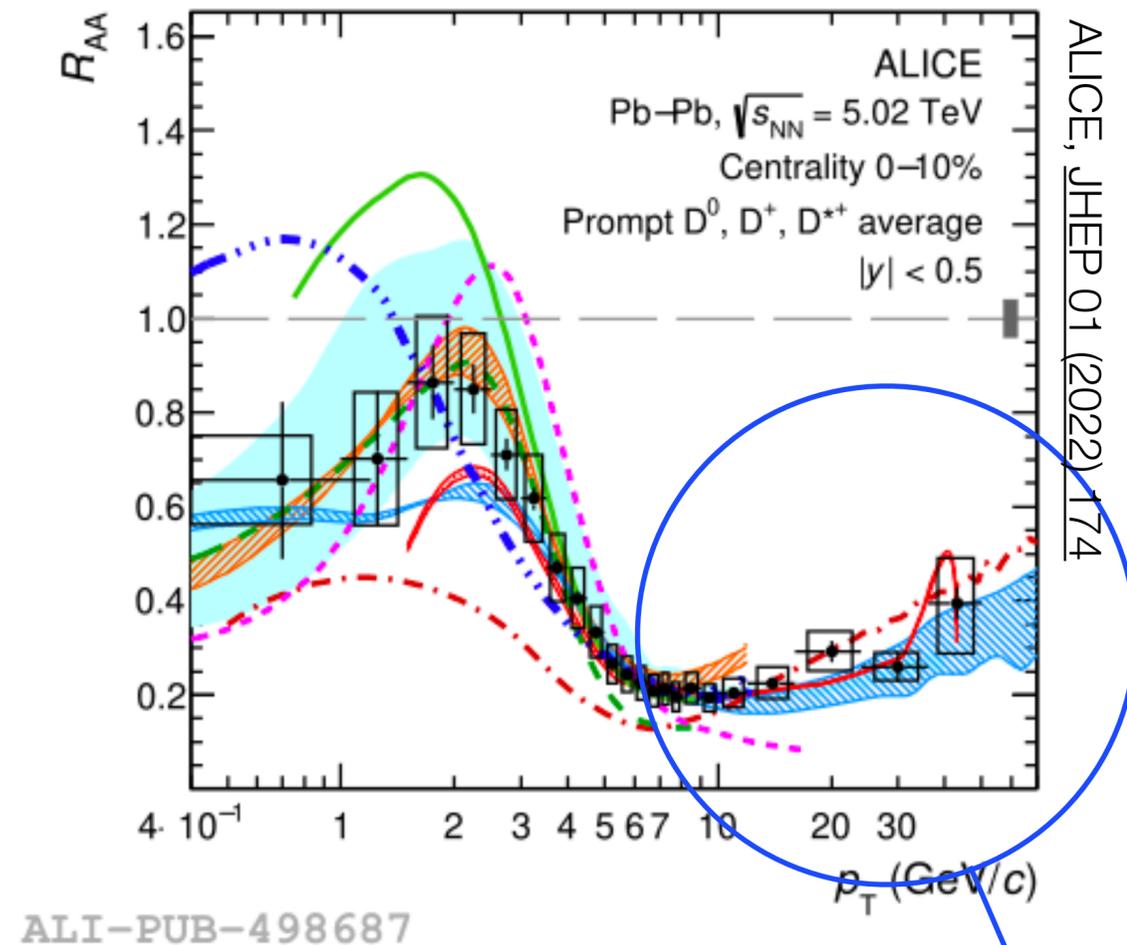
Elliptic flow v_2



Nuclear modification and elliptic flow of D mesons

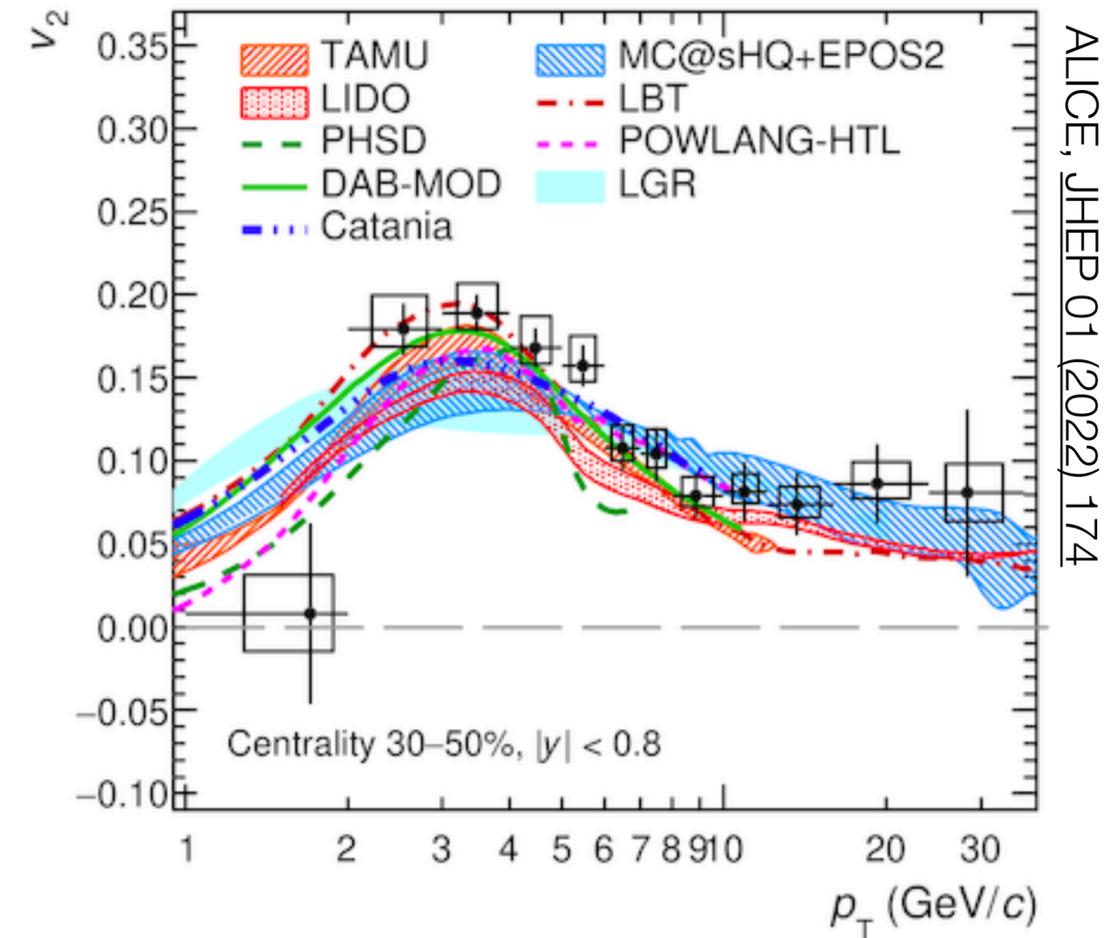
D mesons contain a charm quark $m \gg T$, that is produced in an initial hard scattering

Nuclear modification factor



High- p_T suppression:
due to energy loss

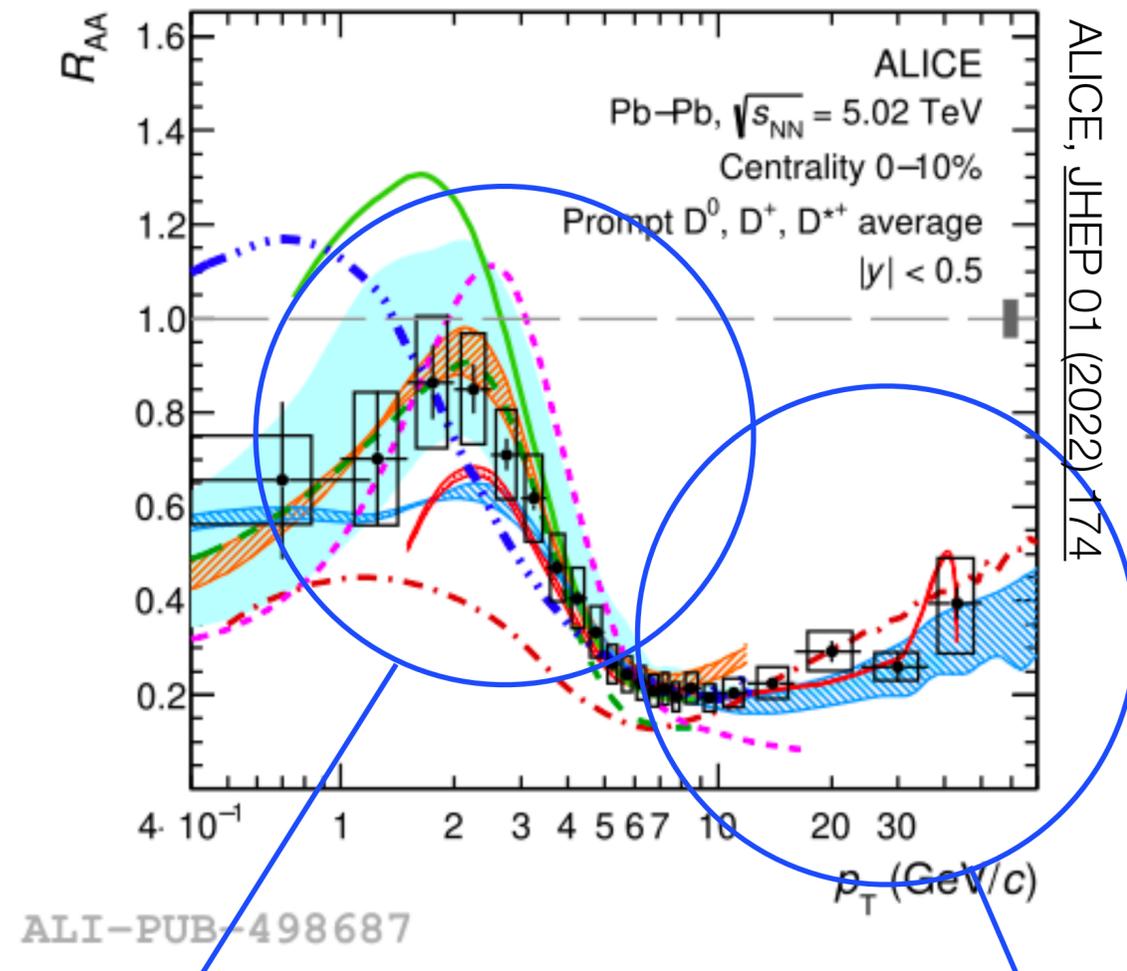
Elliptic flow v_2



Nuclear modification and elliptic flow of D mesons

D mesons contain a charm quark $m \gg T$, that is produced in an initial hard scattering

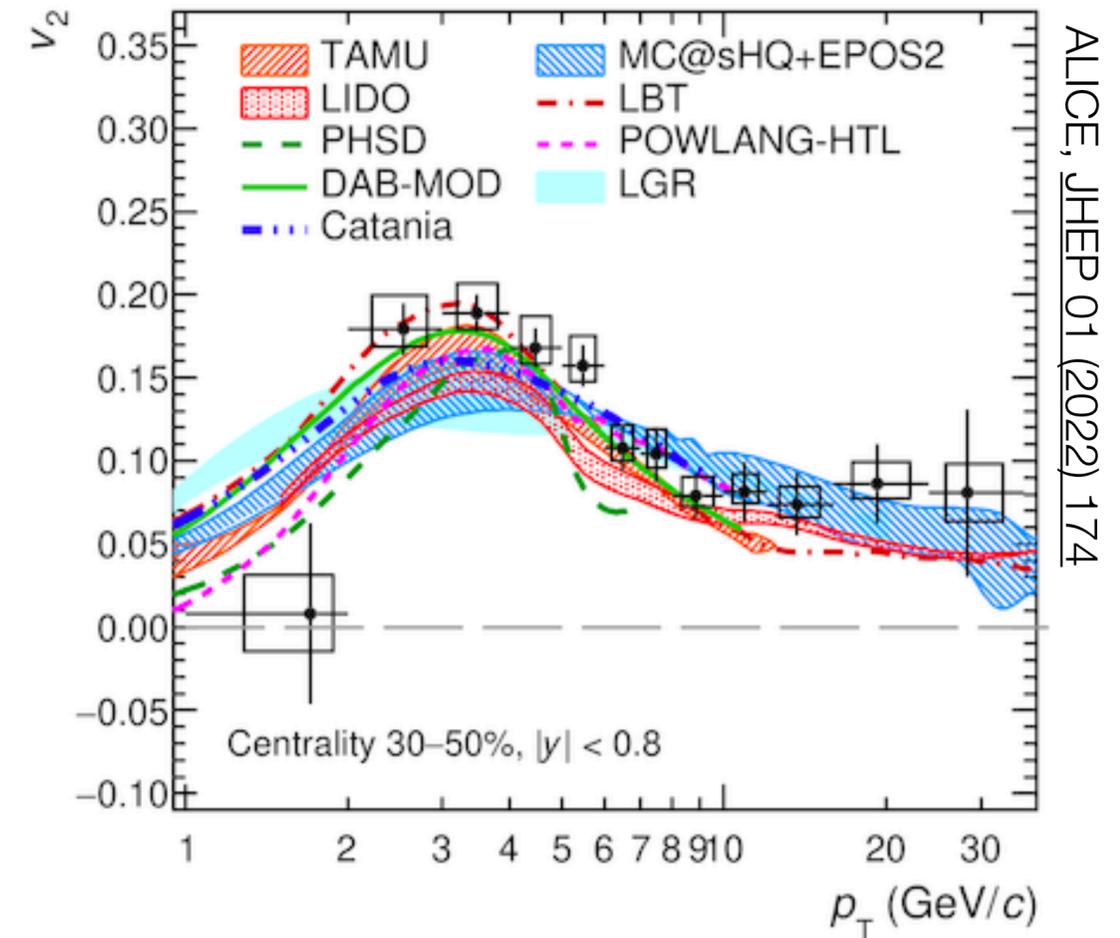
Nuclear modification factor



Low p_T : no change/enhancement:
charm conservation + diffusion

High- p_T suppression:
due to energy loss

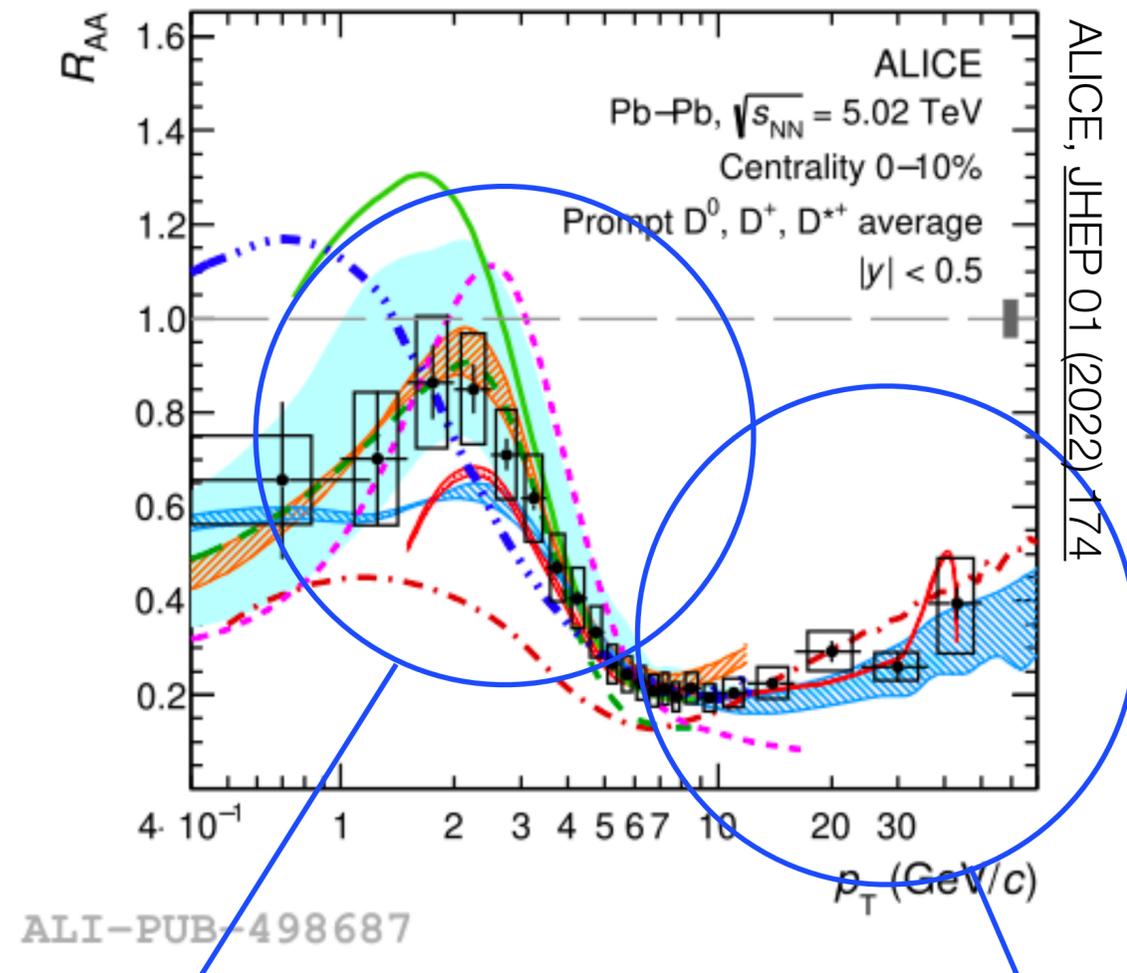
Elliptic flow v_2



Nuclear modification and elliptic flow of D mesons

D mesons contain a charm quark $m \gg T$, that is produced in an initial hard scattering

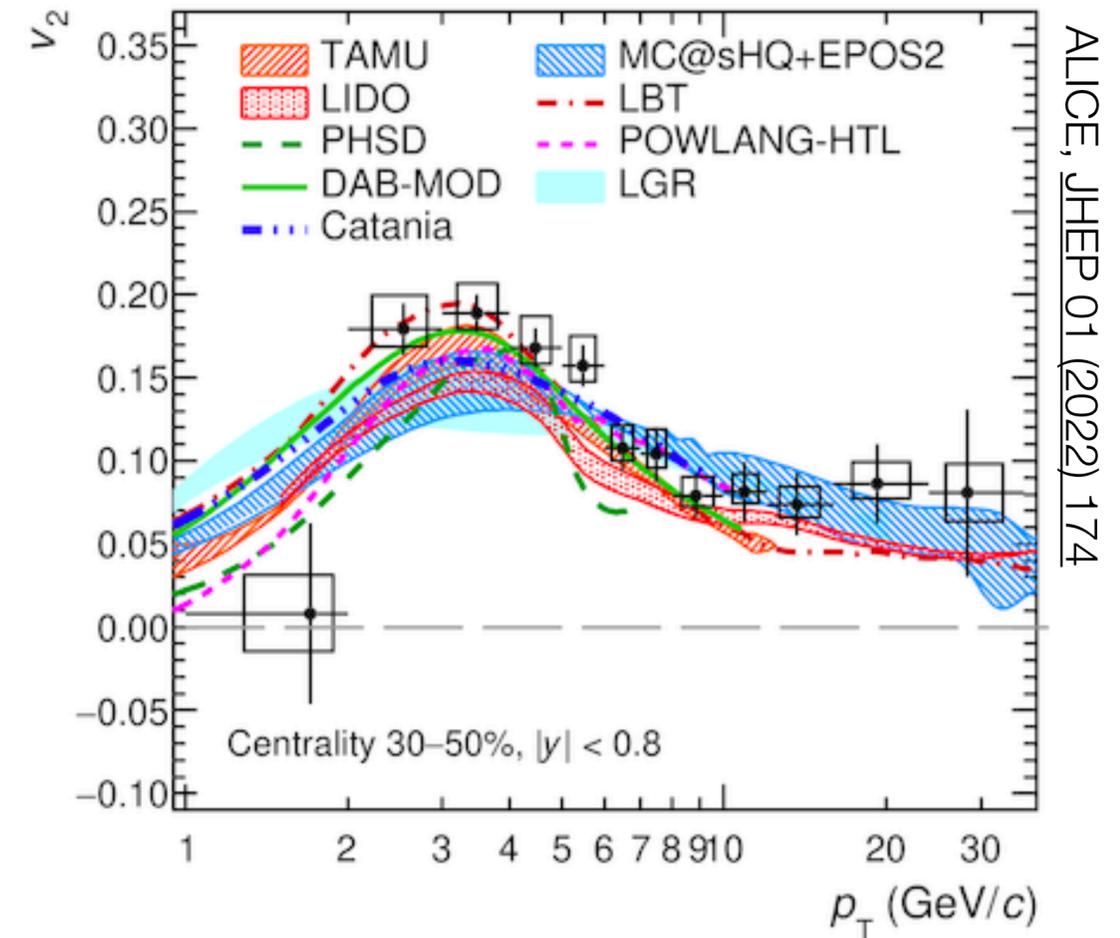
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due to energy loss

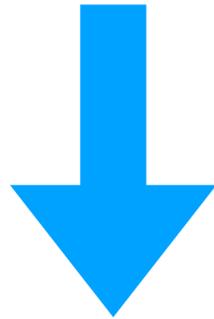
Elliptic flow v_2



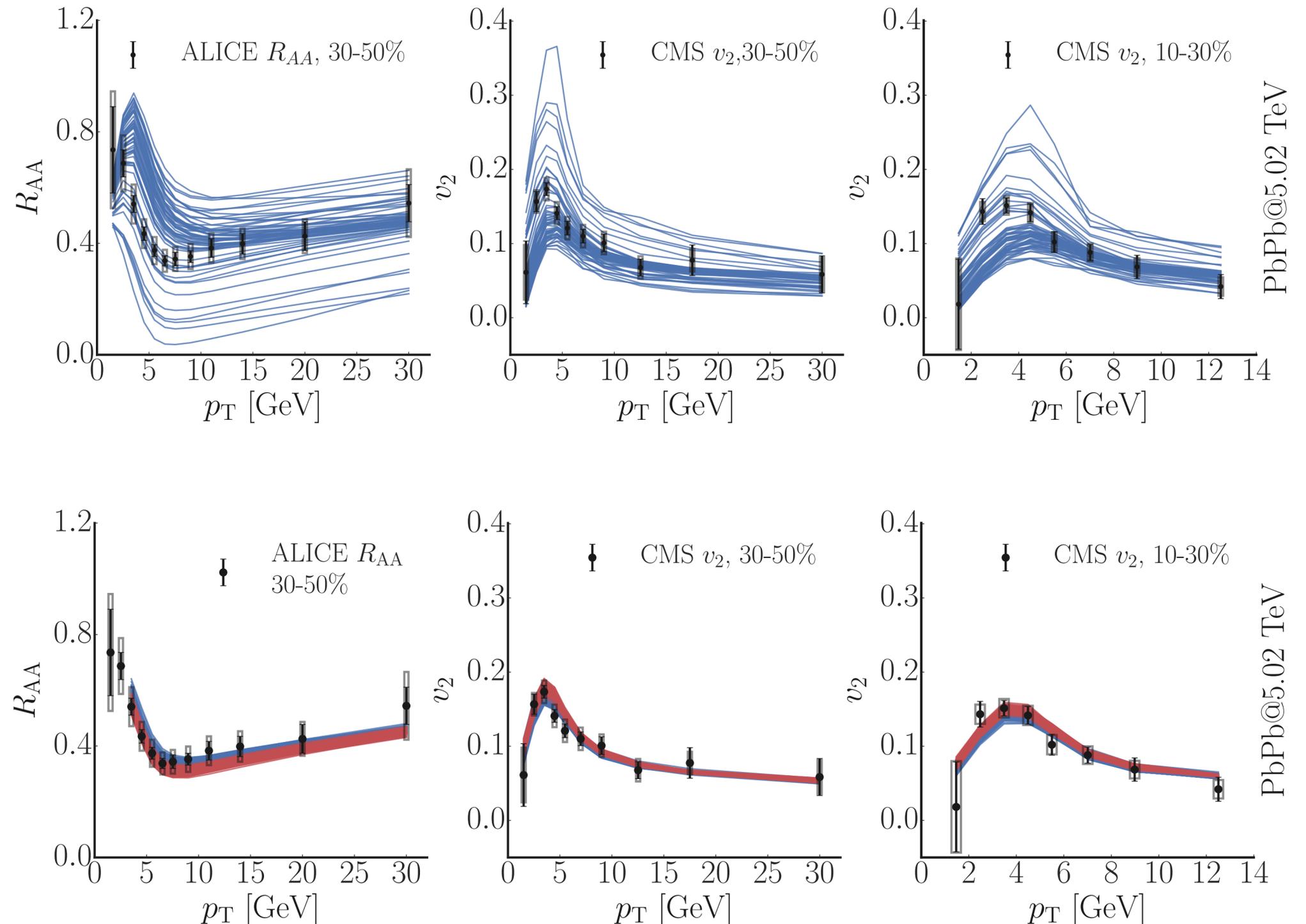
Azimuthal anisotropy:
Full effect generated by interactions

Determining the transport coefficients

Run model for different parameter settings
interpolate with Gaussian process emulator

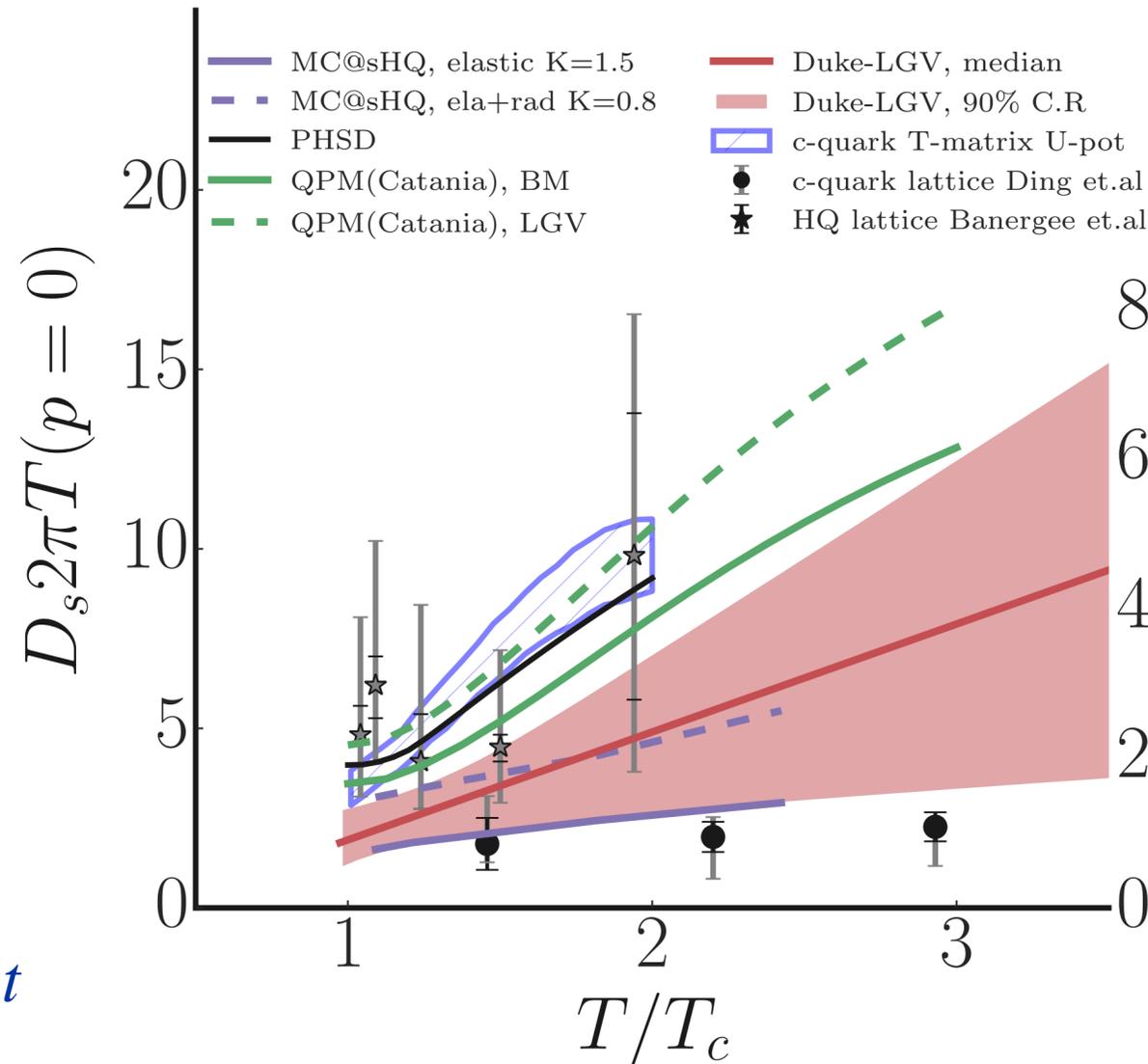


Posterior
Range of model settings that agree with data

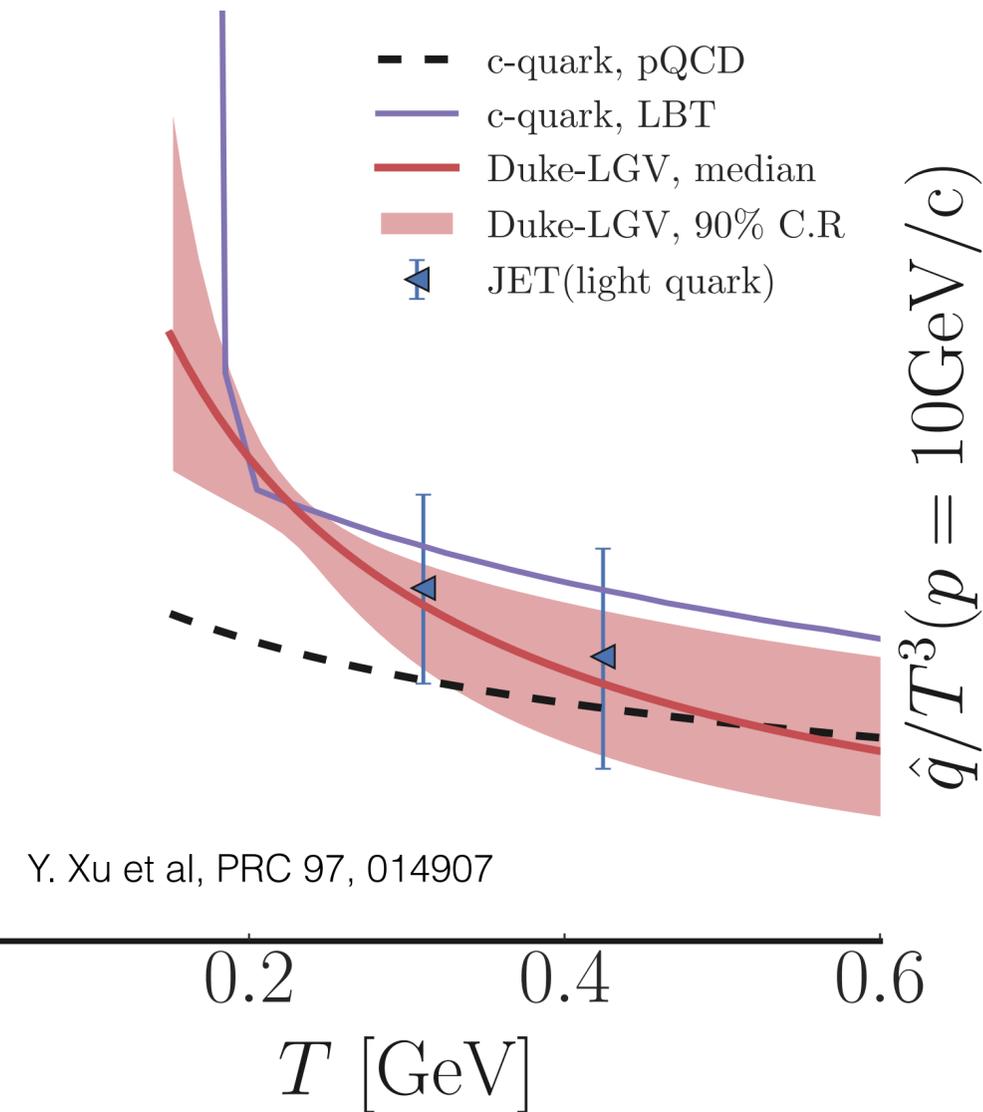


Heavy flavor transport coefficient: Bayesian fit

Diffusion coefficient D_s



Transport coefficient \hat{q}



$$\langle r^2 \rangle = 6 D_s t$$

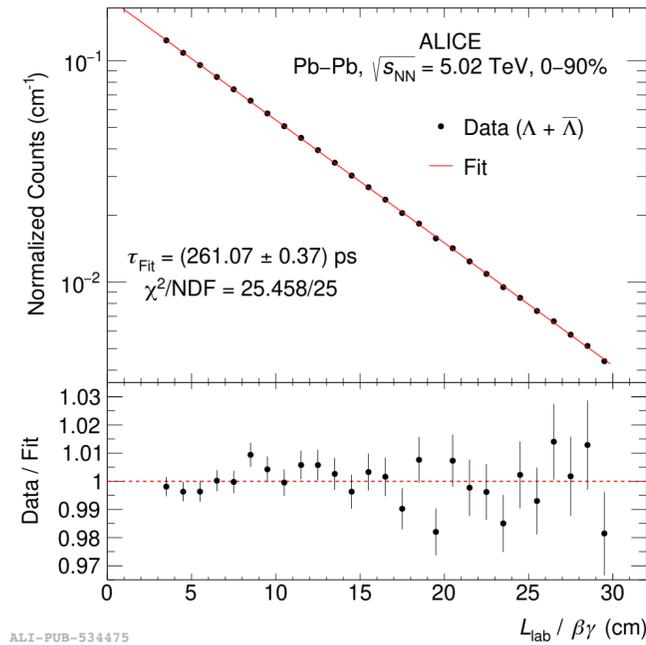
$$\hat{q} = \frac{\langle q_{\perp}^2 \rangle}{\lambda}$$

Data constrain transport properties of the QGP
 Results agree with lattice QCD/pQCD expectations
 and between light and heavy flavour sector

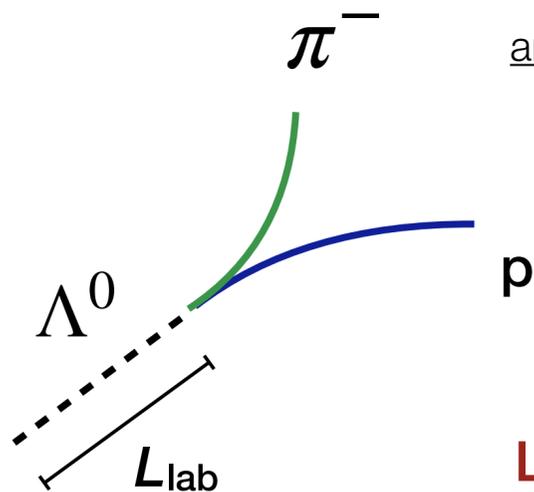
Heavy-ion collisions as a laboratory for nuclear and hadron physics

Example: life time of strange baryons and nuclei

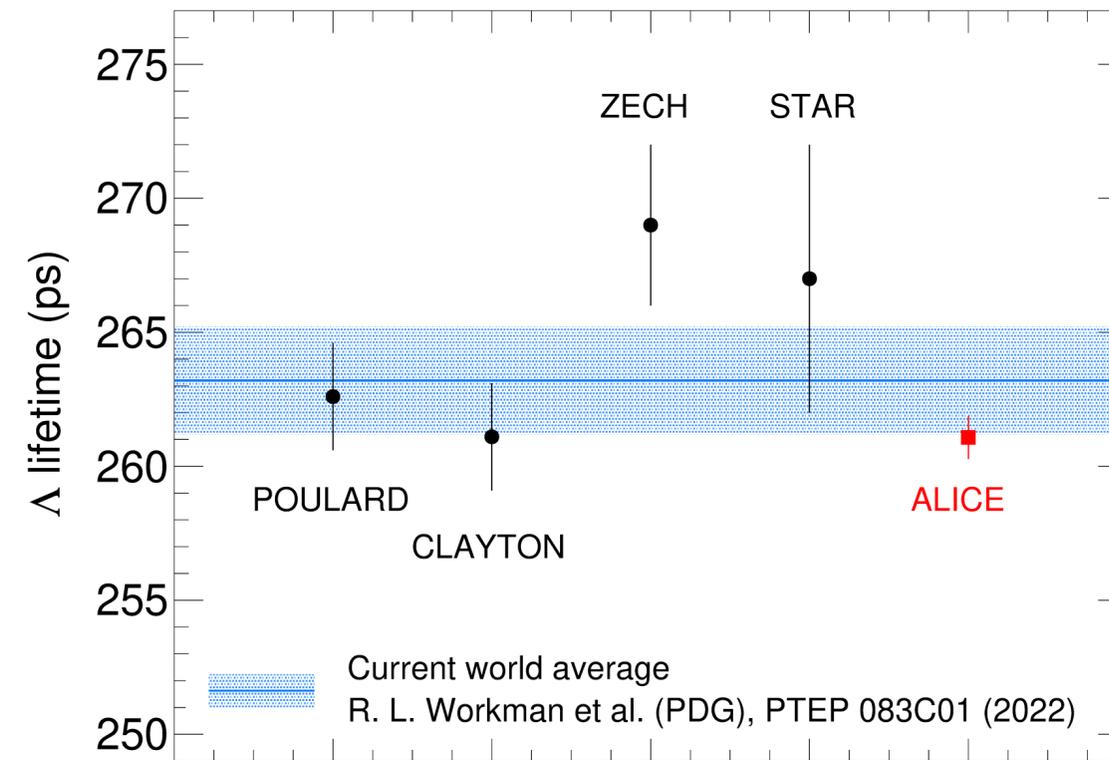
Proper decay length distribution



[arXiv:2303.00606](https://arxiv.org/abs/2303.00606)

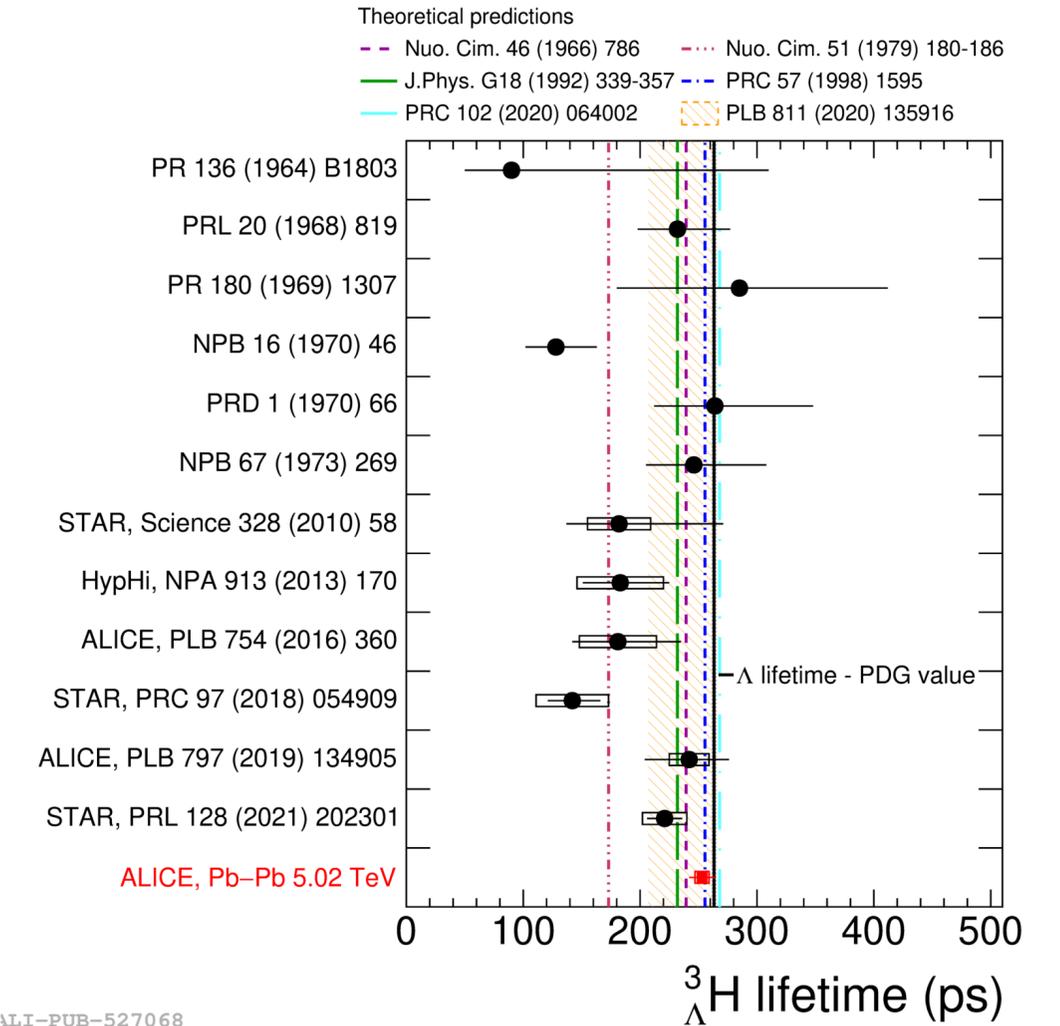


Life time: comparison to existing results

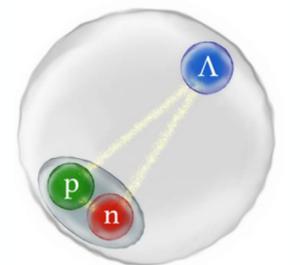


$$\tau_{\Lambda+\bar{\Lambda}} = [261.07 \pm 0.37(\text{stat.}) \pm 0.72(\text{syst.})] \text{ ps.}$$

Hypernuclei life time



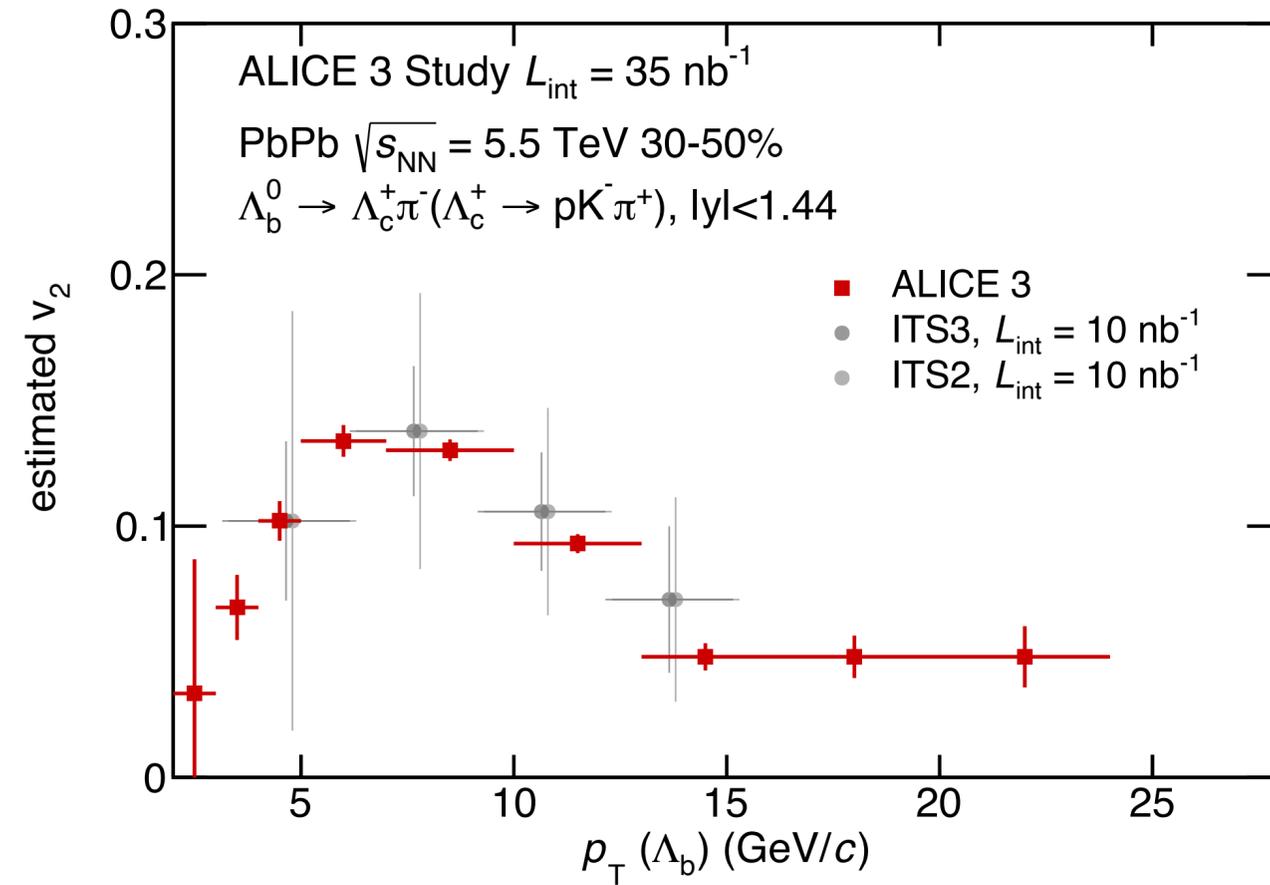
[arXiv:2209.07360](https://arxiv.org/abs/2209.07360)



Life time measurements of hyperons and hypernuclei competitive with world data

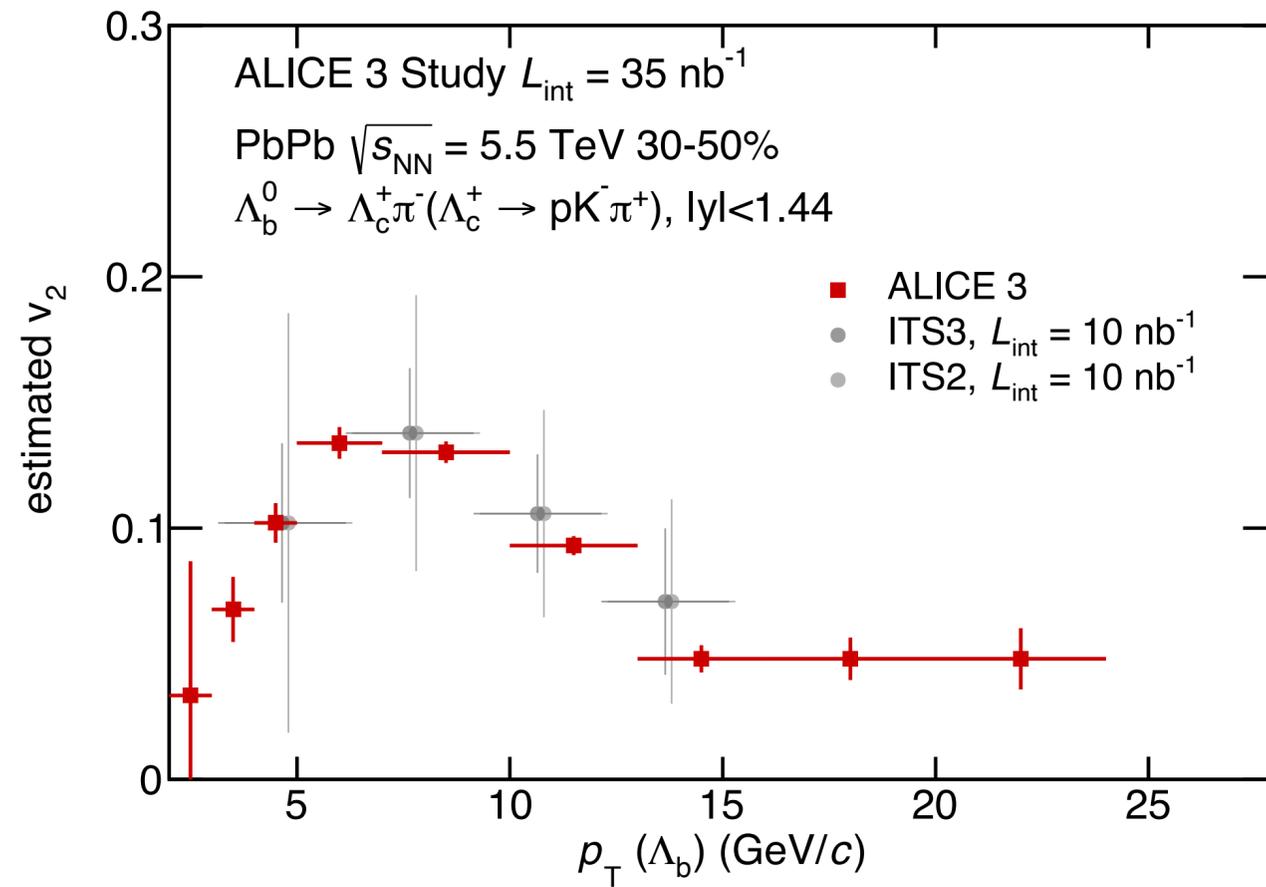
Thermalisation of heavy quarks

Λ_b v_2 performance

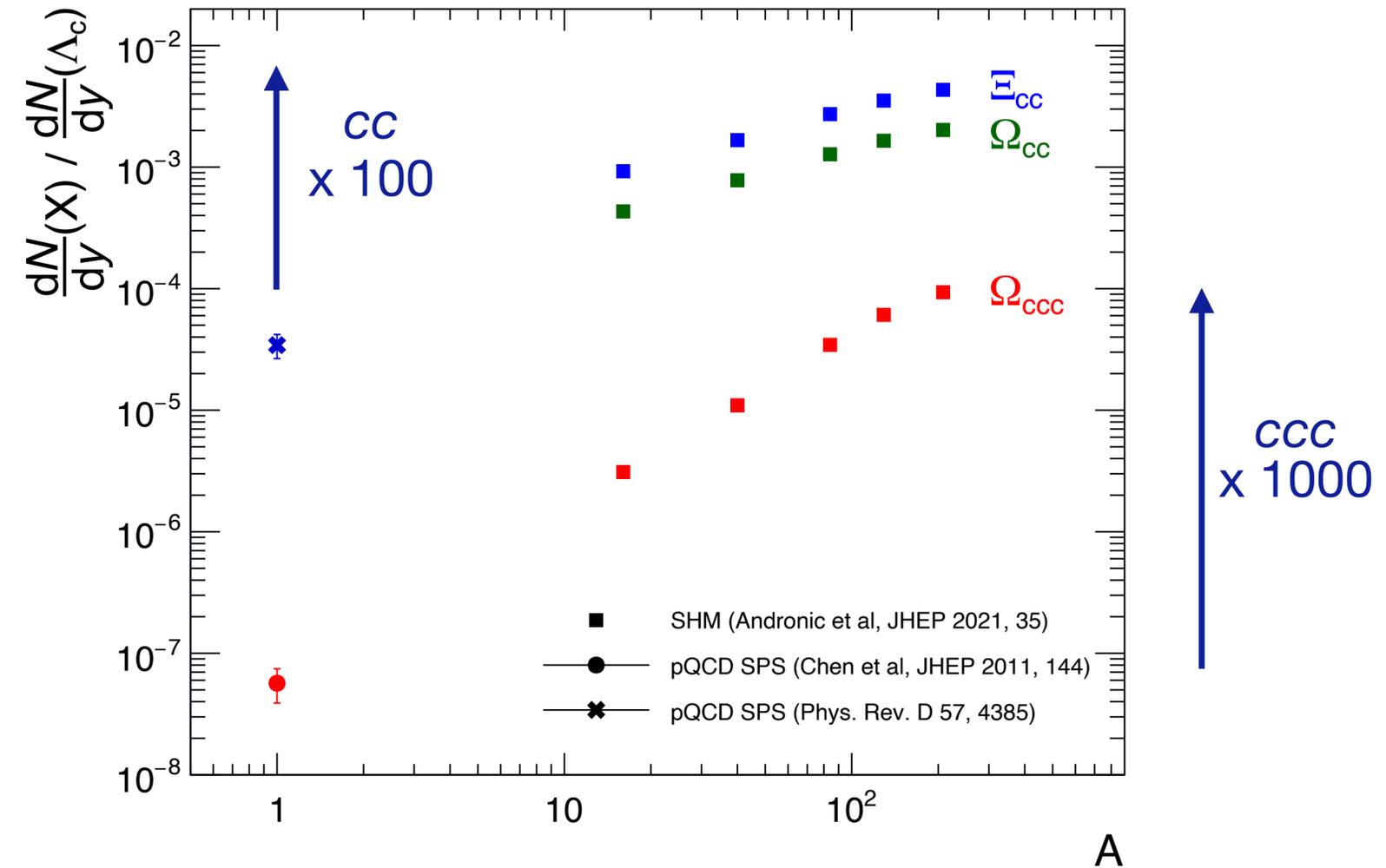


Thermalisation of heavy quarks

Λ_b v_2 performance



Multi-charm baryon yields



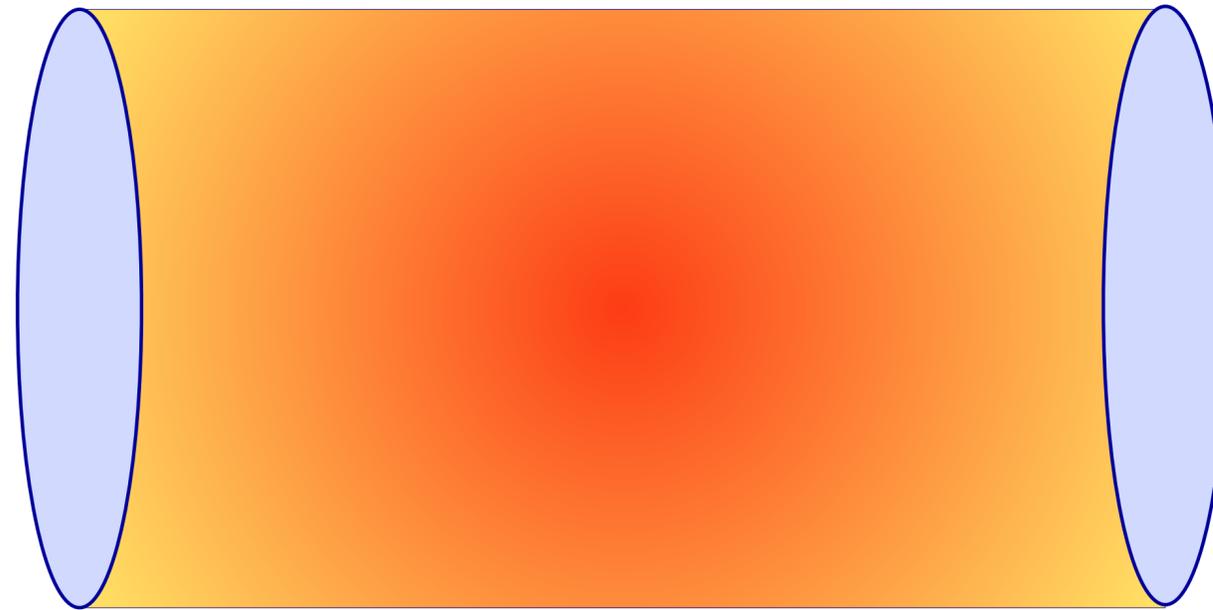
Multi-charm baryons: unique probe

- Large expected enhancement
- Theoretically clean: charm quarks conserved

Messengers of the Plasma: soft and hard processes

Soft probes: particles produced by the QGP

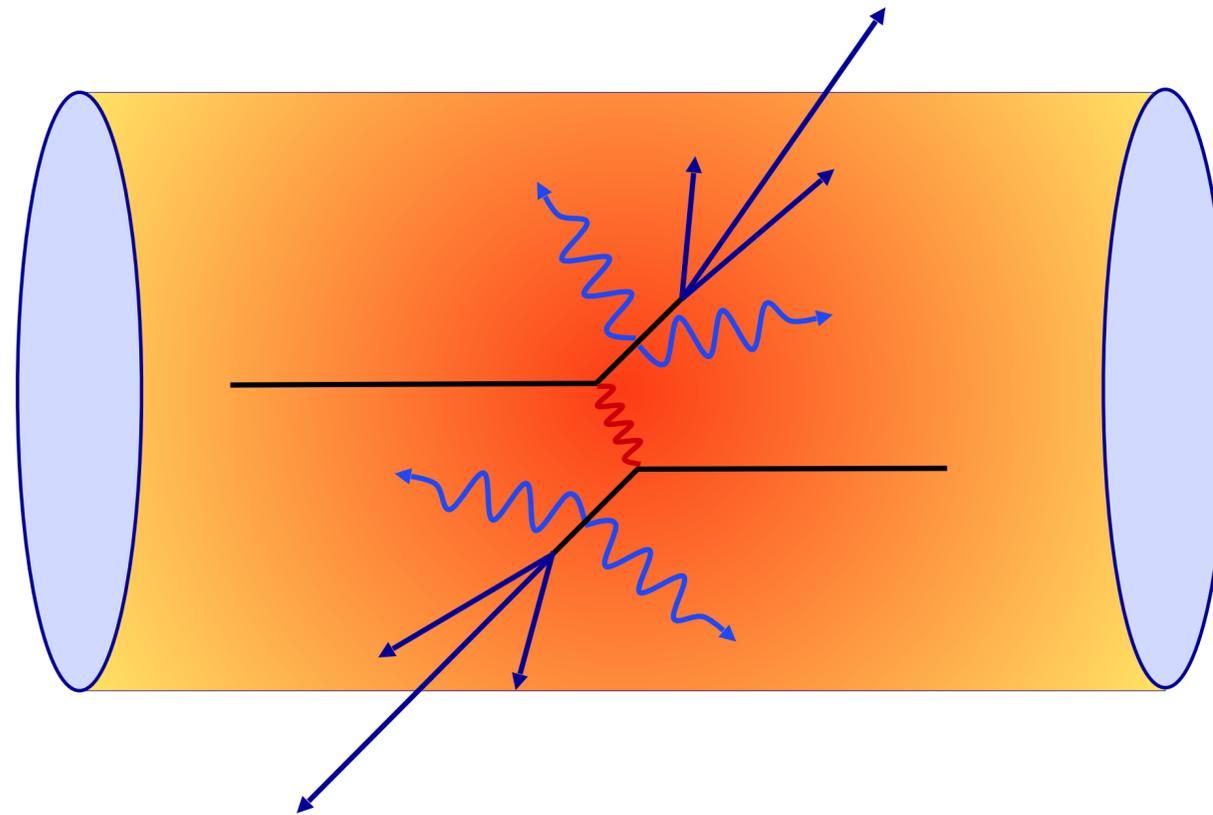
- Azimuthal anisotropy
- Light-flavour particle ratios
- Thermal radiation



Messengers of the Plasma: soft and hard processes

Soft probes: particles produced by the QGP

Azimuthal anisotropy
Light-flavour particle ratios
Thermal radiation



Hard scattering products probe the QGP as they propagate out

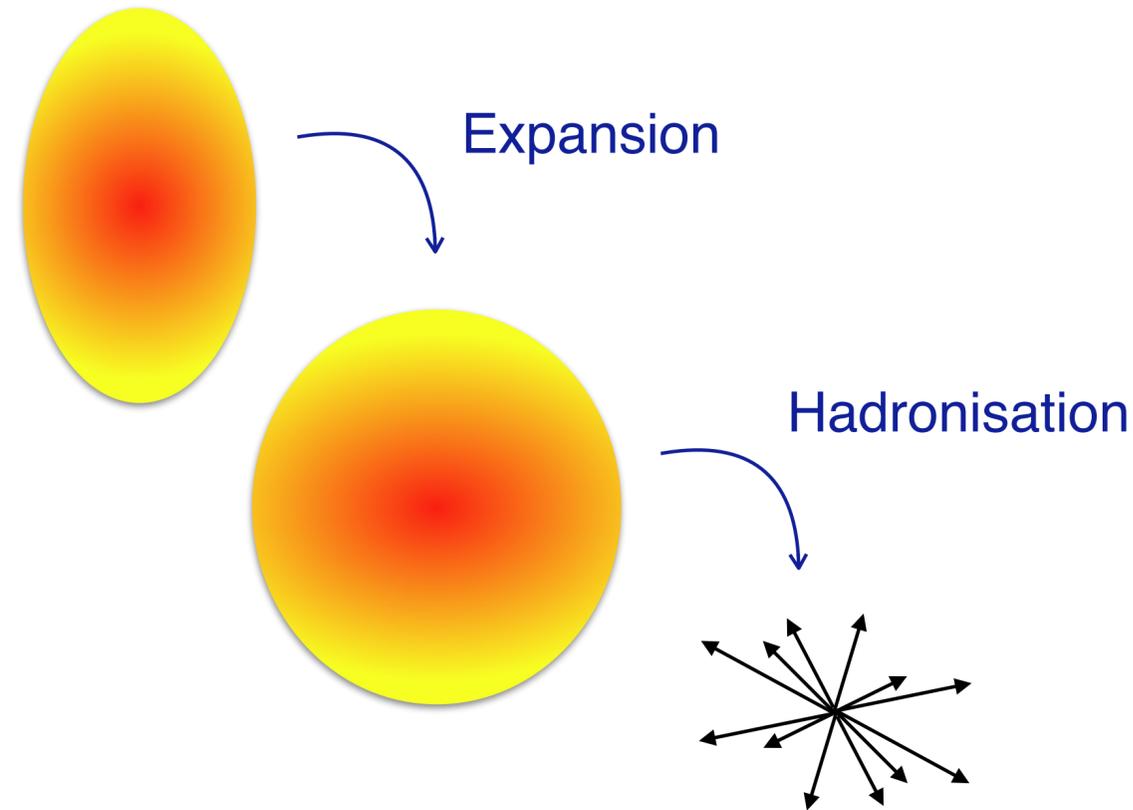
Heavy quarks charm and beauty:

- $m \gg T$: Only produced in initial hard scattering
- Flavour conserved during evolution

Azimuthal anisotropy: two mechanisms

Hydrodynamical expansion

Conversion of pressure gradients into momentum space anisotropy

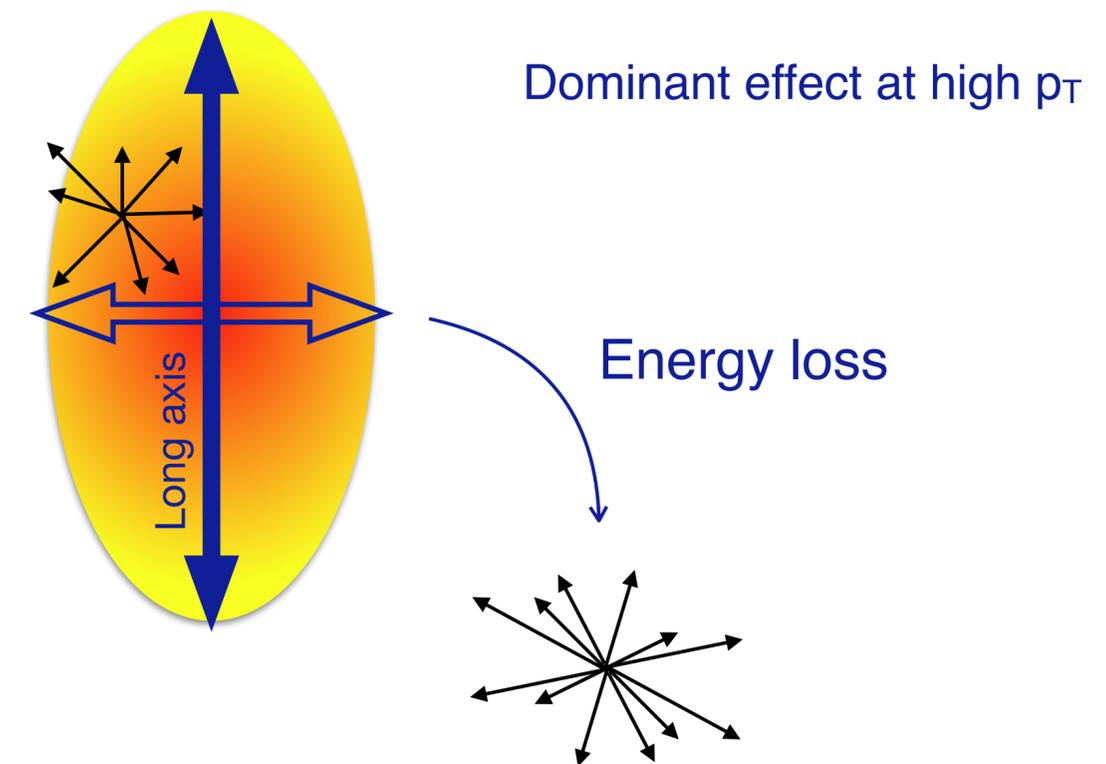


$$\nabla p = \rho \frac{d\vec{v}}{dt}$$

Dominant effect for late formation times:
light flavour at low p_T

Parton energy loss

Anisotropy due to energy loss and path length differences

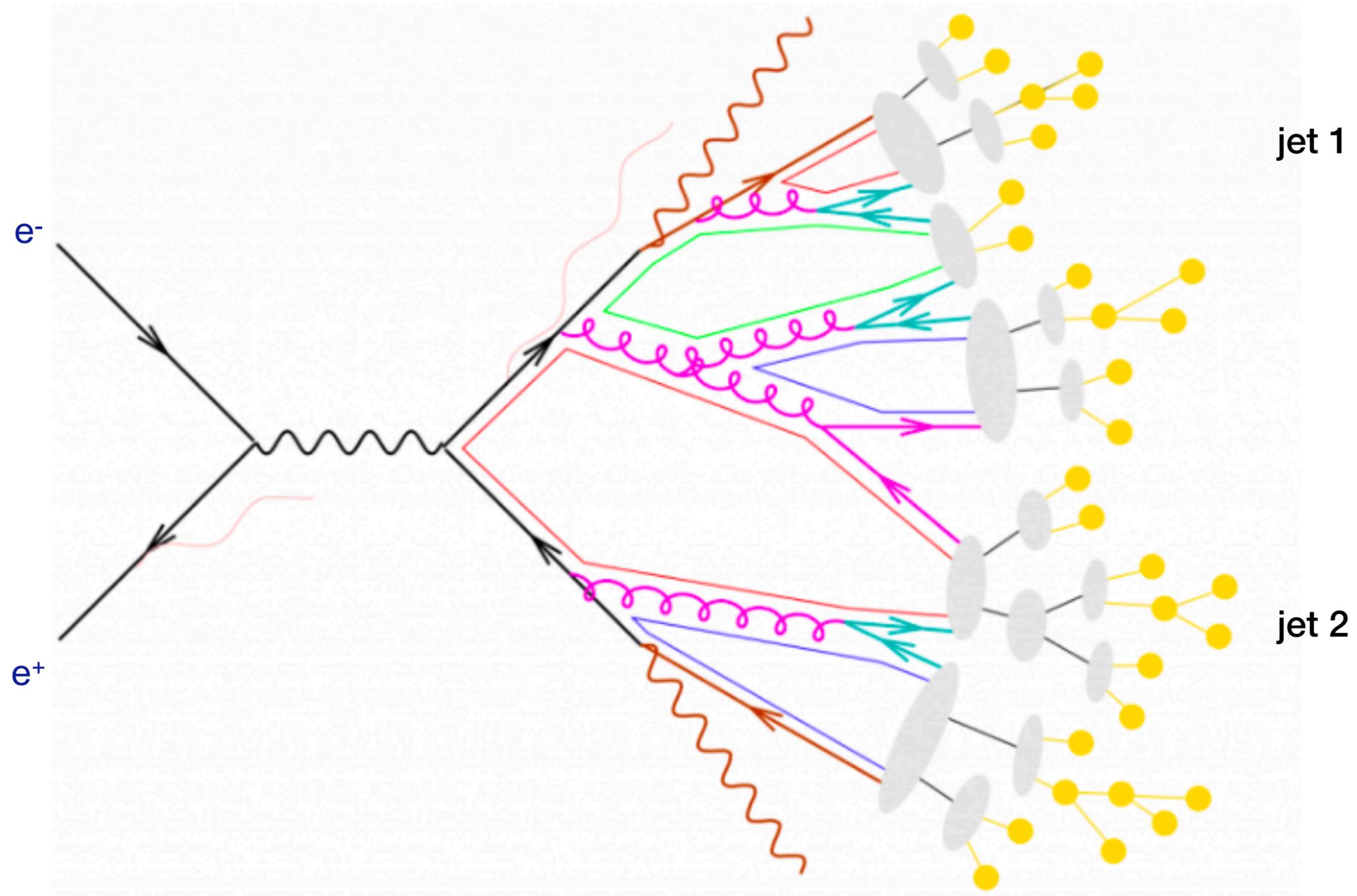


More energy loss along
long axis than short axis

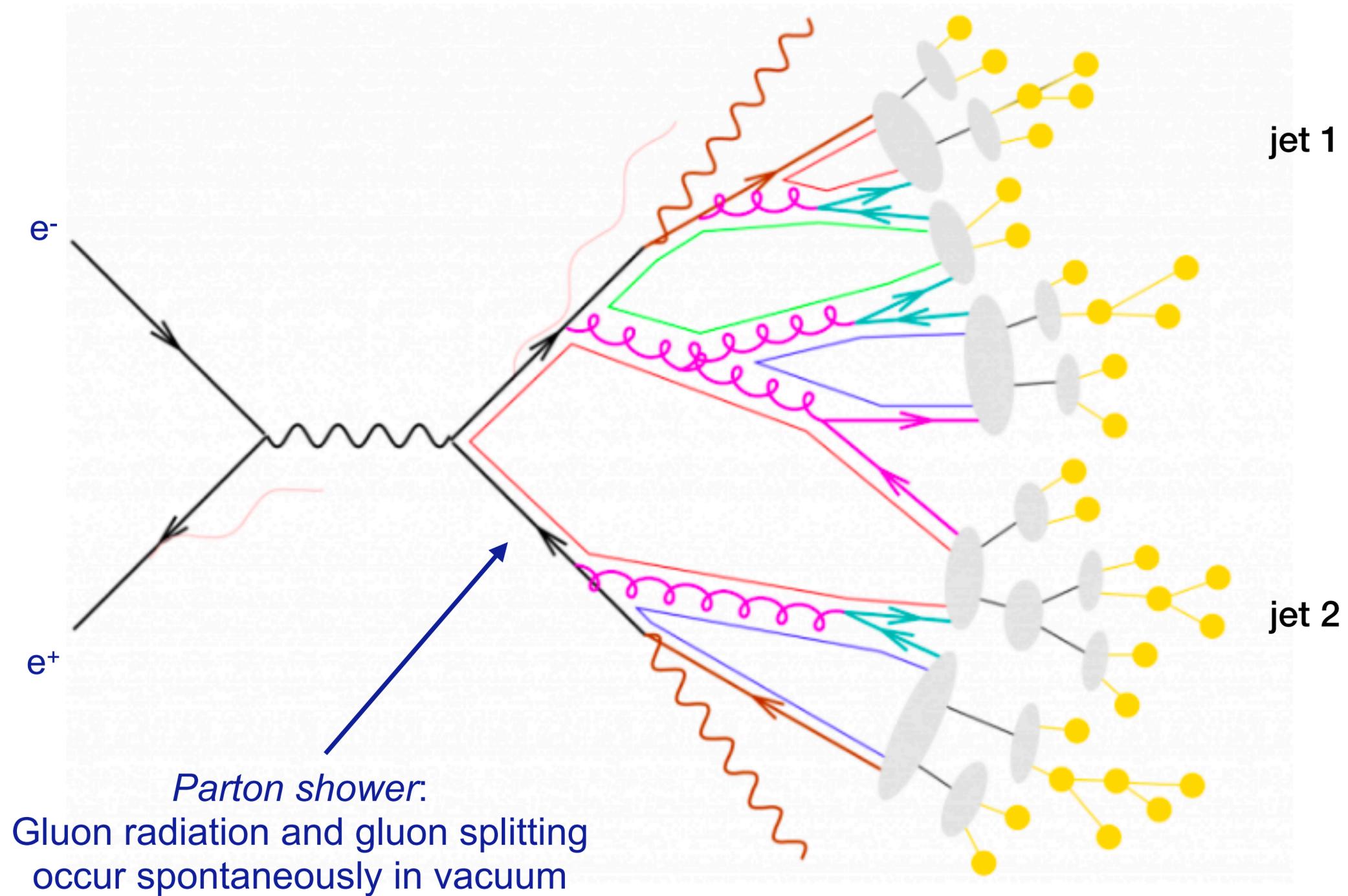
$$\Delta E_{med} \sim \alpha_s \hat{q} L^2$$

Dominant effect for early formation times:
heavy flavour, high p_T probes

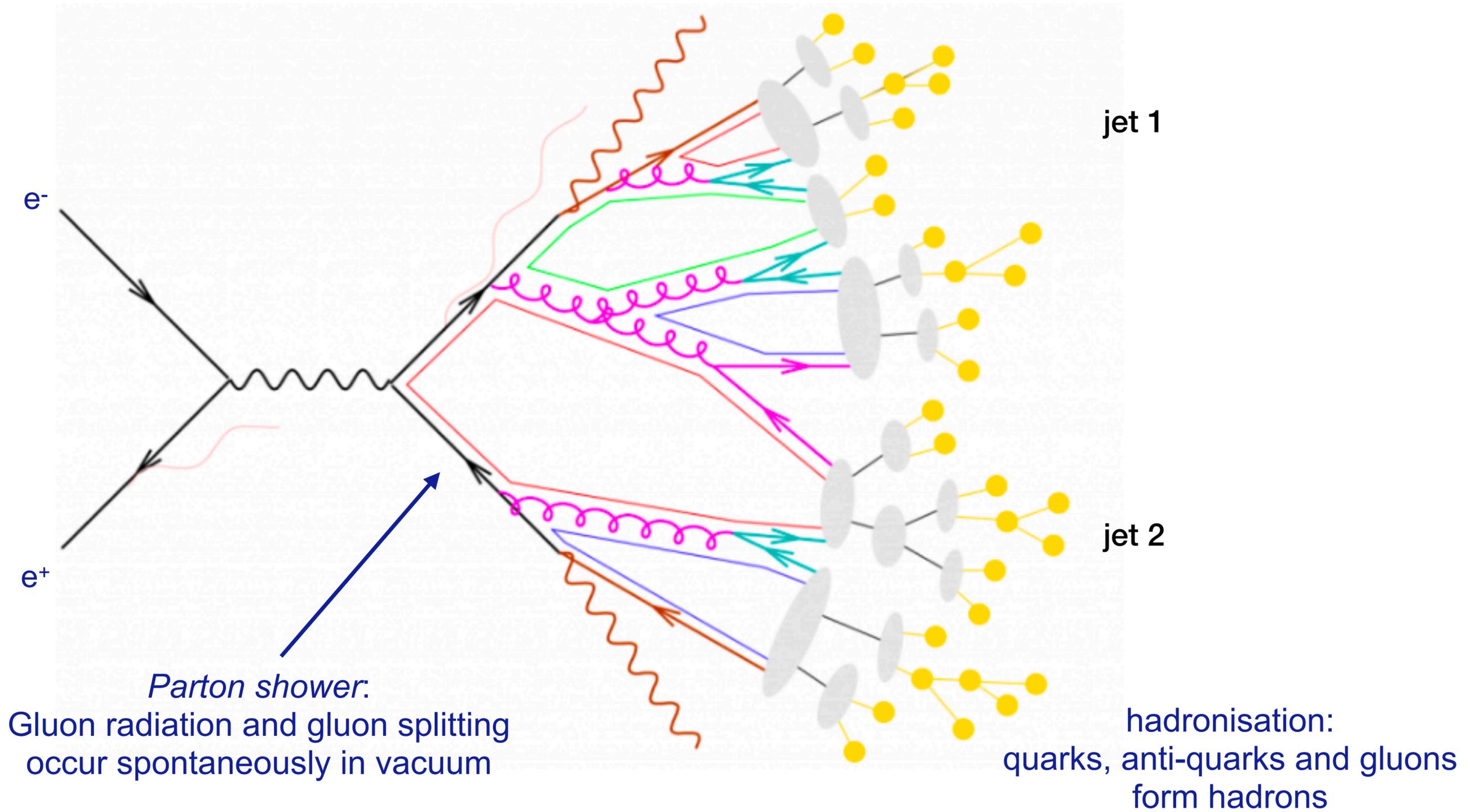
Jets and parton showers



Jets and parton showers

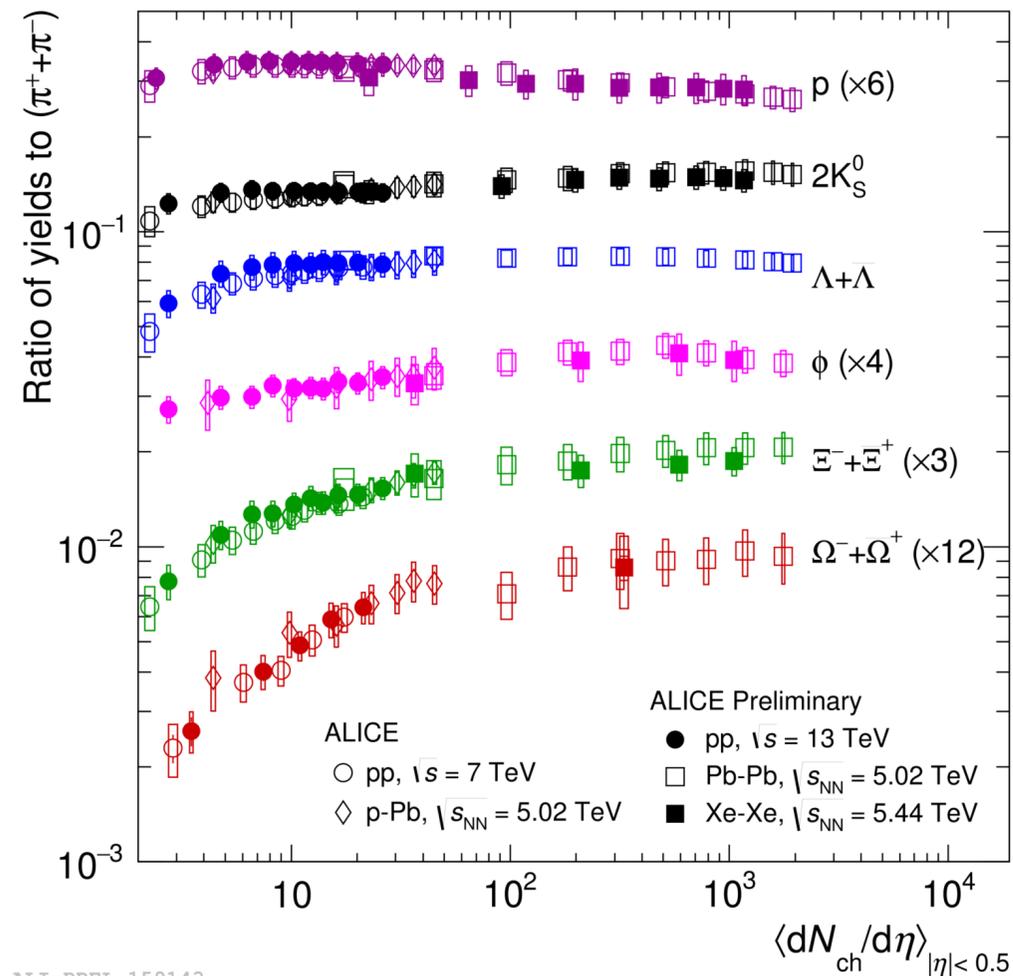


Jets and parton showers



Yields, mean p_T

Particle/pion ratios

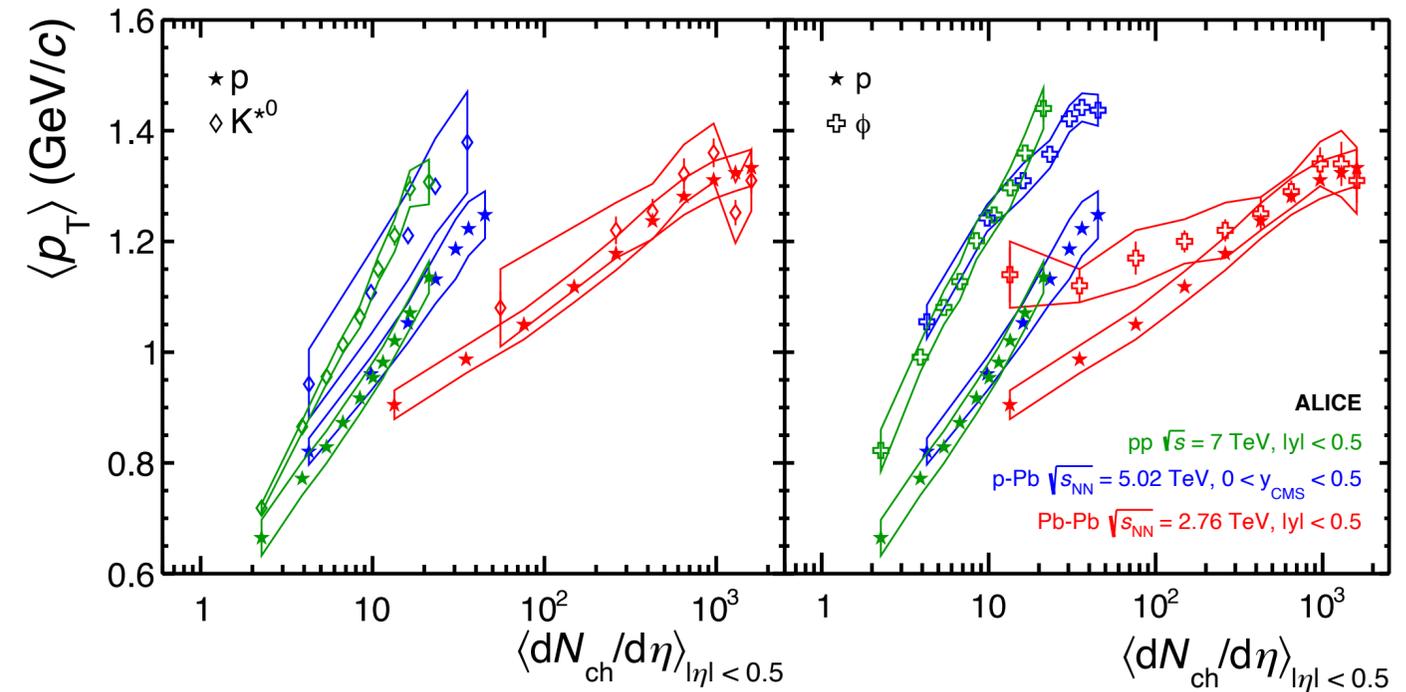


ALI-PREL-159143

Yields 'scale' with multiplicity

Yield, mean p_T are convenient to summarise results, but a lot of physics enters

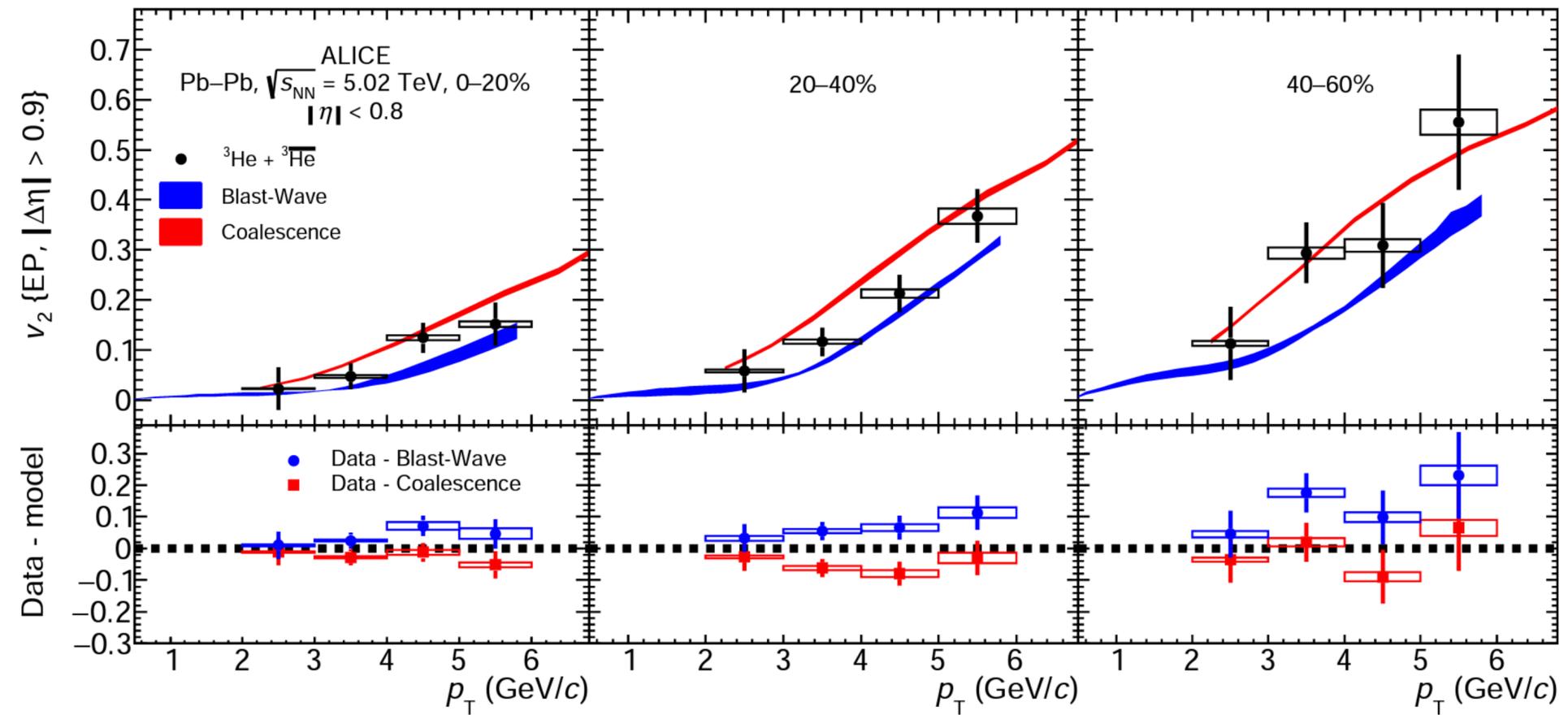
Mean p_T



Mean p_T : disconnect between high-mult small systems and PbPb

A multiplicity selection introduces more bias in small system

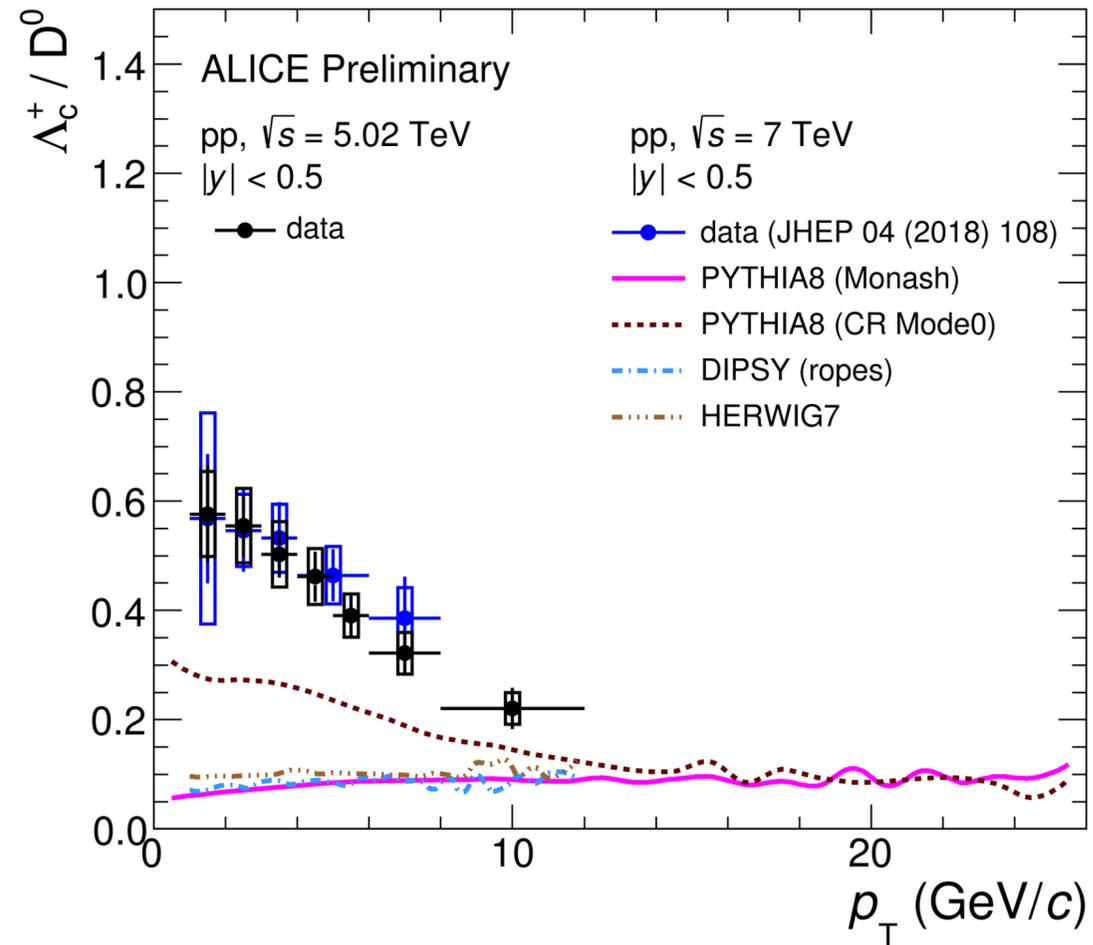
Flow of nuclei: coalescence vs thermal production



ALICE, arXiv:1910.09718

Λ_c production in pp and Pb-Pb

Λ_c/D in pp

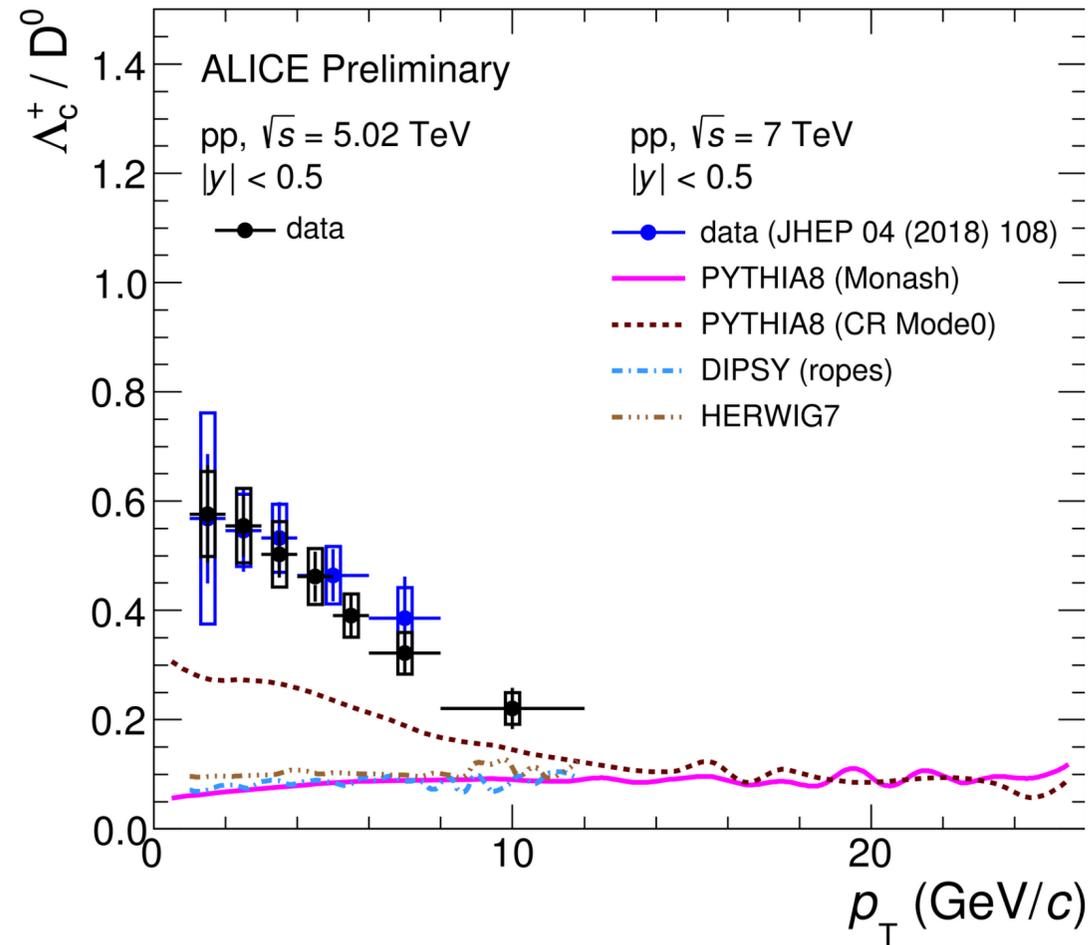


ALI-PREL-311156

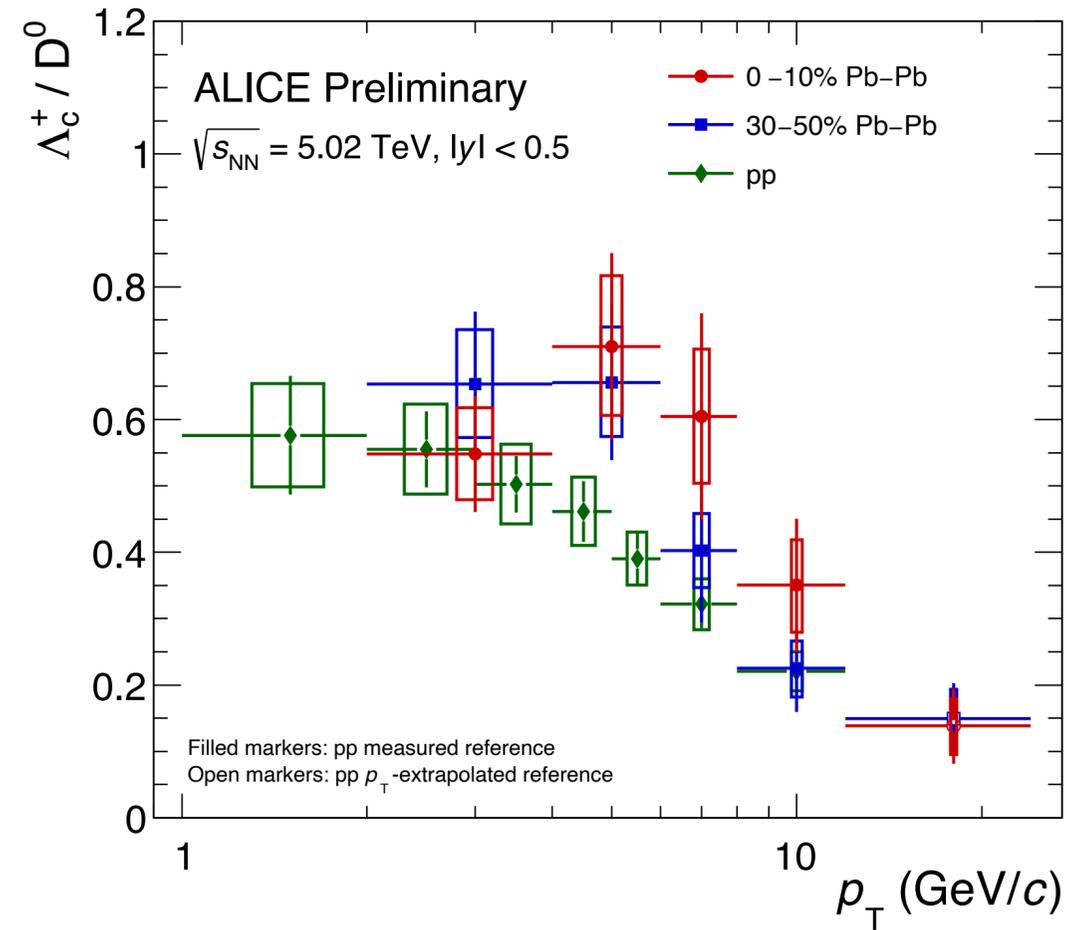
Λ_c/D^0 in pp significantly larger
than expected from e^+e^-

Λ_c production in pp and Pb-Pb

Λ_c/D in pp



Λ_c/D in pp, Pb-Pb

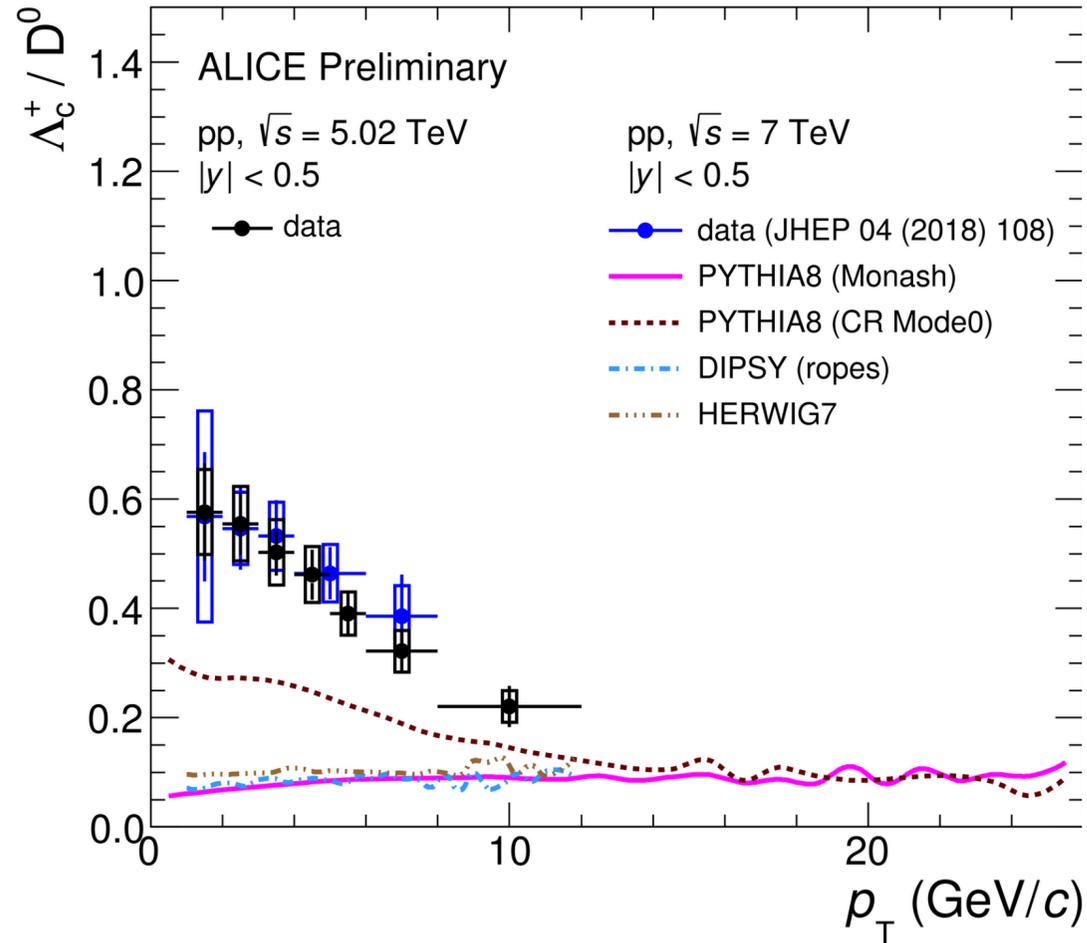


Λ_c/D^0 in pp significantly larger than expected from e^+e^-

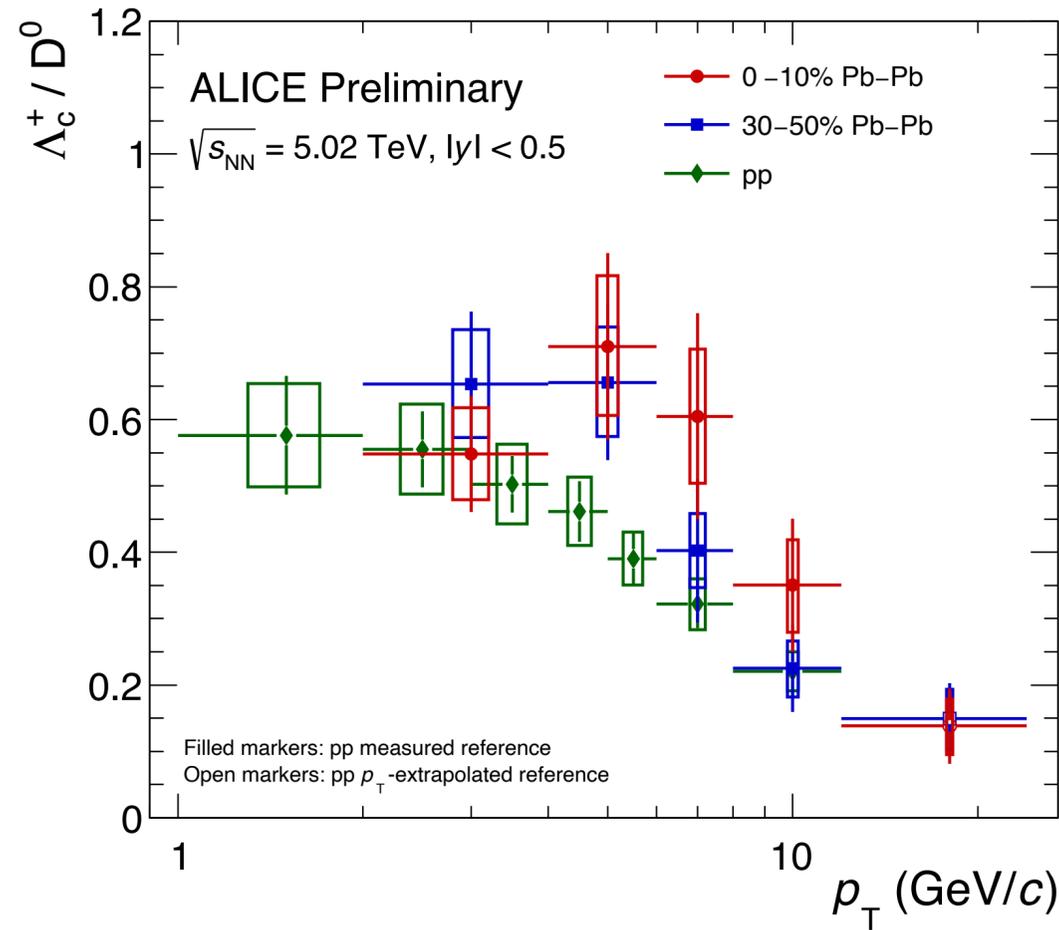
New result: Λ_c in Pb-Pb;
 Λ_c/D similar or slightly larger than in pp

Λ_c production in pp and Pb-Pb

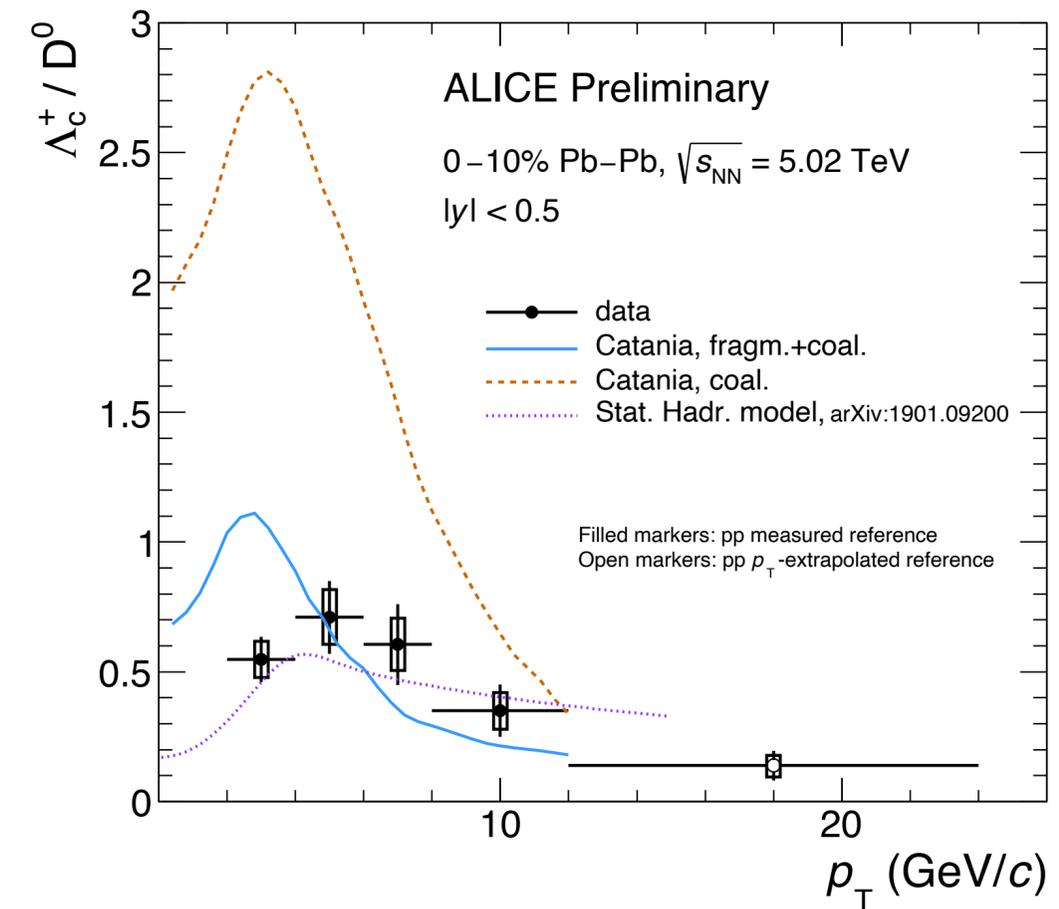
Λ_c/D in pp



Λ_c/D in pp, Pb-Pb



Λ_c/D in Pb-Pb



Λ_c/D^0 in pp significantly larger than expected from e^+e^-

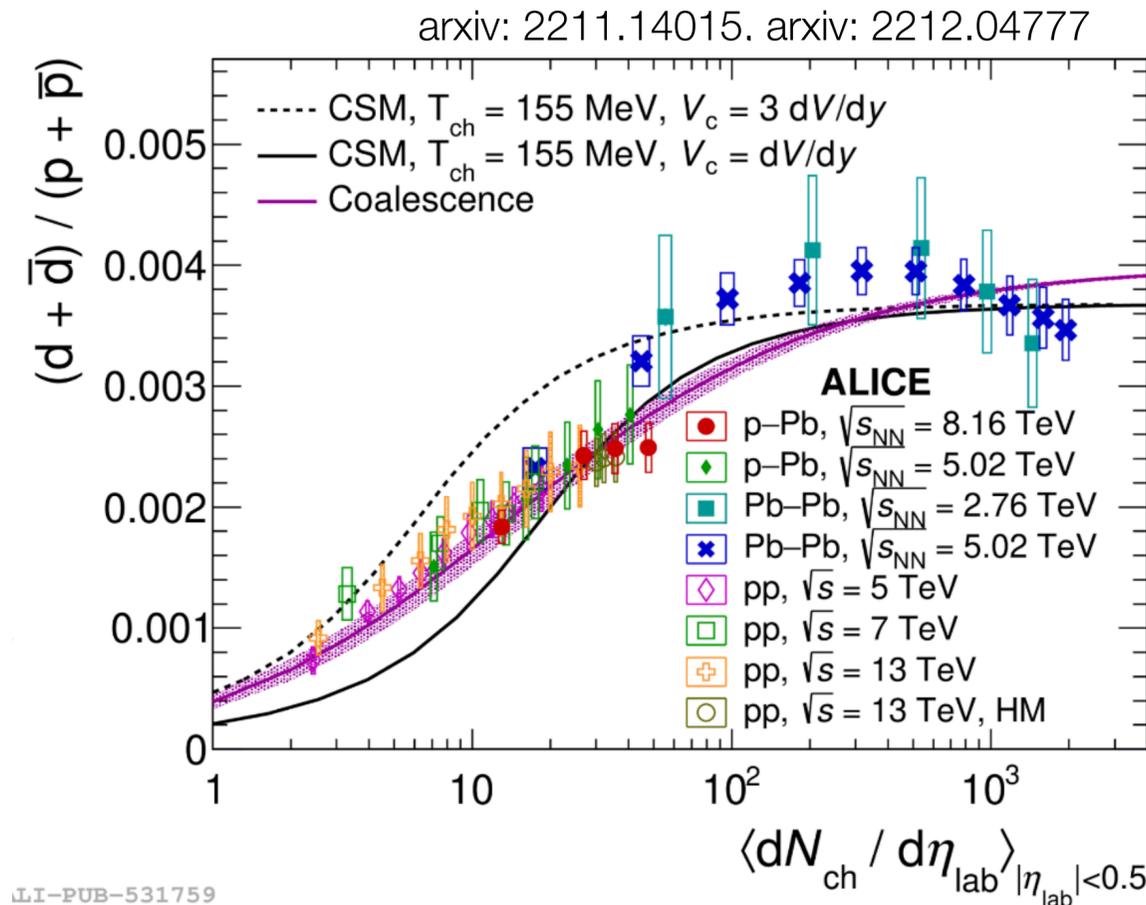
New result: Λ_c in Pb-Pb;
 Λ_c/D similar or slightly larger than in pp

Does hadronisation by recombination play a role? Or 'just' fragmentation?

Production of light (anti-)nuclei

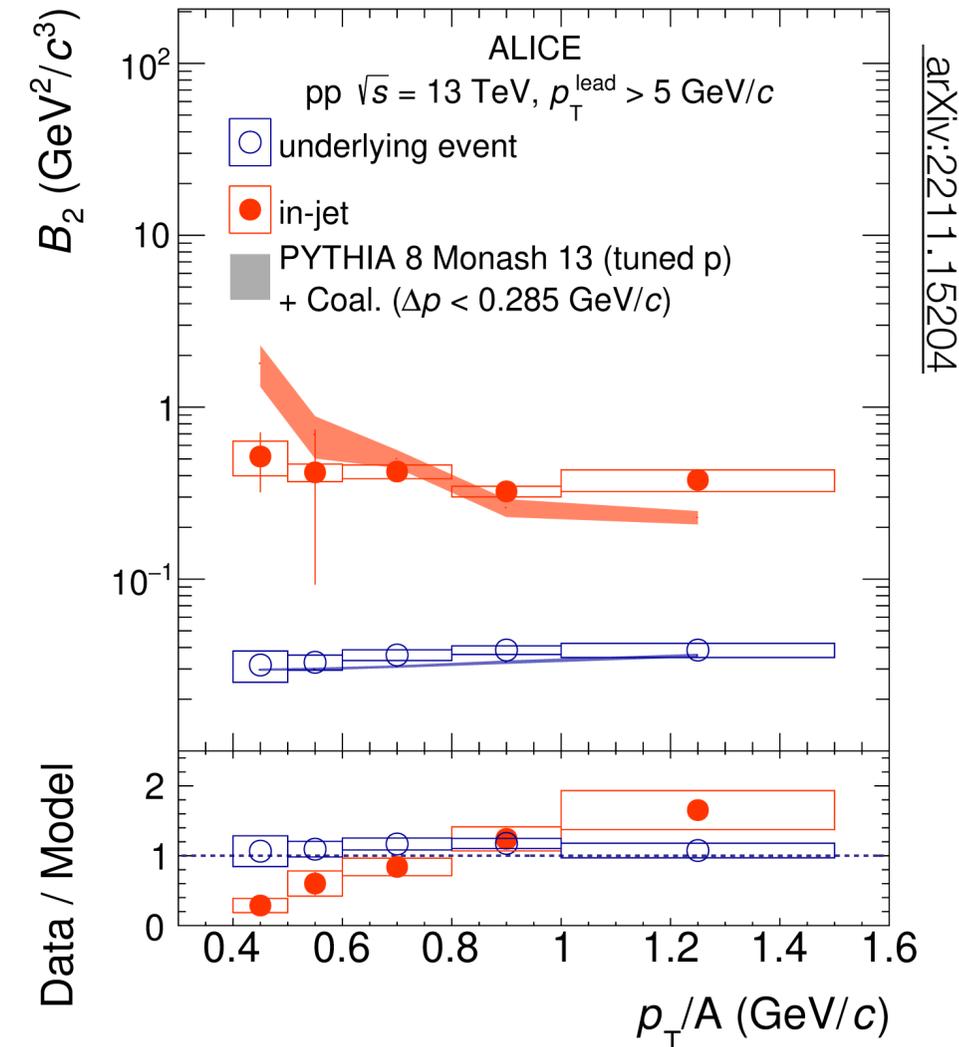
- Increased production in Pb-Pb collisions
- Two possible explanations:
 - Coalescence: density impacts coalescence rate
 - Thermal-statistical (CSM): baryon number conservation suppresses multi-baryon states in pp

Deuteron production rate in pp, Pb-Pb



- New results in pp, p-Pb collisions as a function of multiplicity
 - Smooth trend as a function of multiplicity
 - Data more in line with coalescence model

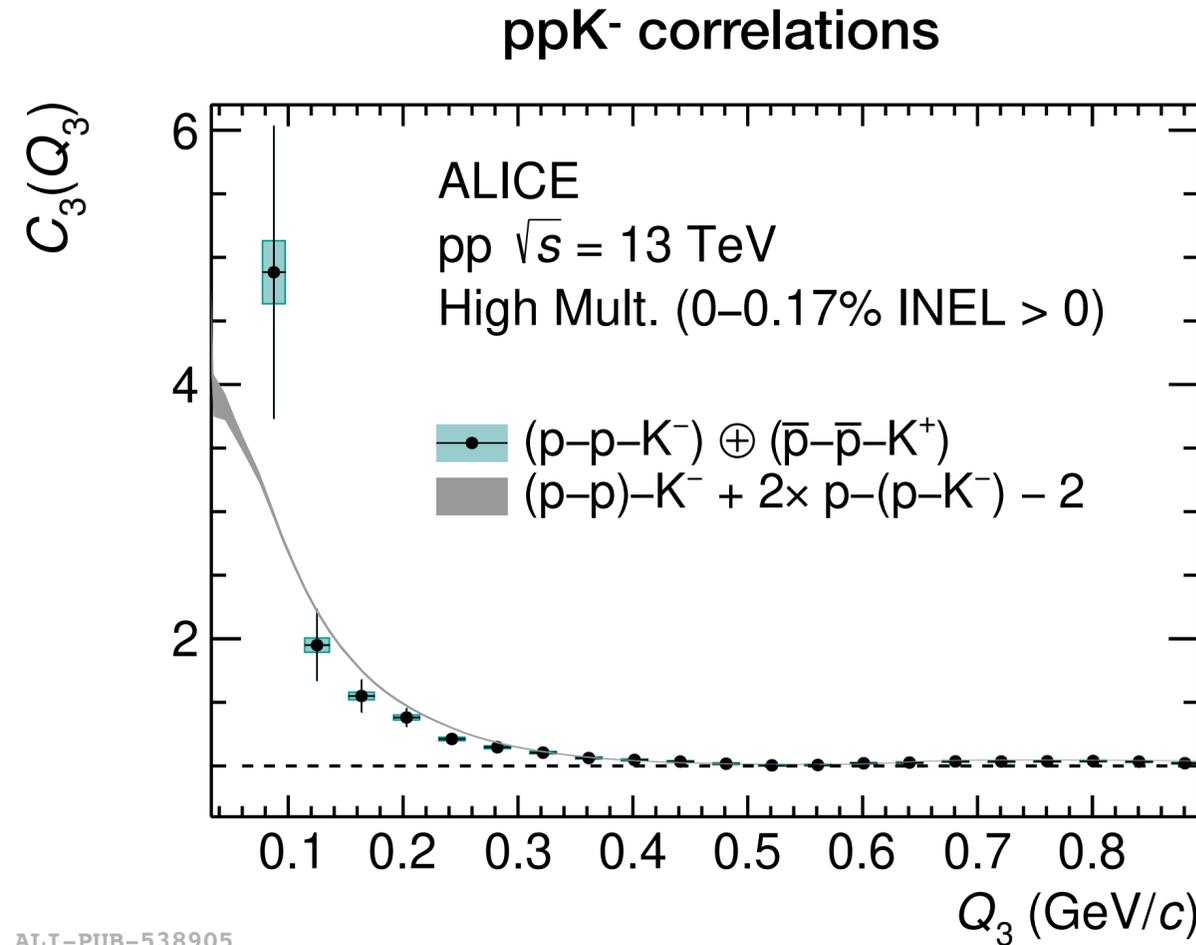
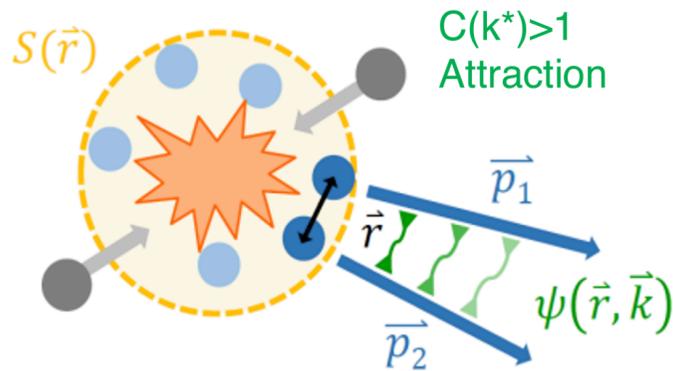
Deuteron production in jets pp collisions



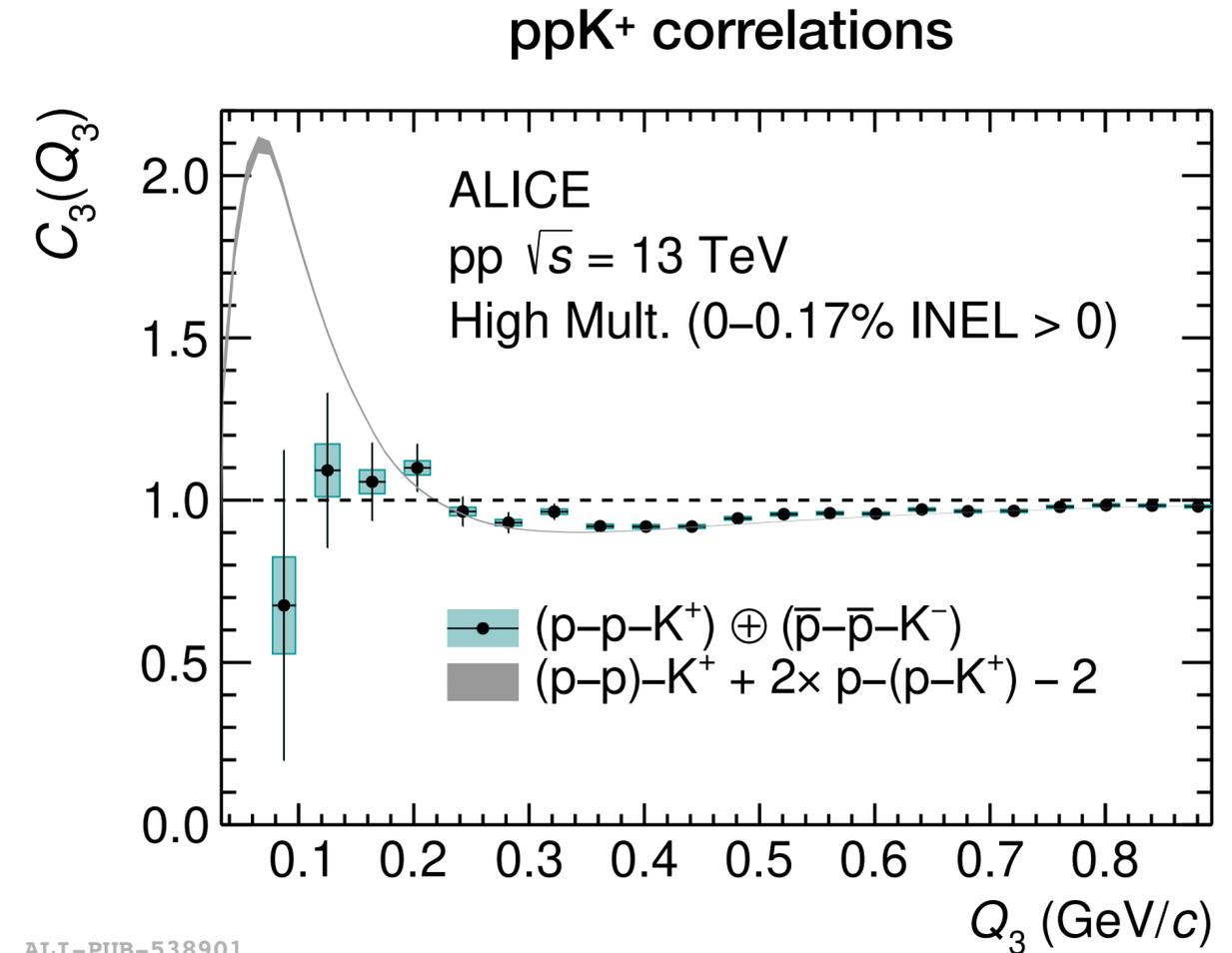
Deuteron production in jets: increase of rate due to higher local density in line with coalescence model expectation

Hadron interactions: 3-particle correlations

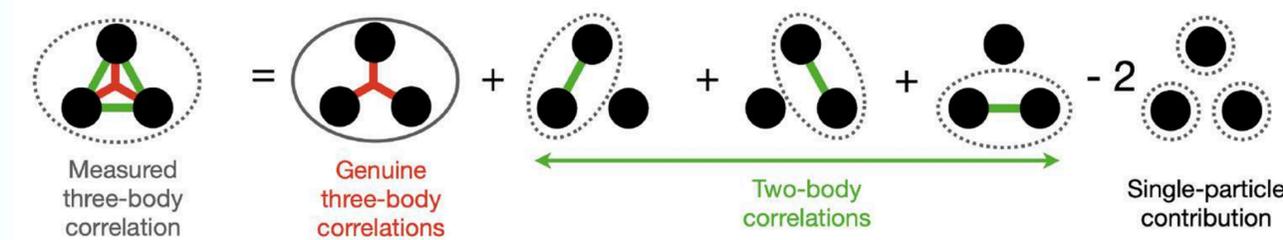
arXiv:2303.13448



ALI-PUB-538905



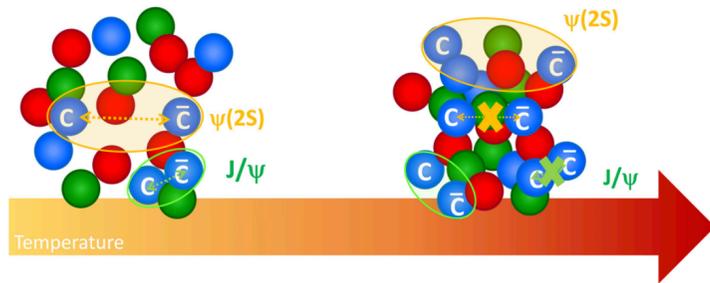
ALI-PUB-538901



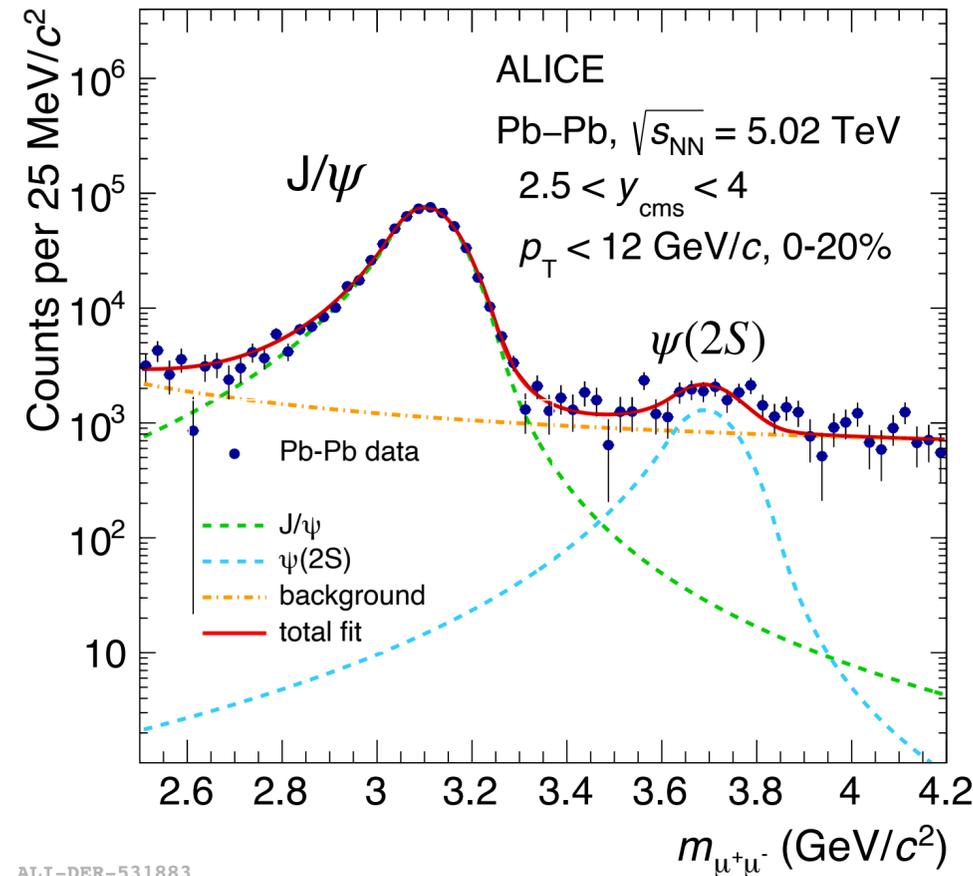
- Use correlations of pairs with similar momentum to measure potential
- First direct study of 3-body potential ppK
- No evidence of true 3-body force; 2-body interactions fully explain correlation signal

Melting and regeneration of charmonia: $\psi(2S)$ vs J/ψ

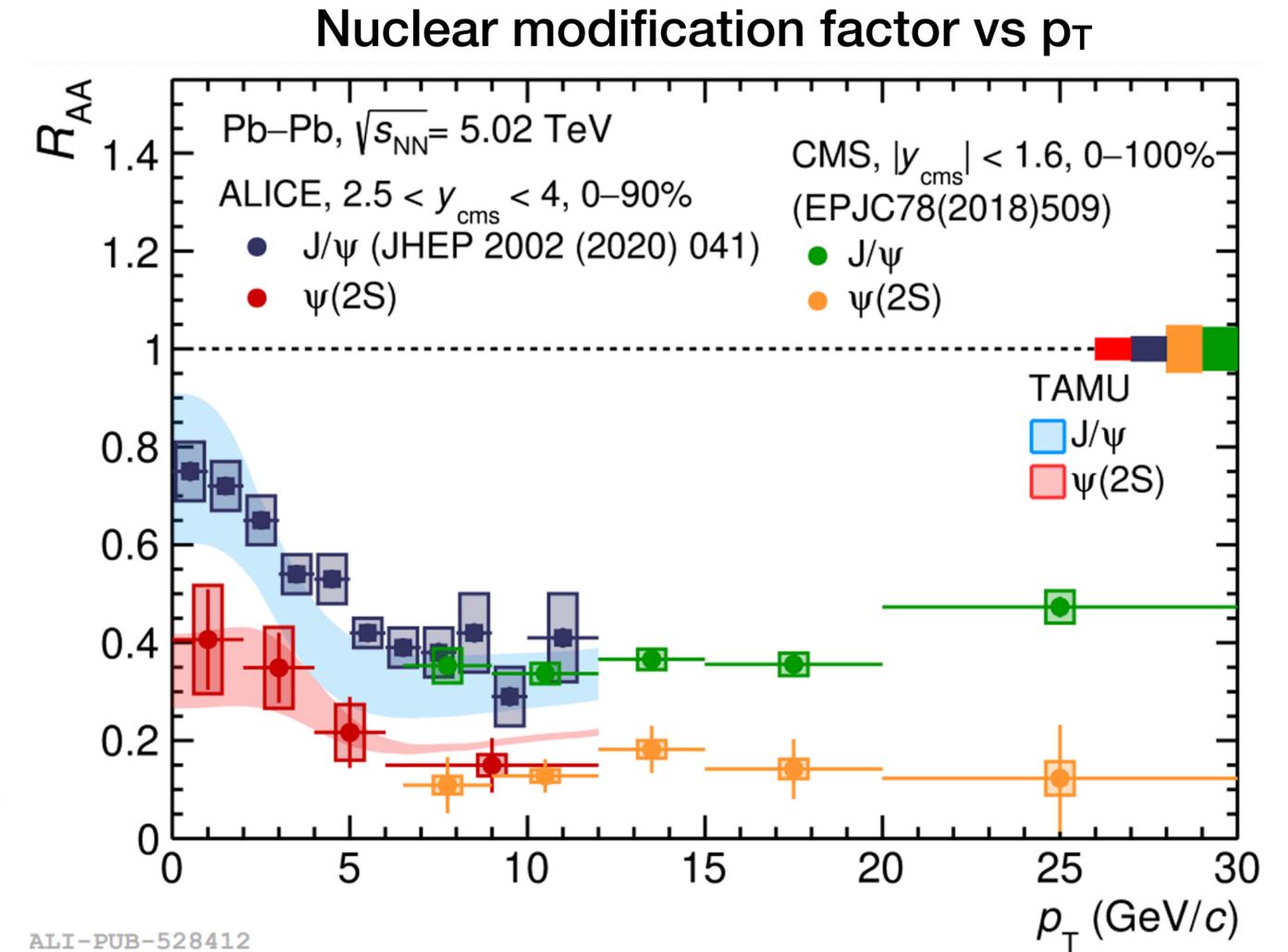
arXiv:2210.08893



Different $c\bar{c}$ bound states:
 $\psi(2S)$ and J/ψ
 different binding energies, sizes



ALI-DER-531883



ALI-PUB-528412

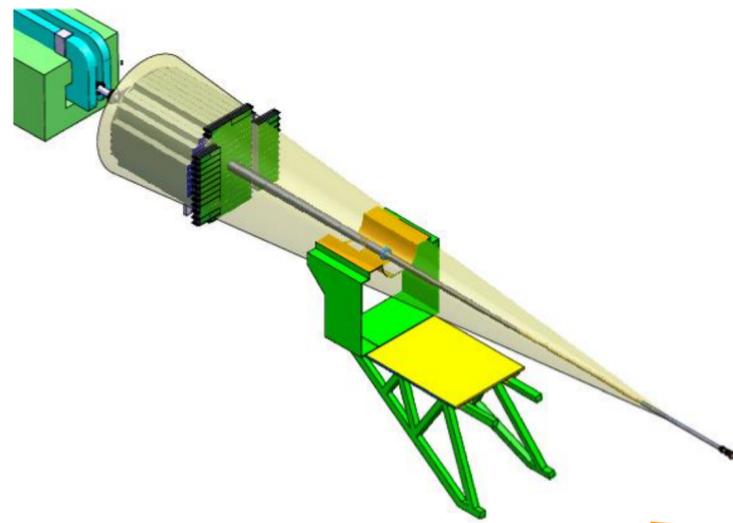
$$R_{AA} = \frac{dN/dp_T|_{AA}}{\langle N_{coll} \rangle dN/dp_T|_{pp}}$$

- High p_T : stronger suppression of $\psi(2S)$ – lower melting temperature
- Low p_T : R_{AA} increases – regeneration similar to J/ψ

ALICE upgrade plans for LS3 and LS4

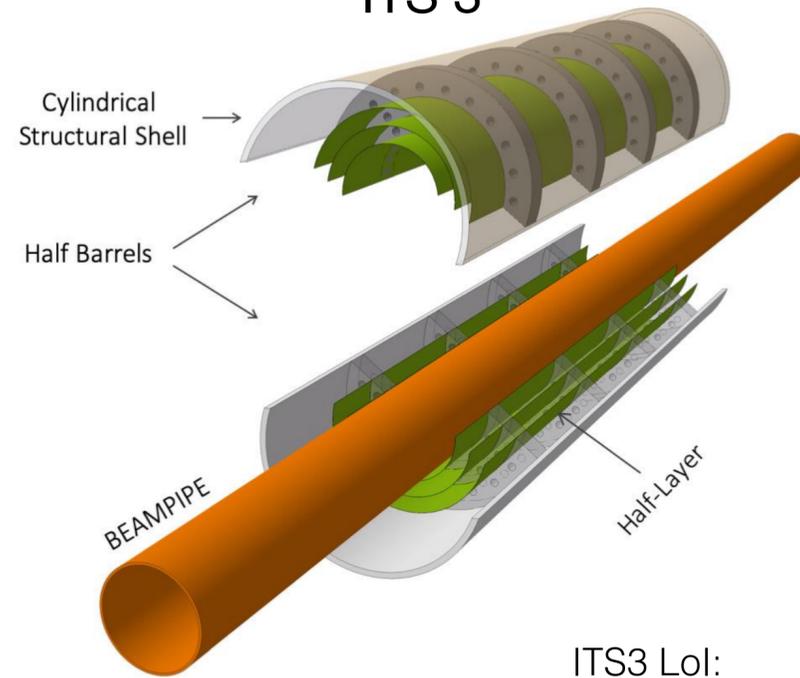
LS3 upgrades

Forward Calorimeter



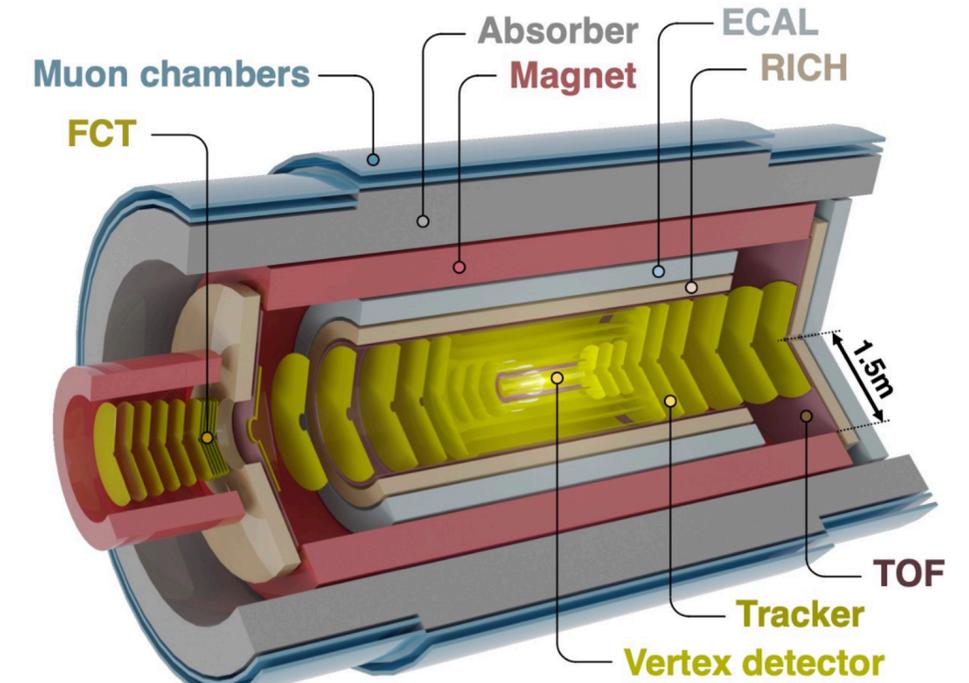
FoCal Lol:
CERN-LHCC-2020-009

ITS 3

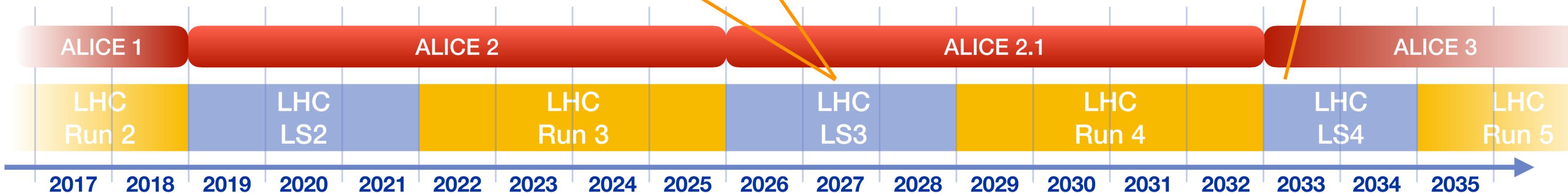


ITS3 Lol:
CERN-LHCC-2019-018

ALICE 3: LS4



ALICE 3 Lol:
CERN-LHCC-2022-009

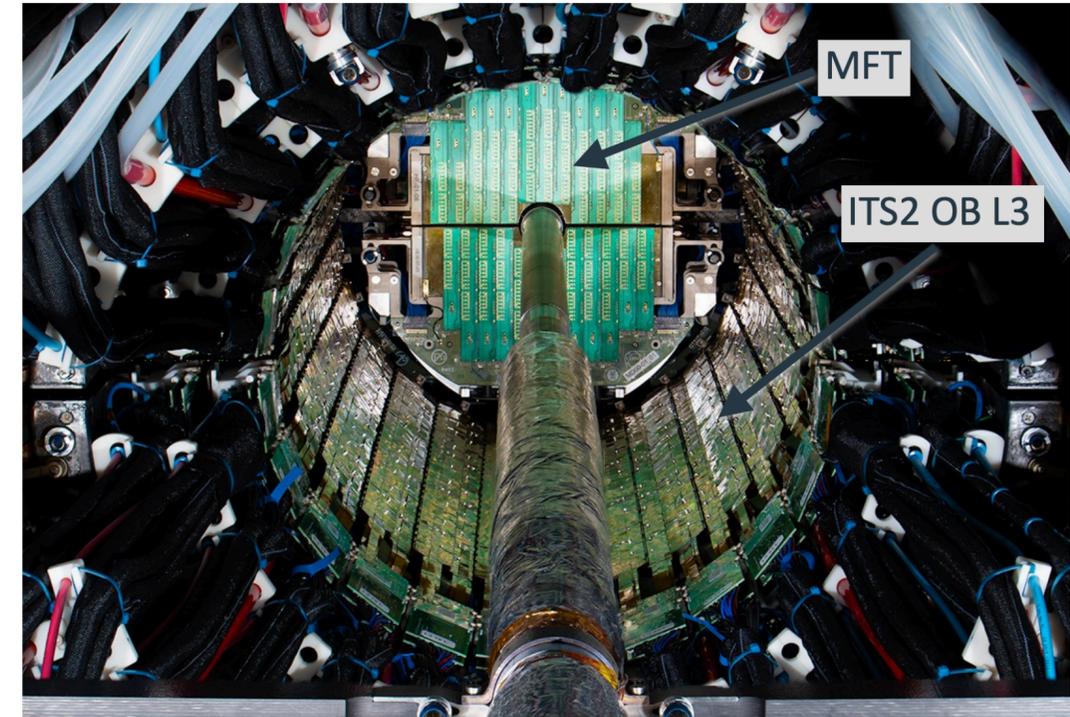


Detector development: silicon pixel sensors

High-energy physics experiments bring together disciplines:

- detector development
- electronics
- computing
- cooling
- mechanics
- super-conducting magnet technology

Upgraded inner tracking system

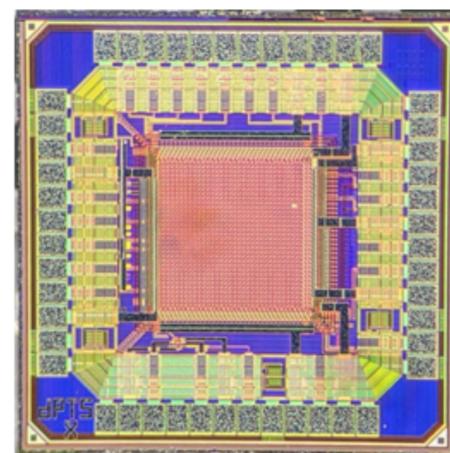
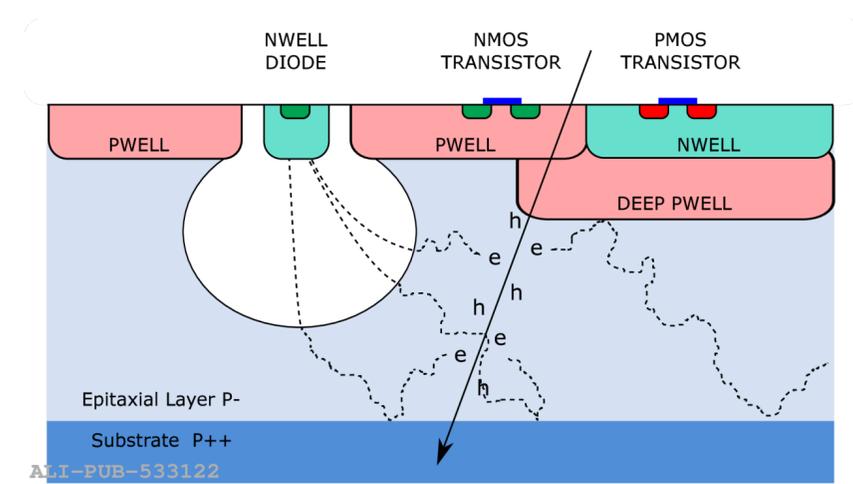


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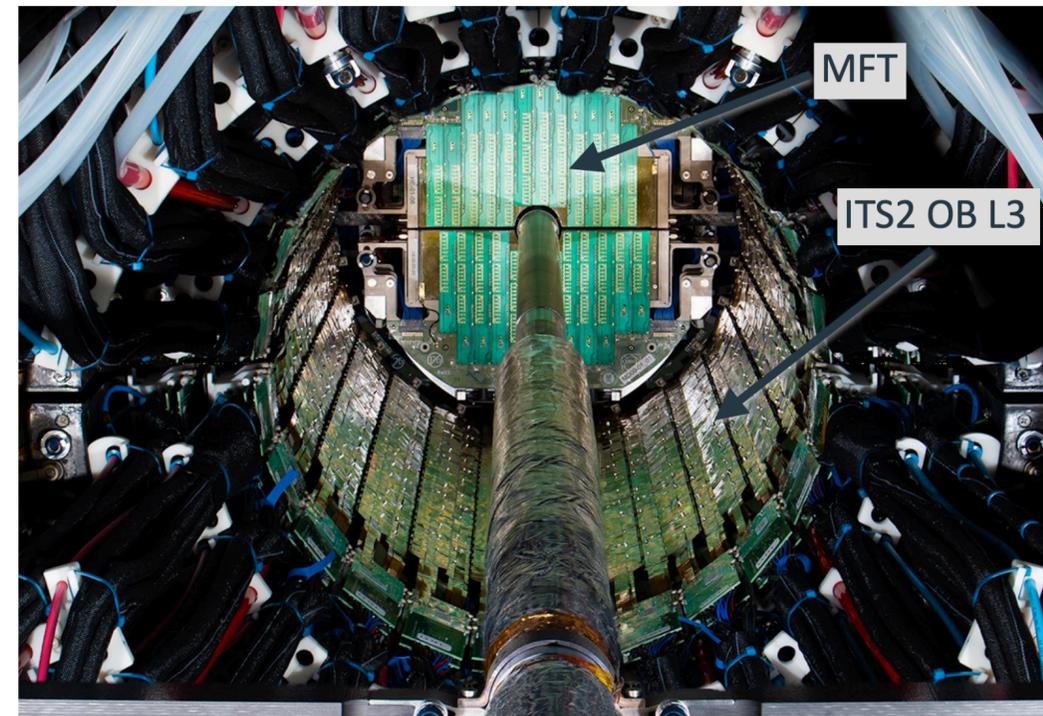
- detector development
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ALICE has spearheaded the development and adaptation of monolithic active pixel sensors



DPTS test paper [arXiv:2212.08621](https://arxiv.org/abs/2212.08621)

Upgraded inner tracking system

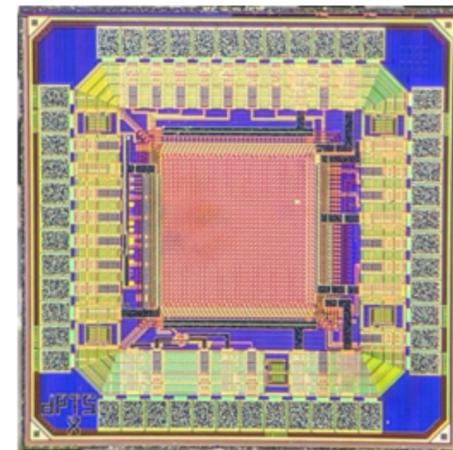
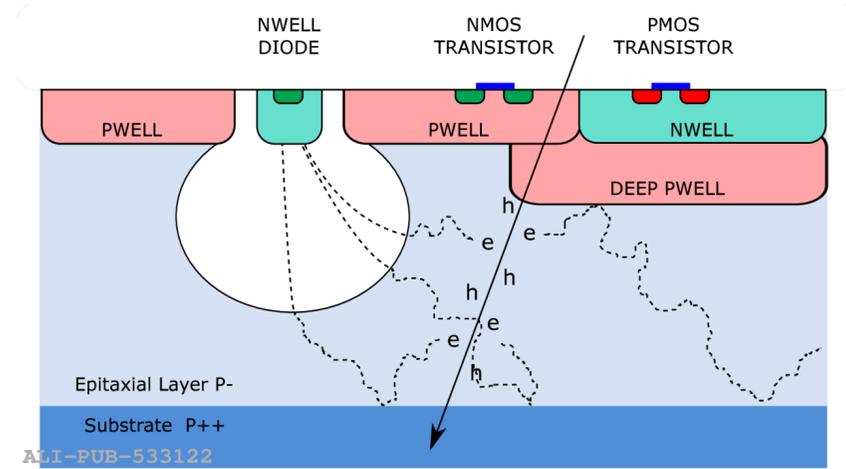


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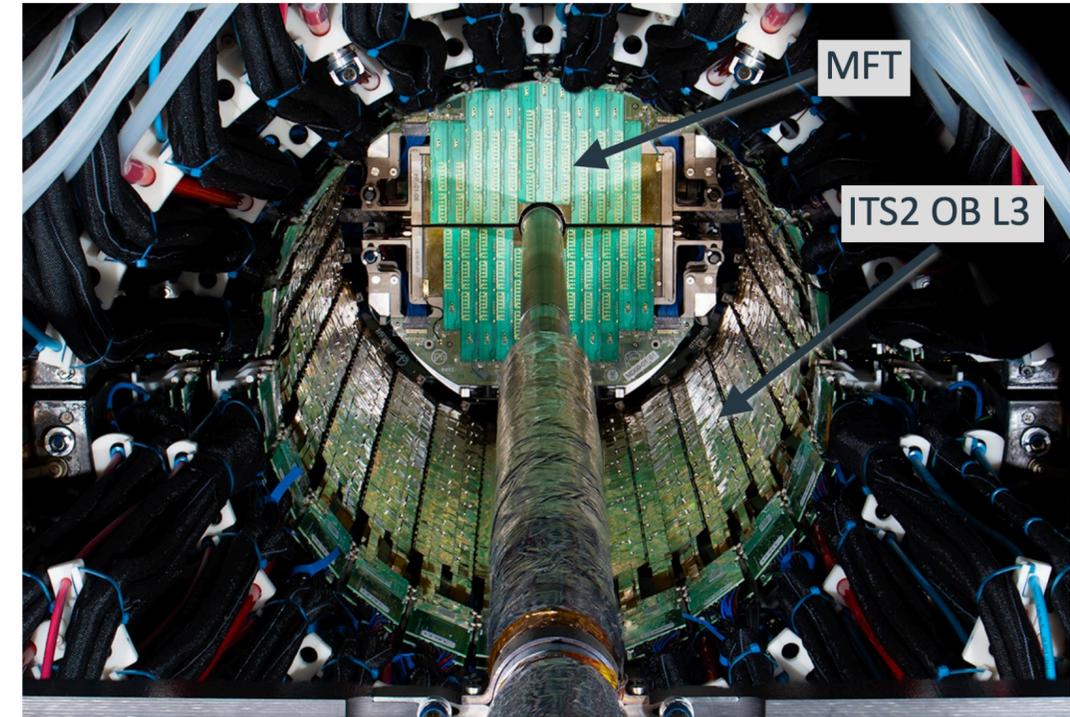
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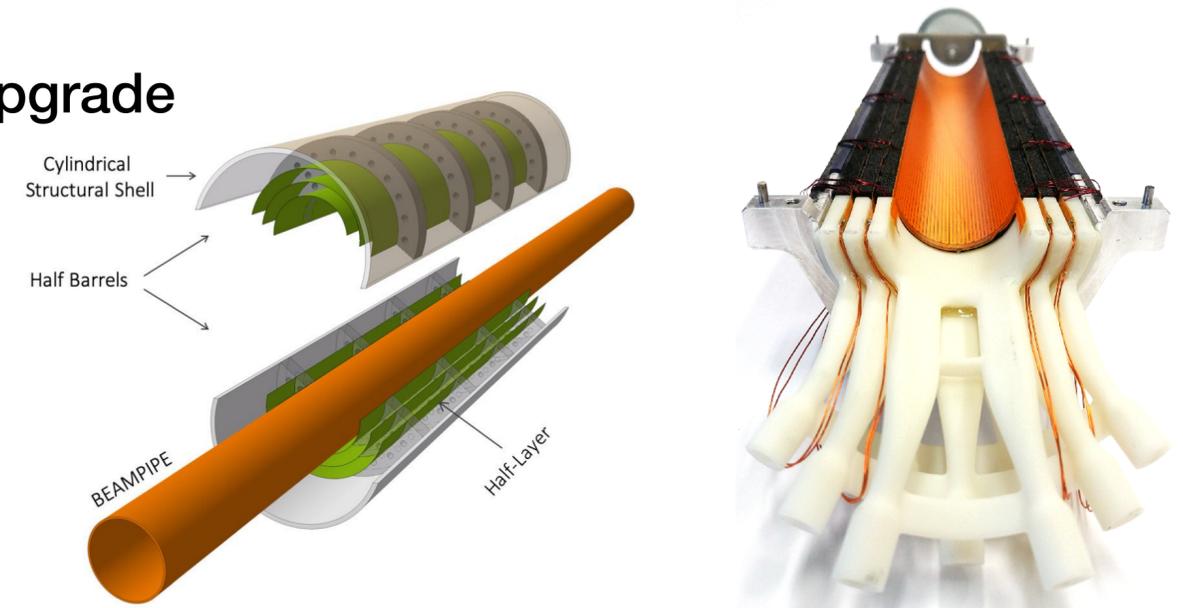


DPTS test paper arXiv:2212.08621

Upgraded inner tracking system



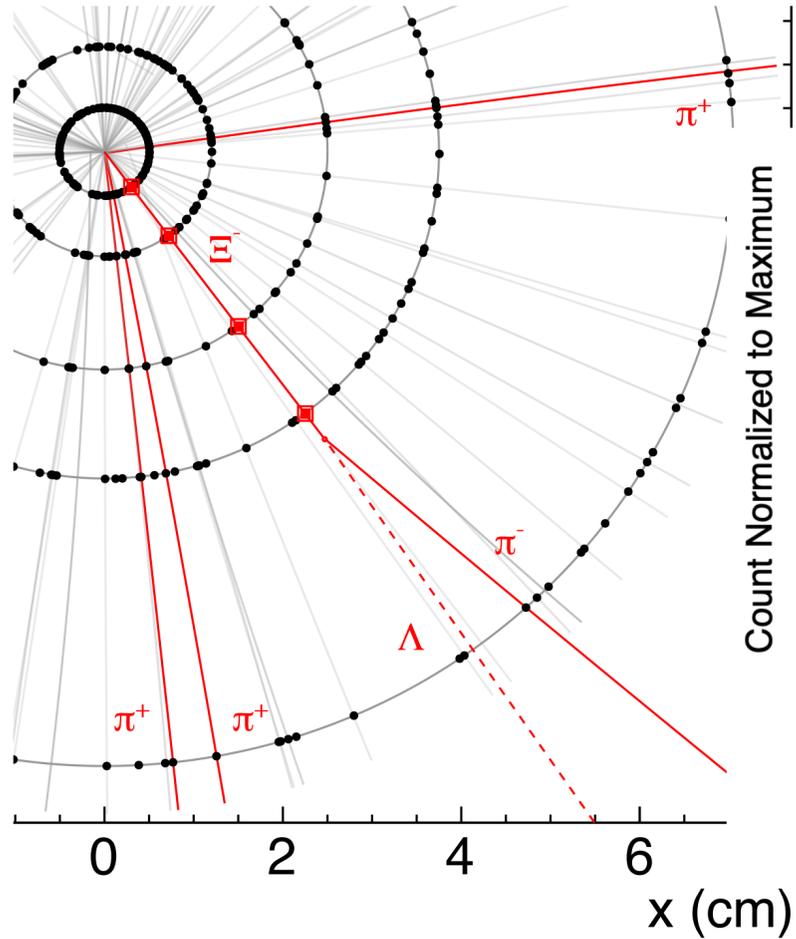
ITS3 upgrade



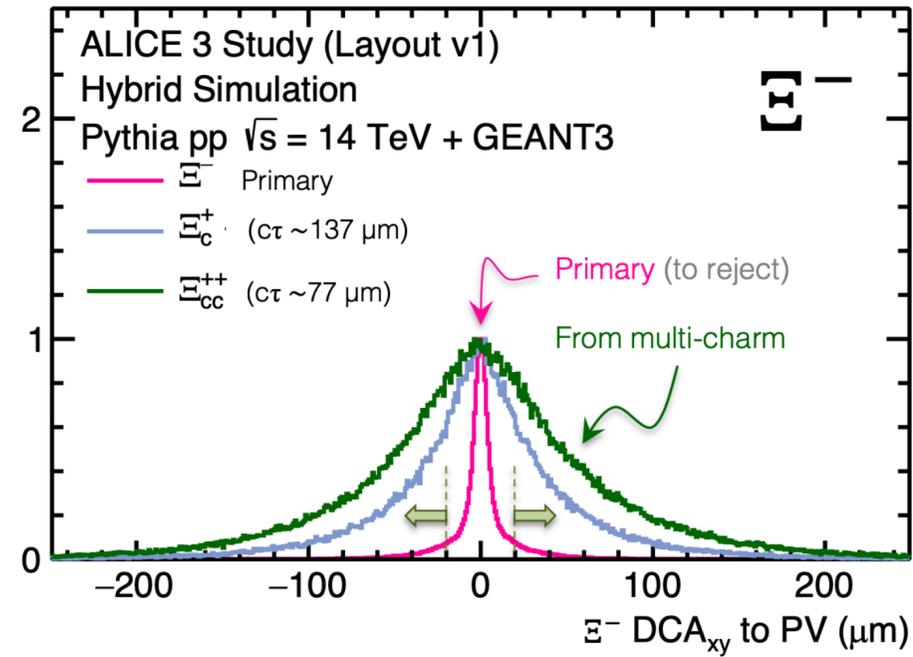
Ultralight inner layers: improved pointing resolution

Multi-charm baryon detection

New technique: strangeness tracking



Impact parameter of Ξ

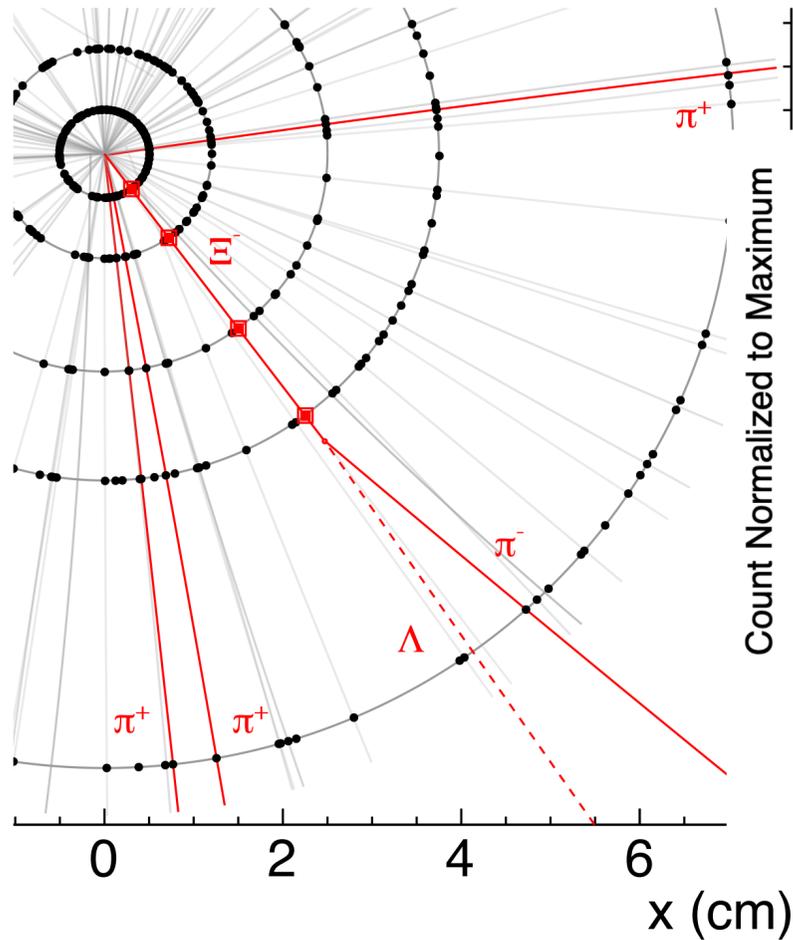


Xi baryon provides high selectivity

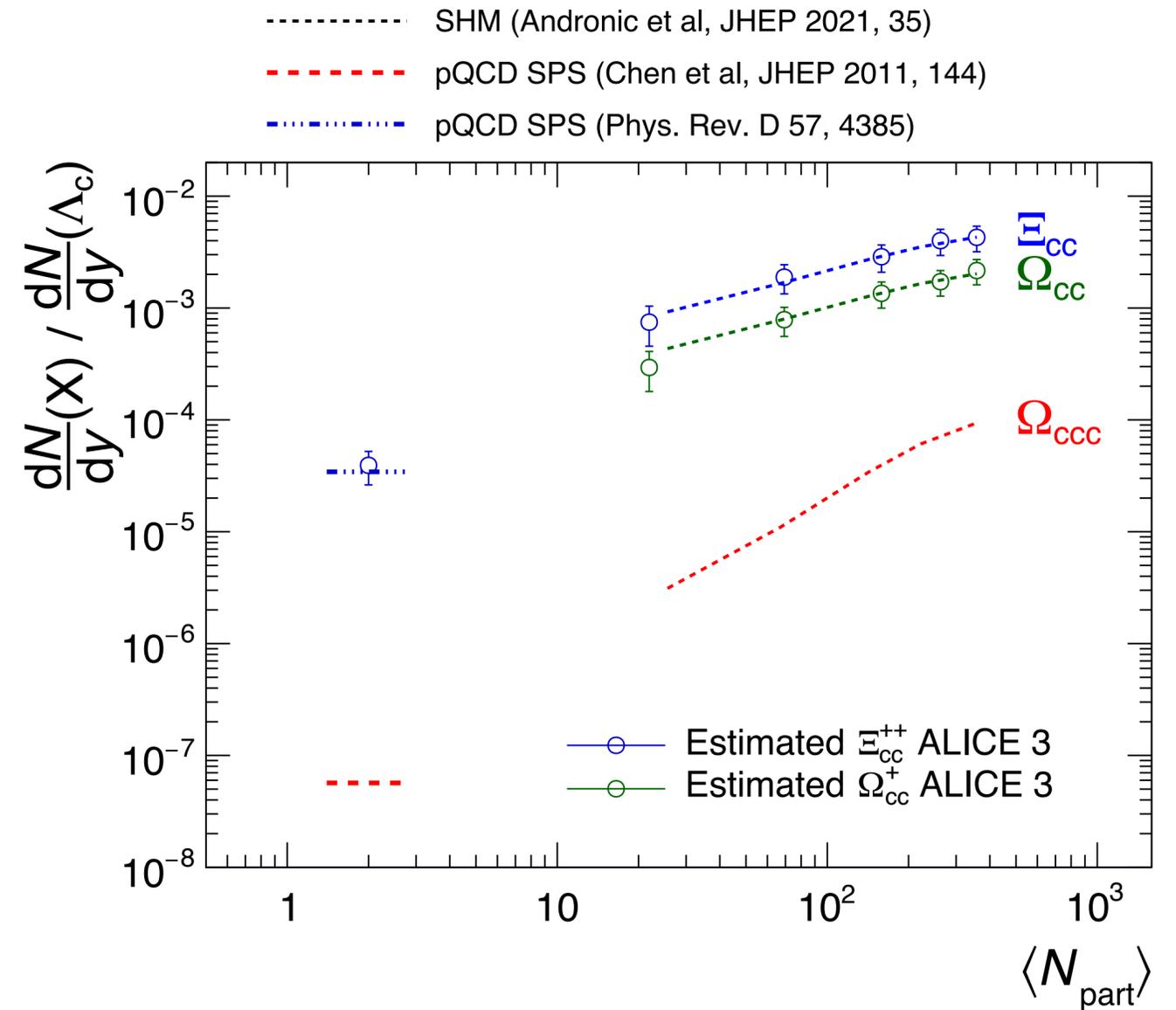
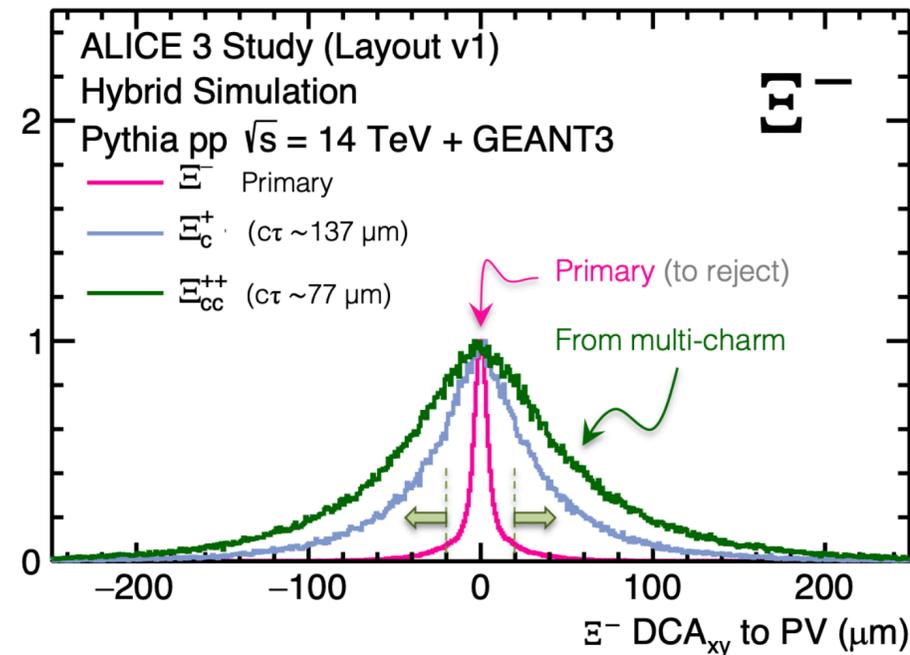


Multi-charm baryon detection

New technique: strangeness tracking



Impact parameter of Ξ



Multi-charm baryon provides high selectivity

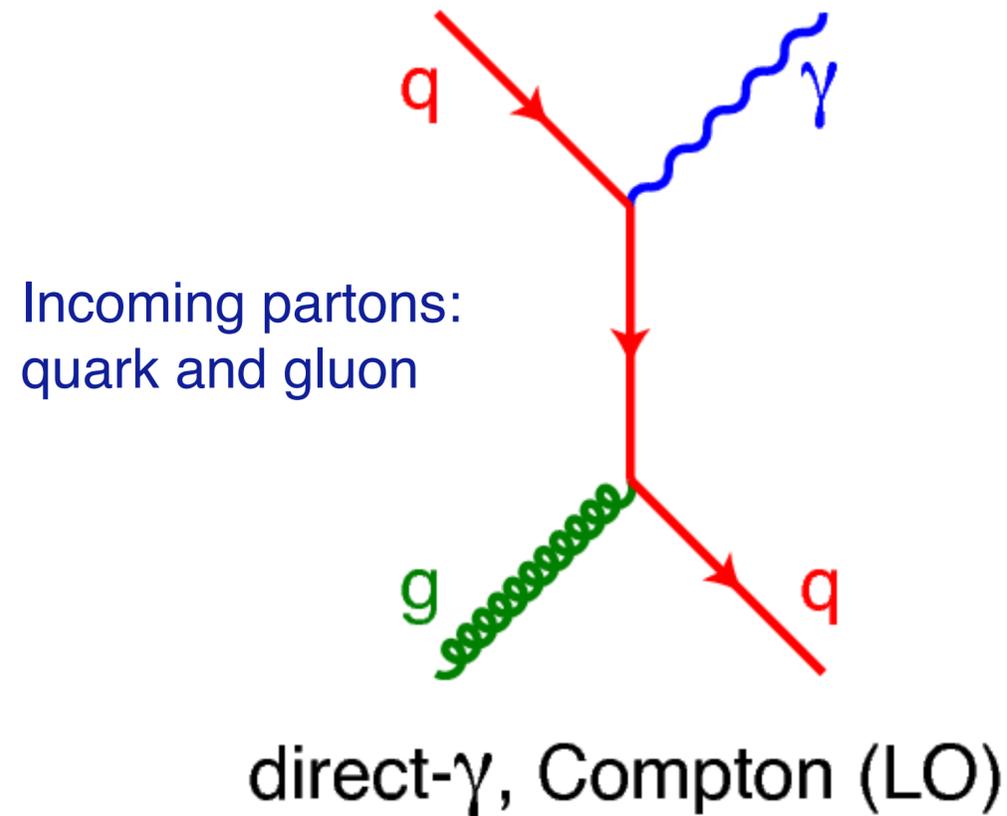


Large enhancements: unique sensitivity to thermalisation and hadronisation dynamics

Unique access in Pb-Pb collisions with ALICE 3

Probing the gluon density in a hadron collider

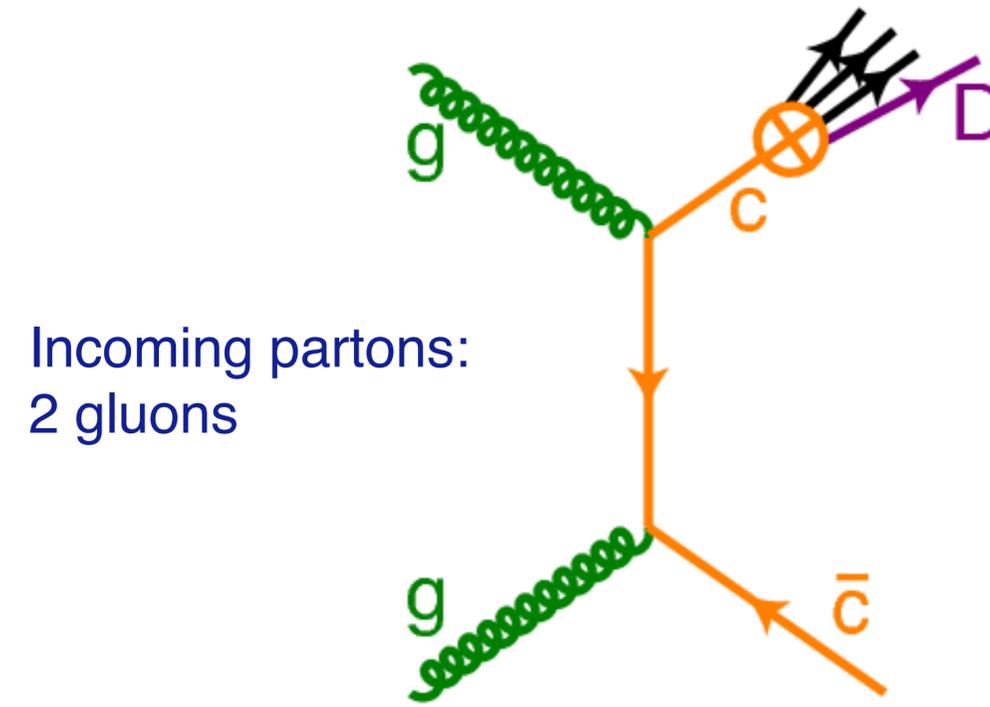
Direct photon production



Sensitive to **gluons at LO**

Photon momentum directly related to incoming partons

Charm production

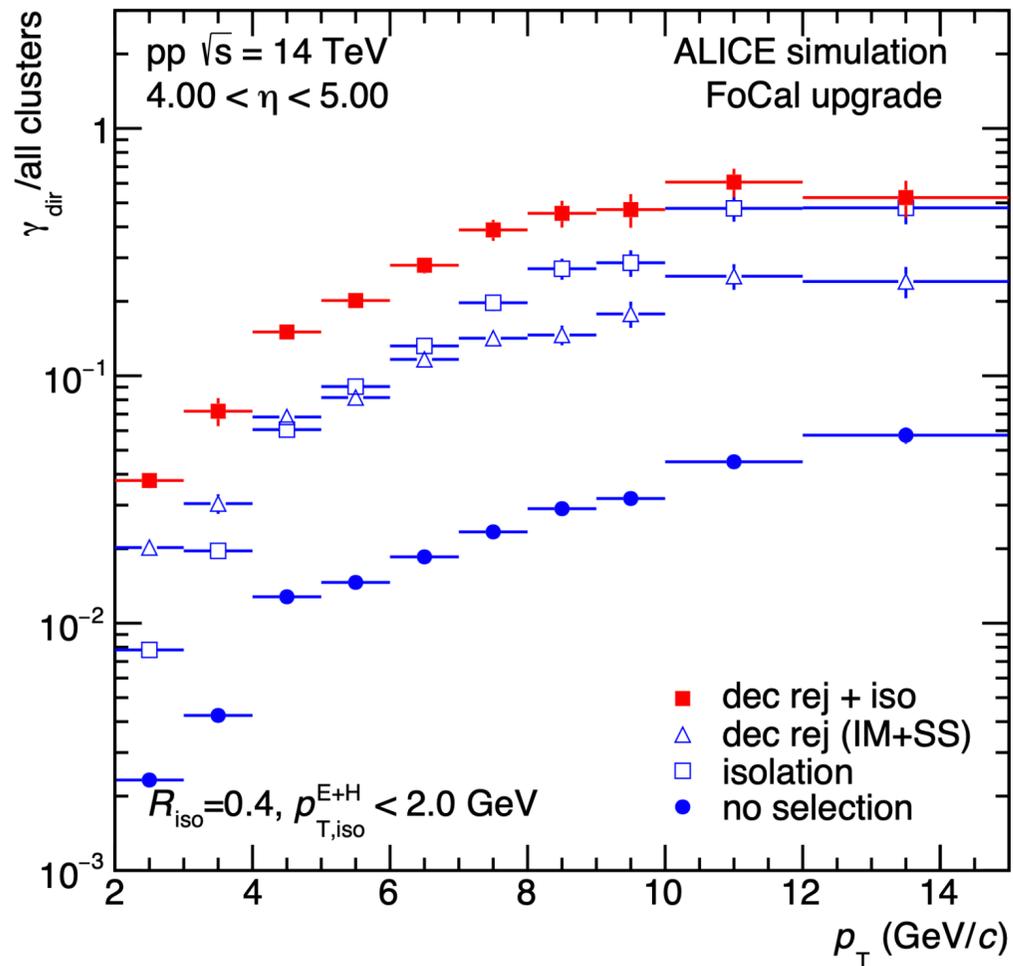


Heavy hadron:
also directly sensitive
but fragmentation reduces
kinematic constraint

More processes contribute, e.g. gluon splitting

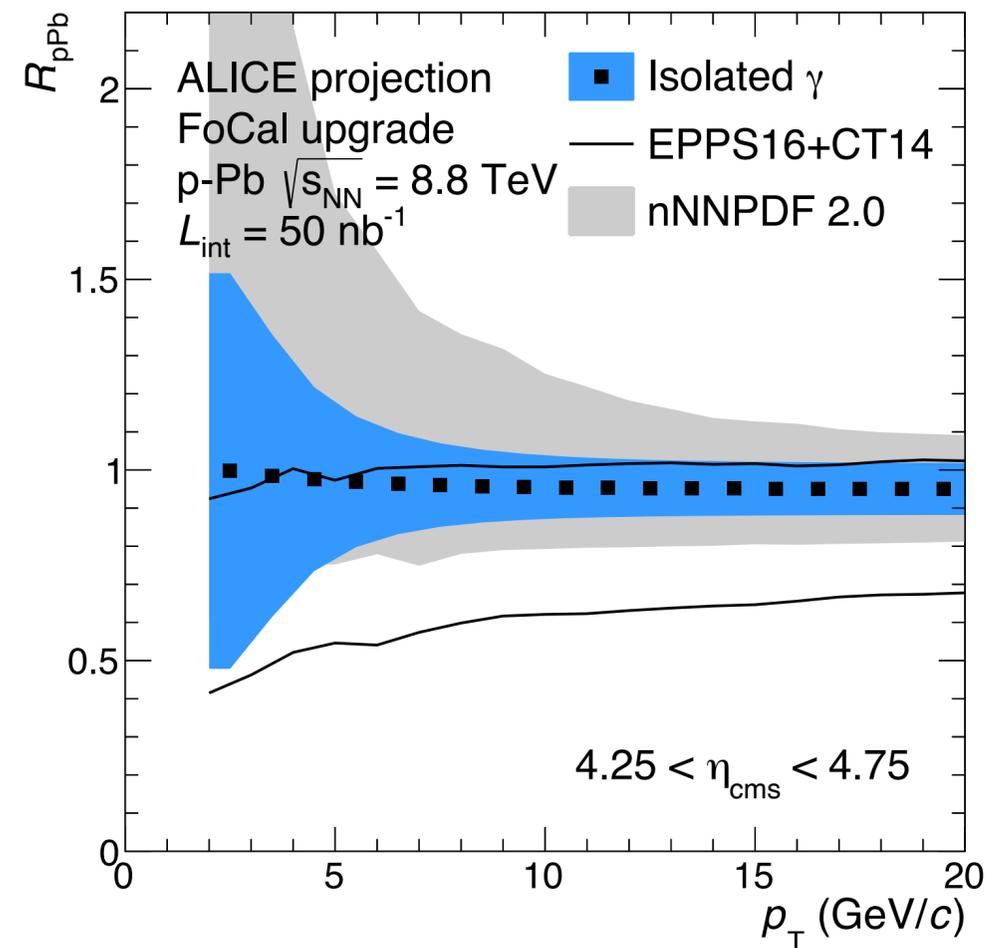
Forward photons with FoCal

Signal photon fraction



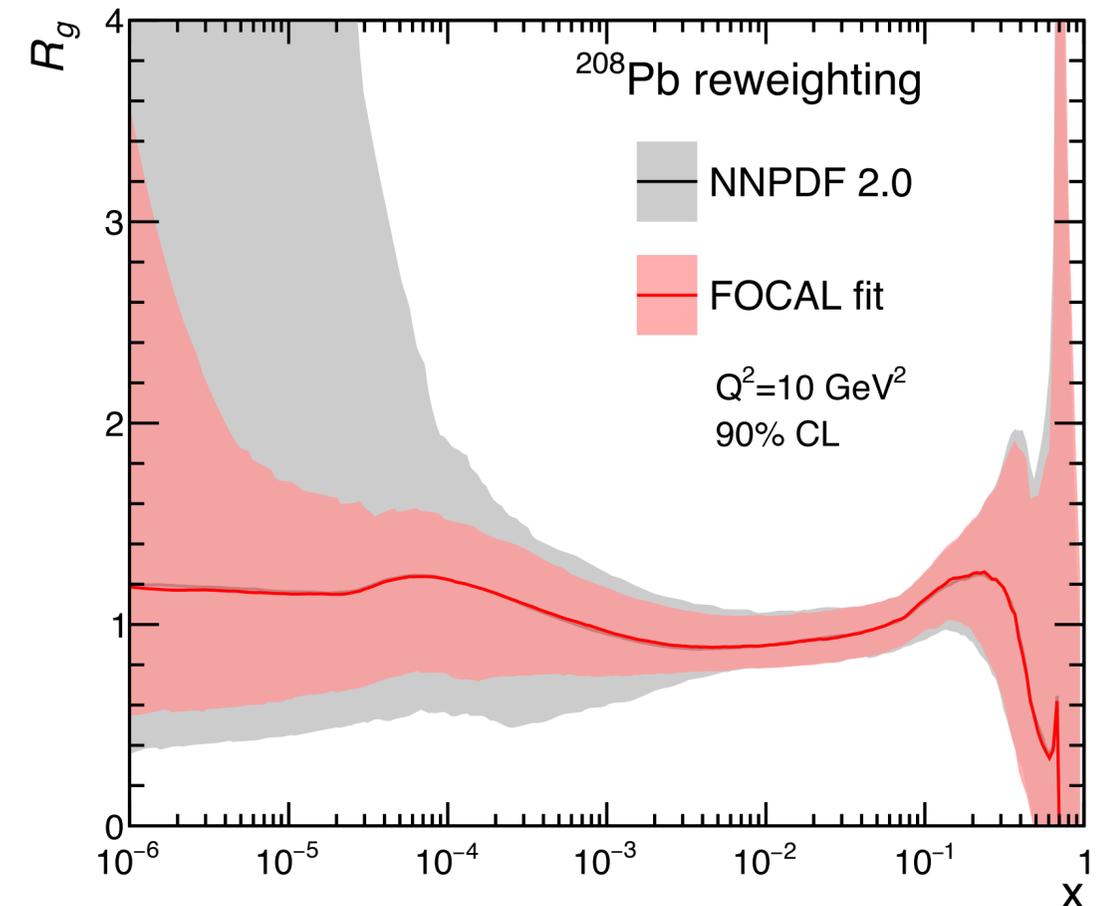
High granularity to reject decay background

Projected photon uncertainties



High precision direct photon measurement down to low p_T

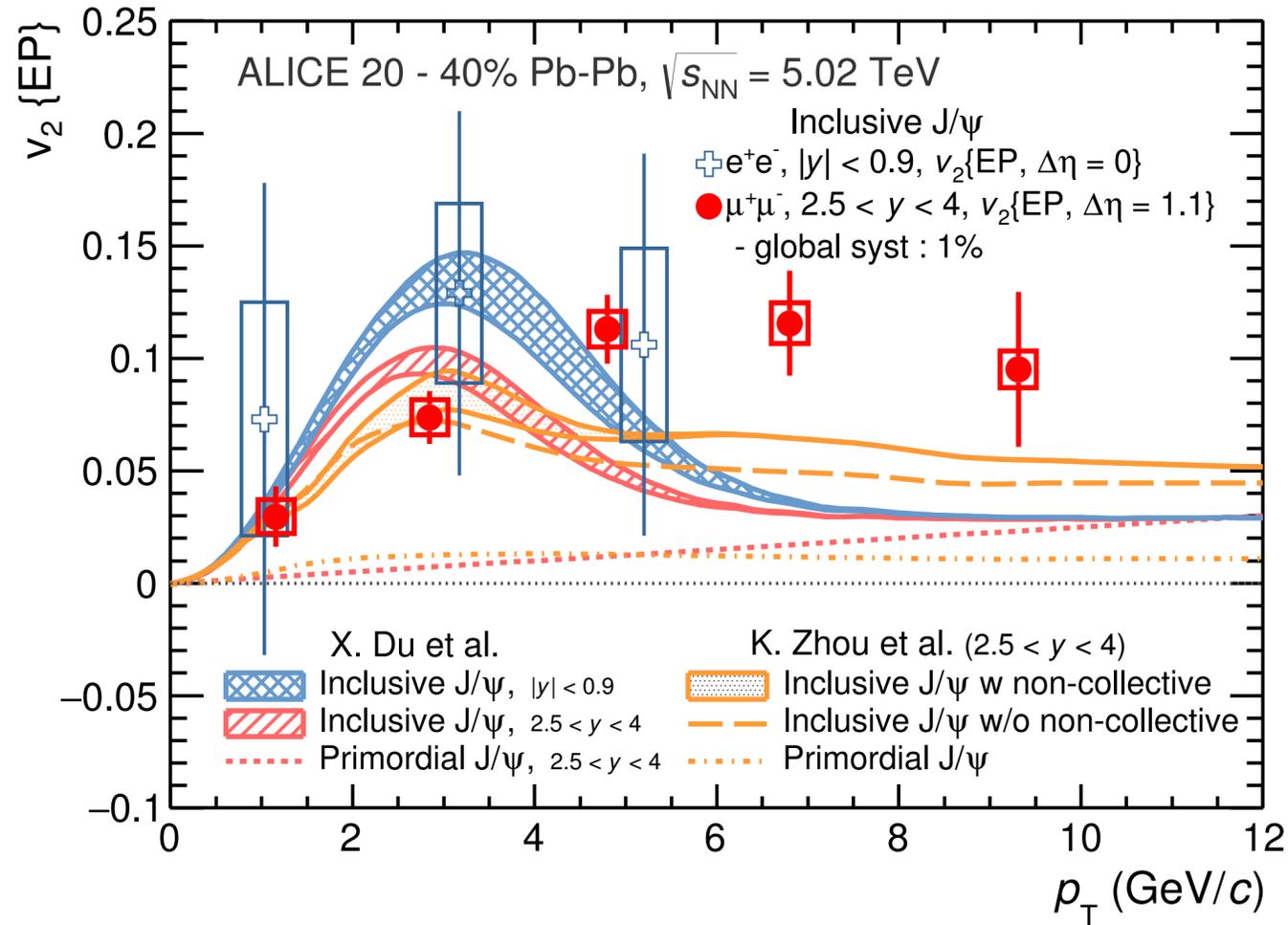
Projected PDF uncertainties



Constrain gluon density in nuclei over a broad range:
 $x \sim 10^{-5} - 10^{-2}$ at small Q^2

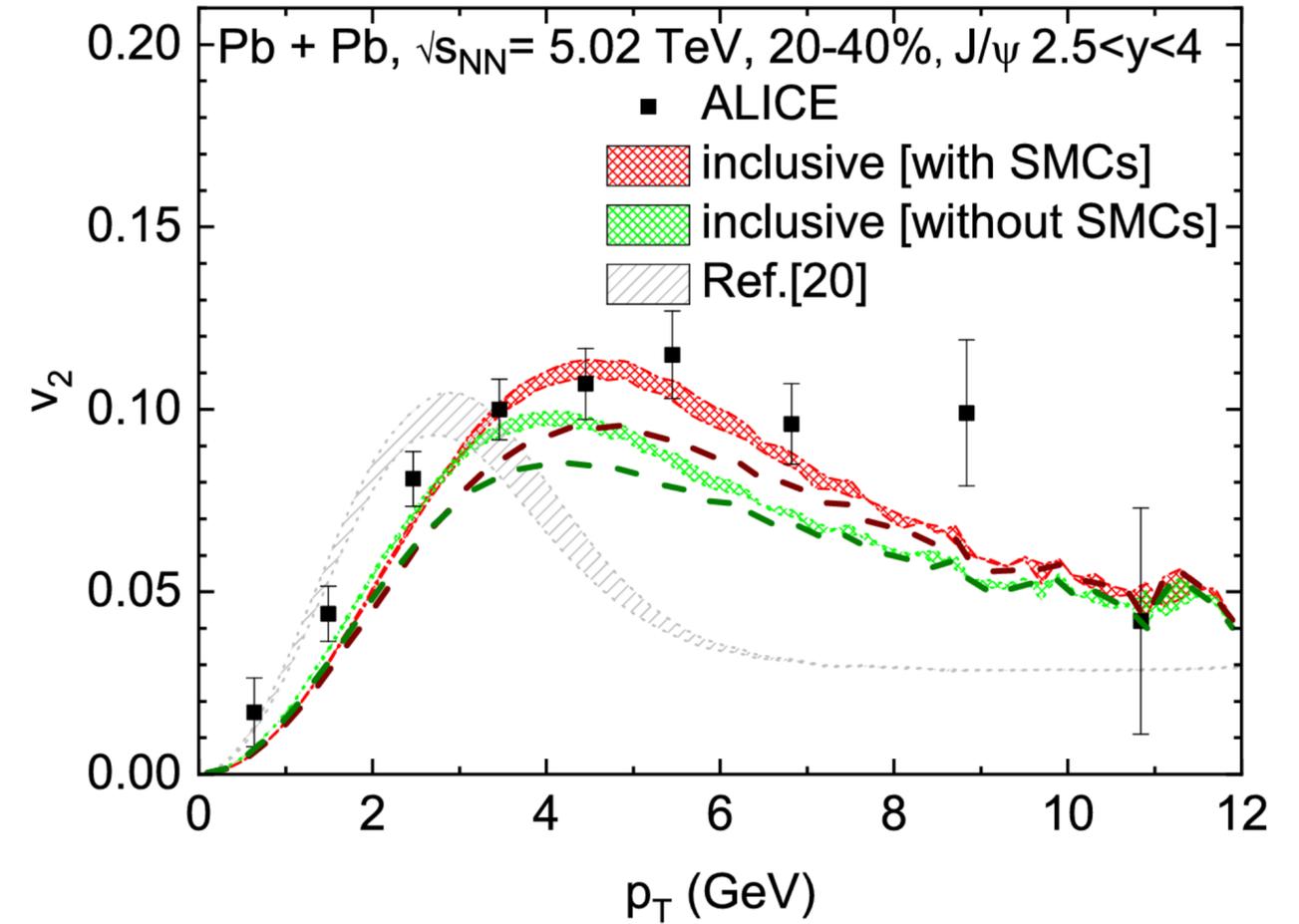
Elliptic flow of quarkonia

J/ψ v_2



PRL 119, 242301

J/ψ v_2 : new model calculation



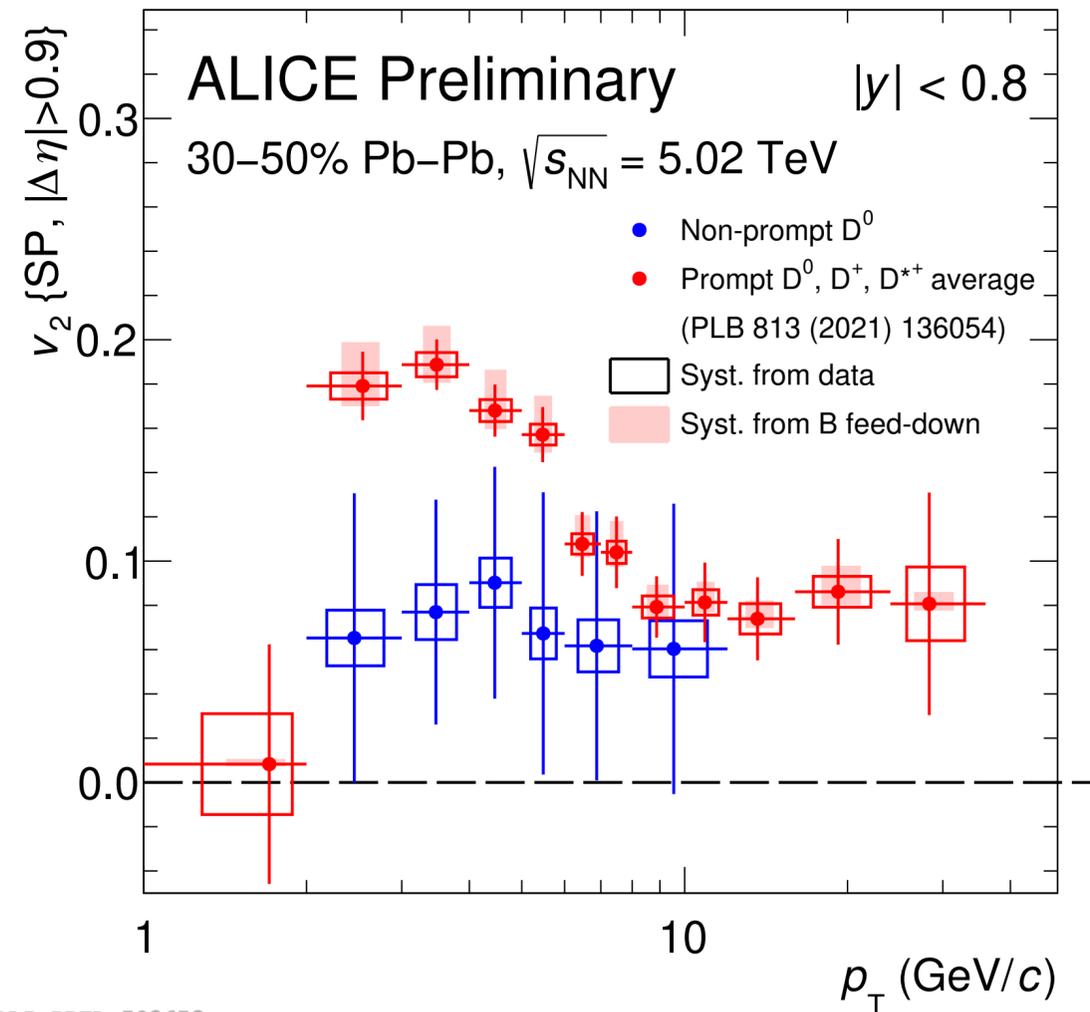
Min He and Ralf Rapp, PRL 128 (2022) 16, 162301

J/ψ melting is not directional; expect no or very little v_2
 Recombination sensitive to v_2 of QGP
 —> Expect non-zero v_2 at low p_T

Updated model: v_2 extends to larger p_T

Interactions of beauty quarks with the QGP: v_2

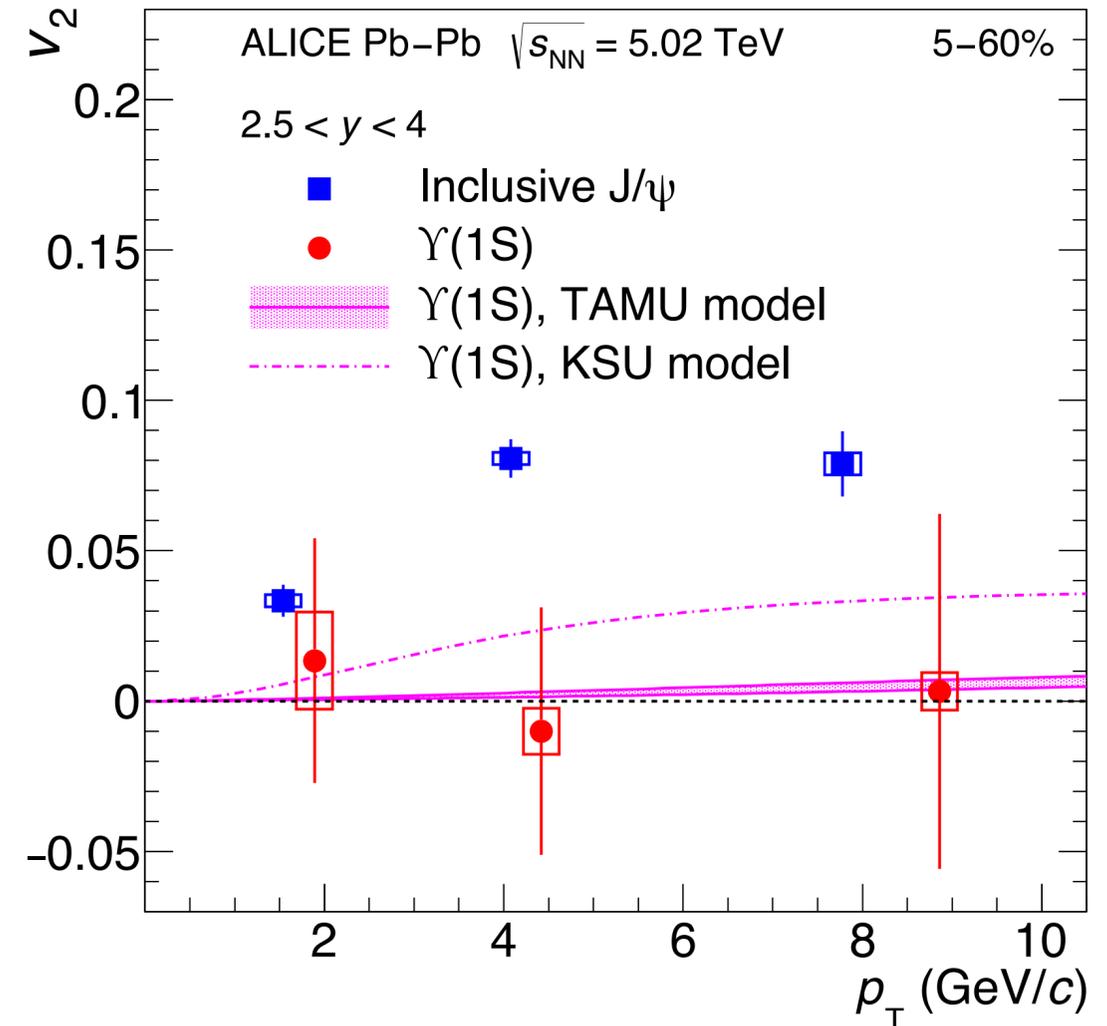
Prompt and non-prompt (from B decay) D^0 v_2



ALI-PREL-502672

Open HF: smaller v_2 for beauty than charm

J/ψ and Υ v_2

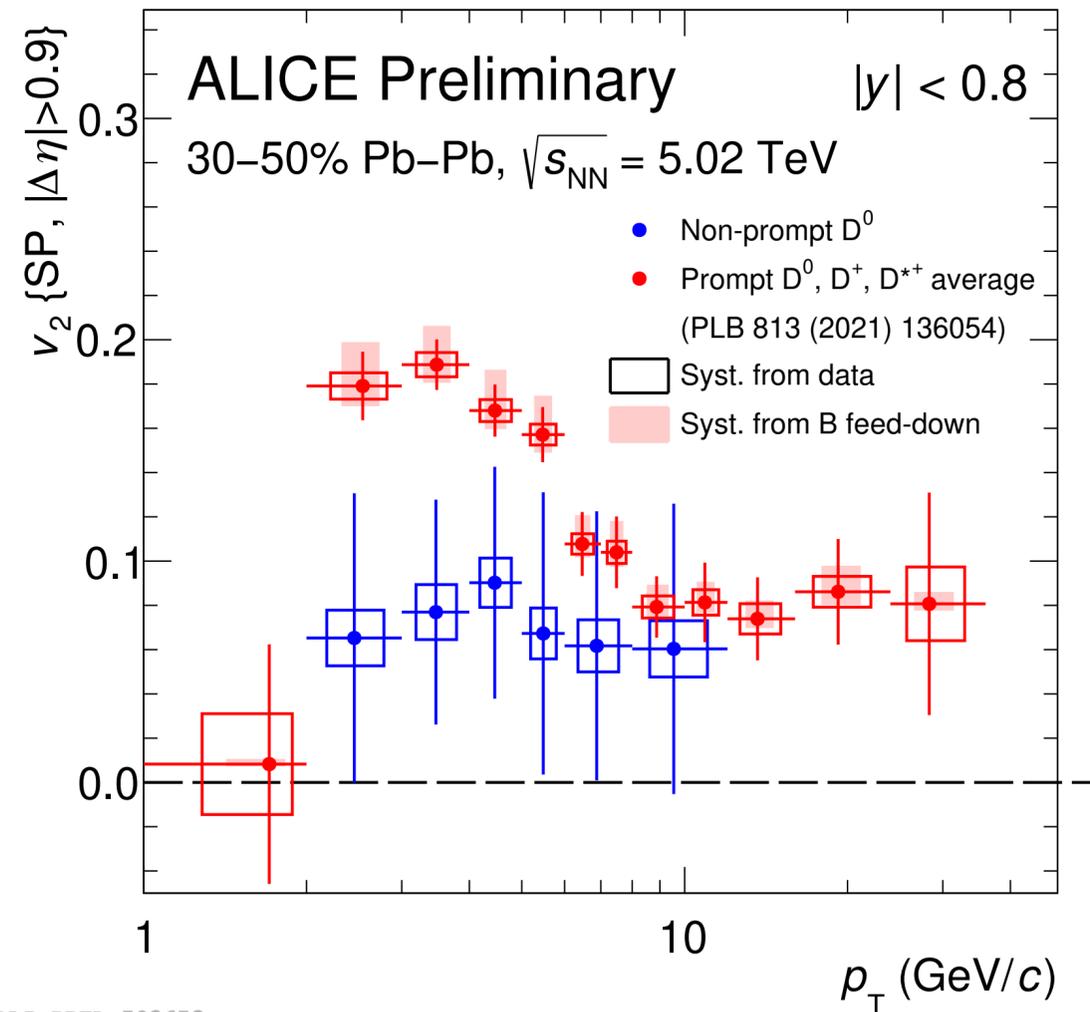


Hidden HF: no beauty v_2 ?

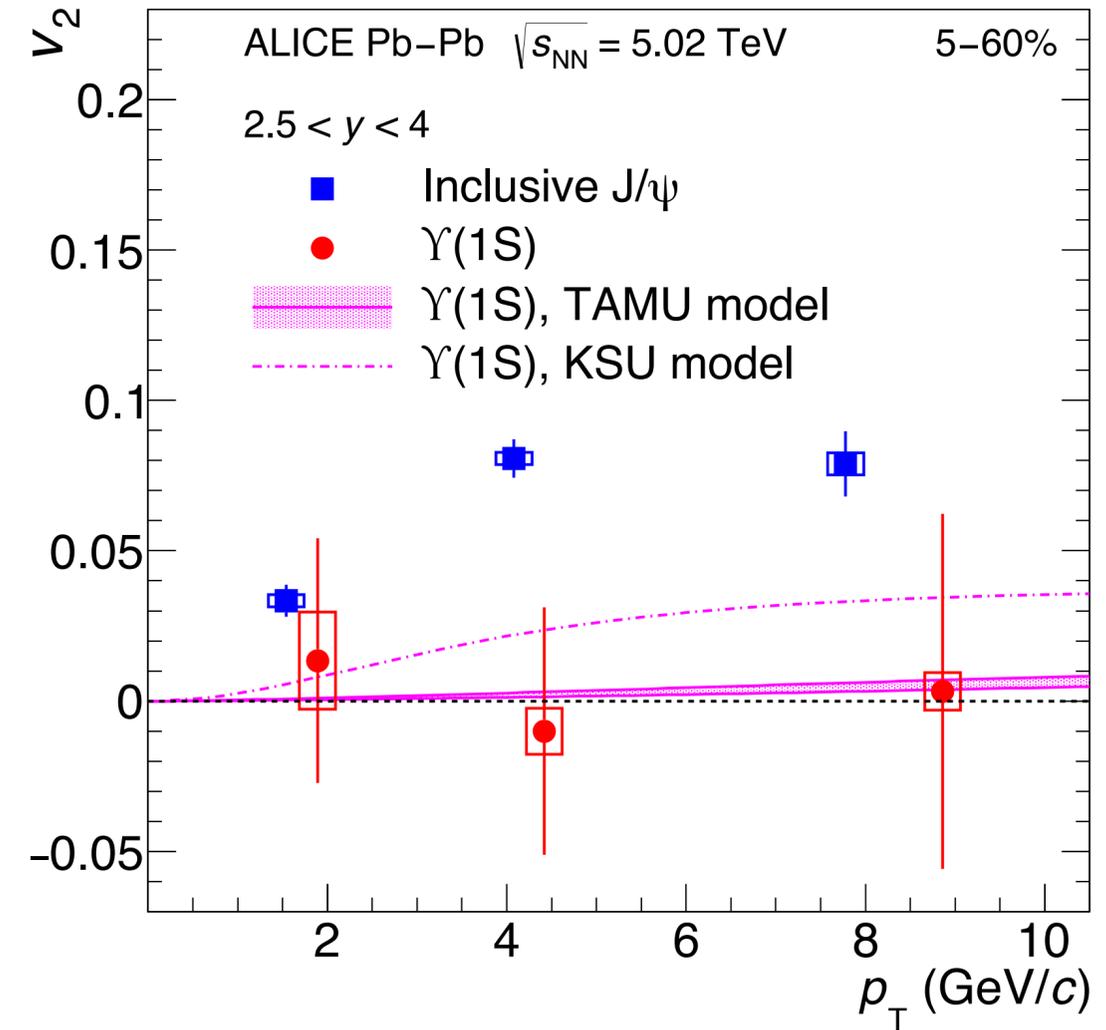
Suggests smaller interaction cross sections for beauty than charm — longer thermalisation time

Interactions of beauty quarks with the QGP: v_2

Prompt and non-prompt (from B decay) D^0 v_2



J/ψ and Υ v_2



ALI-PREL-502672

Open HF: smaller v_2 for beauty than charm

Hidden HF: no beauty v_2 ?

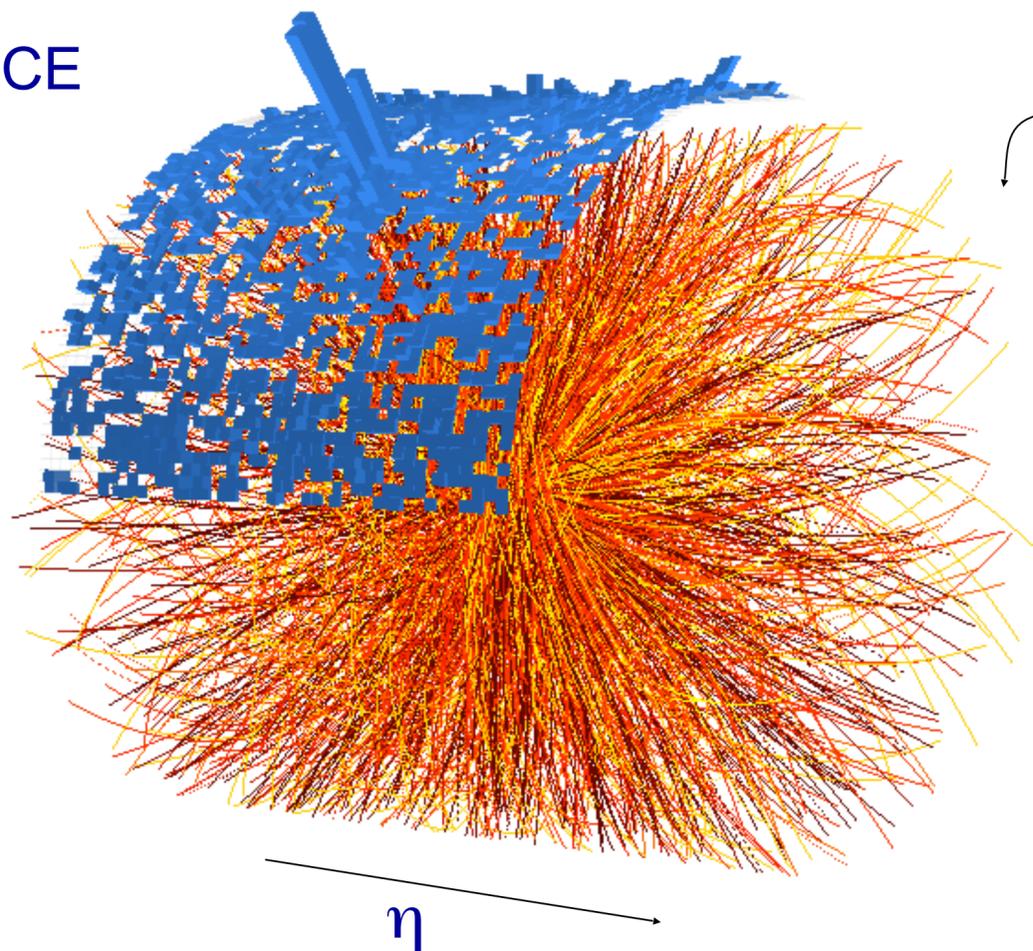
Suggests smaller interaction cross sections for beauty than charm — longer thermalisation time

Light quark v_2 may contribute to open beauty v_2 , via recombination

Uncertainties large — improve precision to conclude

Probing the QGP with jets at LHC

ALICE



ϕ

E_T (GeV)

100
80
60
40
20
0

-4

-2

0

2

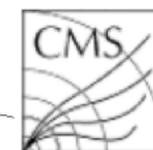
4

-2

2

0

ϕ



CMS Experiment at LHC, CERN
Data recorded: Sun Nov 14 19:31:39 2010 CEST
Run/Event: 151076 / 1328520
Lumi section: 249

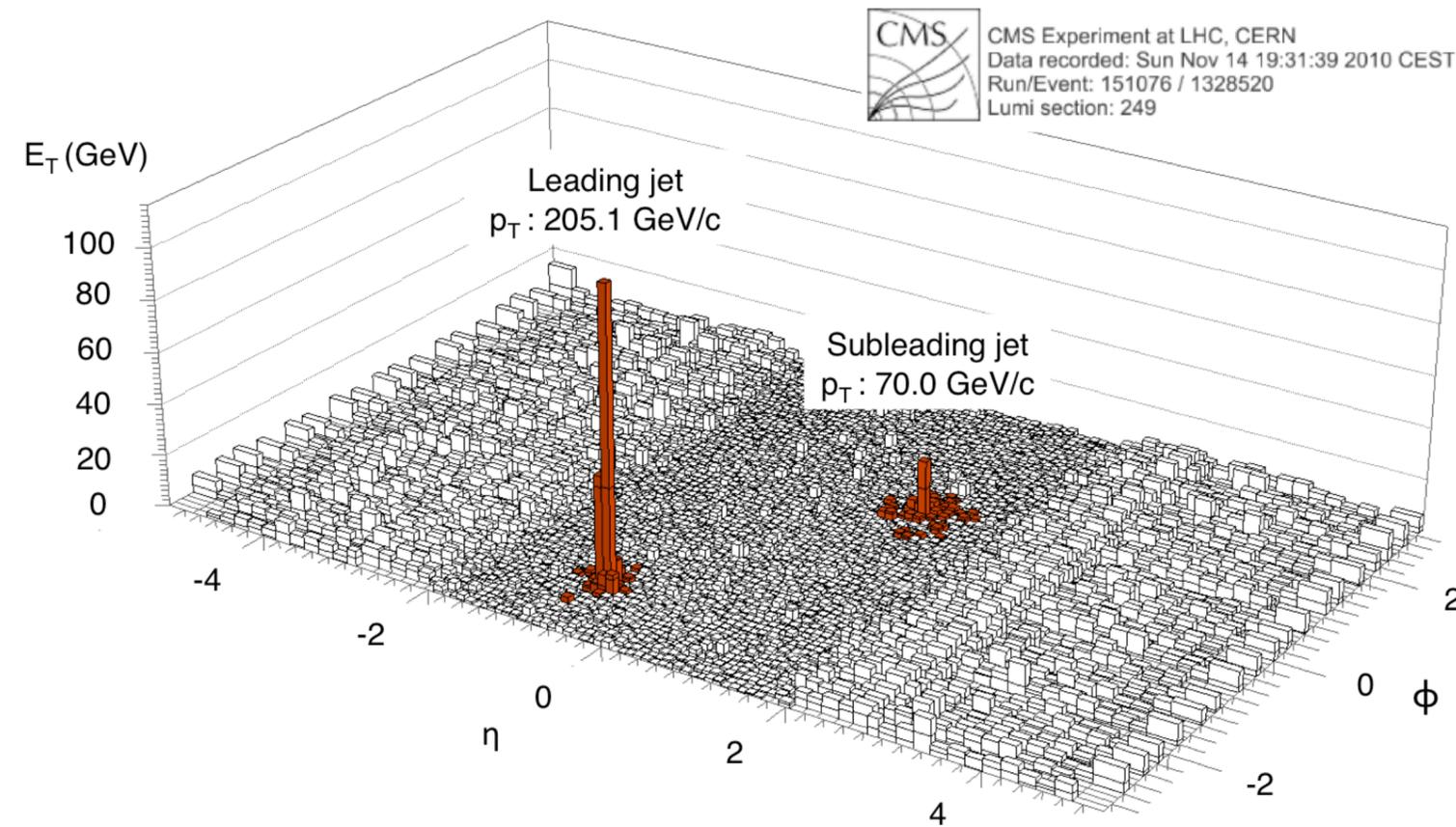
Leading jet
 p_T : 205.1 GeV/c

Subleading jet
 p_T : 70.0 GeV/c

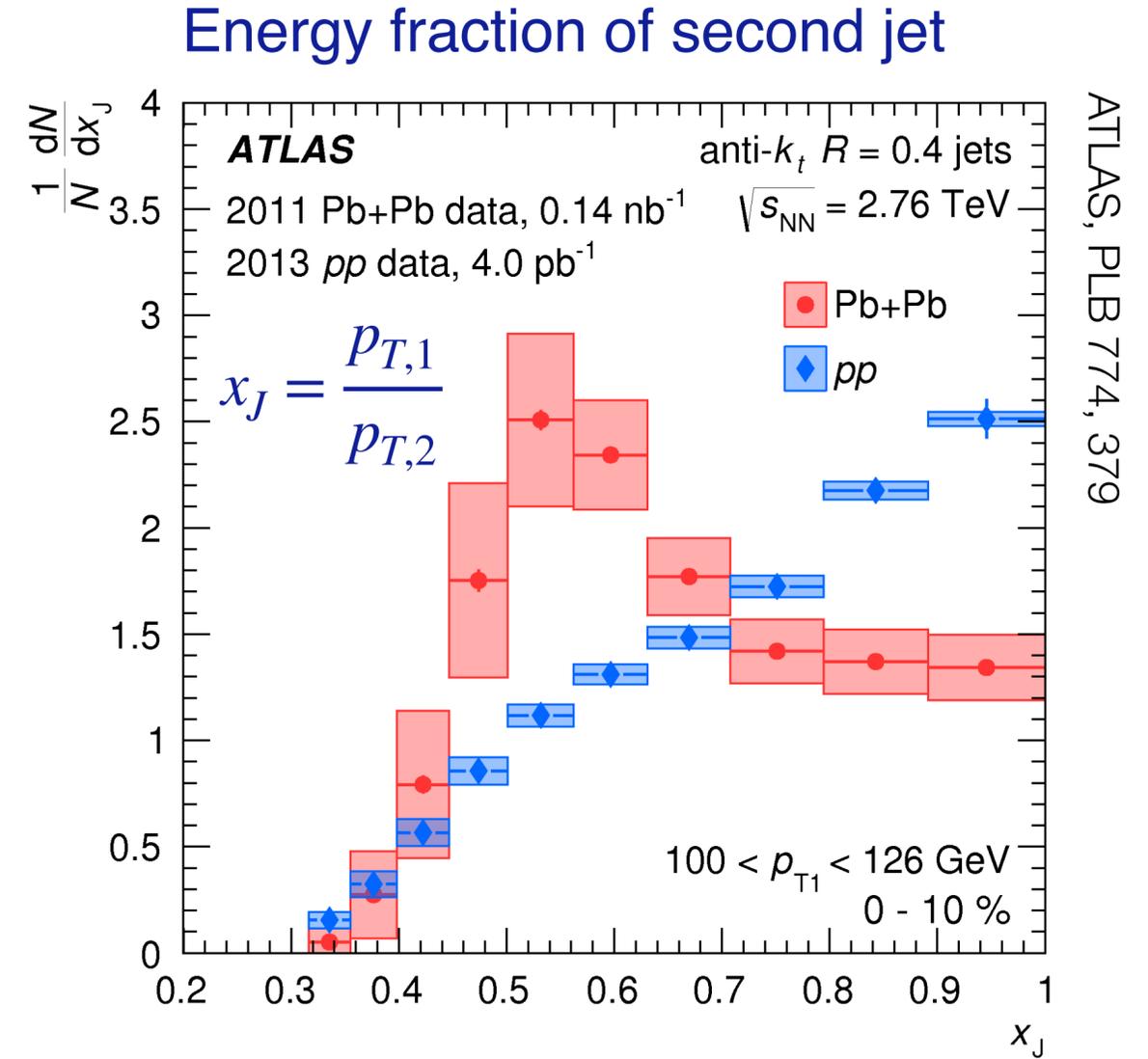
Very clear signals at high p_T : jets stand out above uncorrelated 'soft' background

Interactions with QGP: energy loss, high- p_T suppression

Energy loss: di-jet asymmetry



Single event: p_T not balanced!

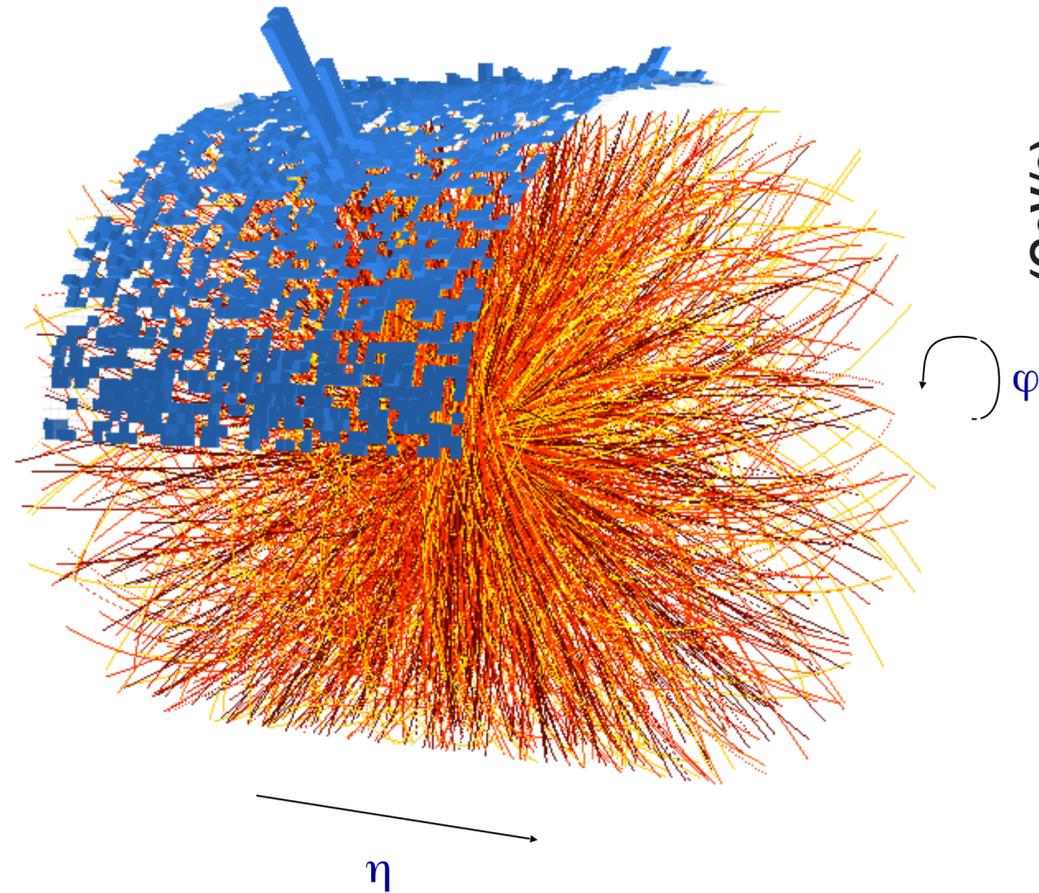


pp: peak at 1 — balanced jets
 PbPb: shift towards lower values

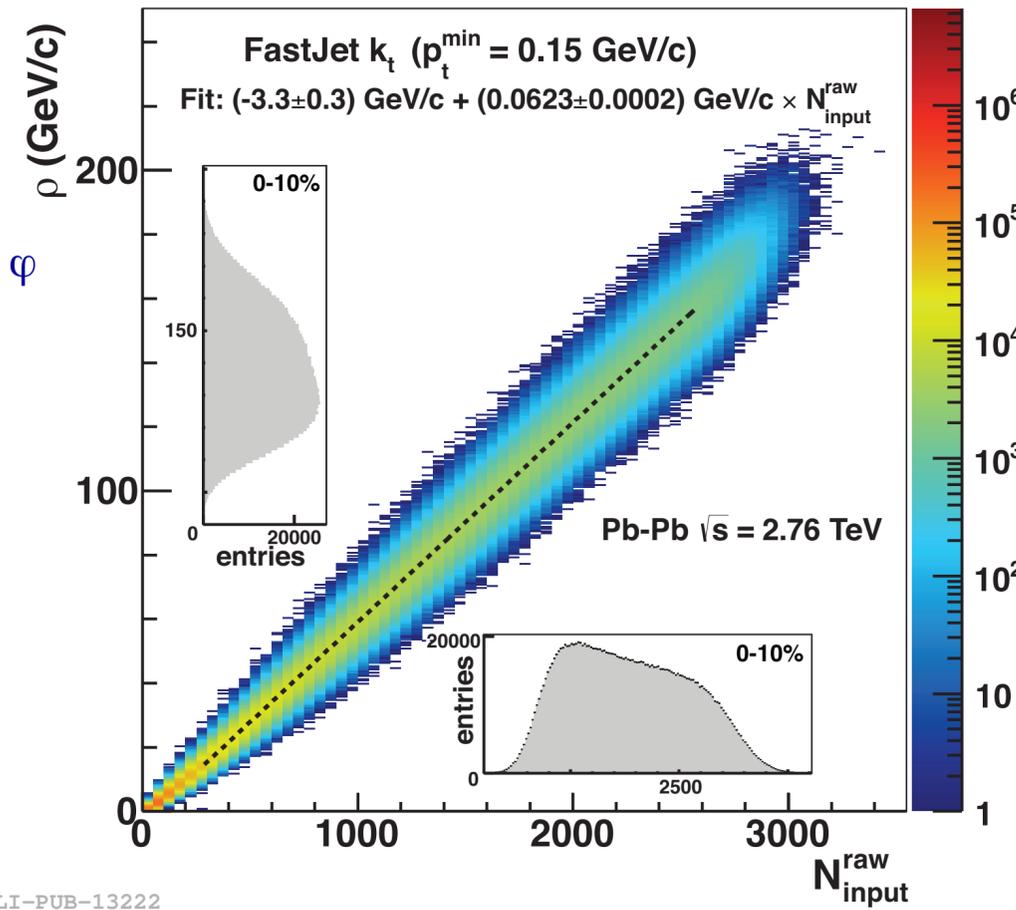
Di-jet energy imbalance: jets lose energy as they propagate through the plasma

Experimental challenge: large combinatorial background in Pb-Pb

Haake and Loizides, PRC99, 064904

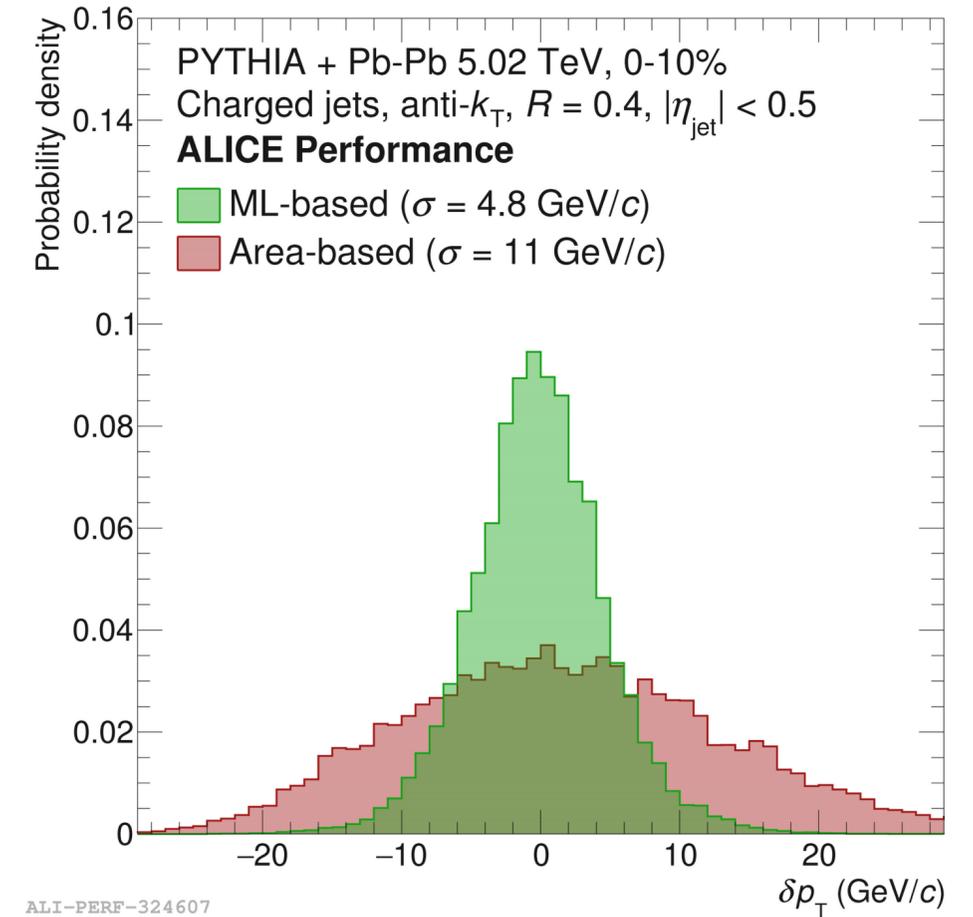


JHEP 03 (2012) 053



ALI-PUB-13222

Machine-learning technique



ALI-PERF-324607

Standard technique: measure background E_T density outside jet 'Area based'

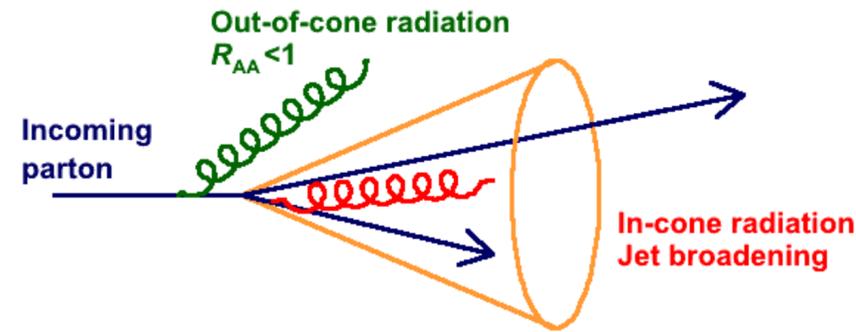
Background subtraction refined based on e.g. leading particle p_T
 Reduces fluctuations

Residual fluctuations due to finite number statistics

Allows to measure jets with larger R , lower p_T

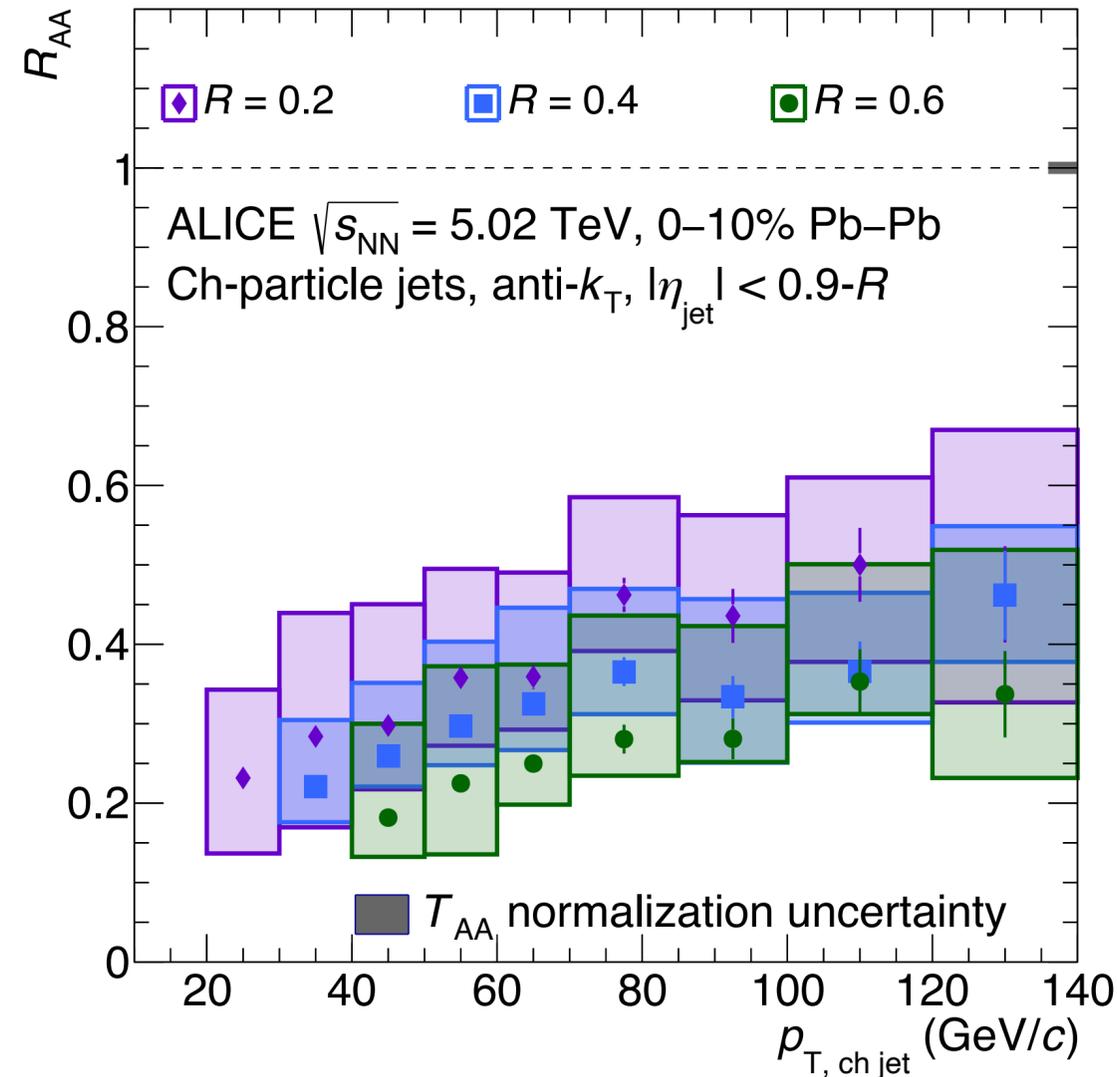
Jet-radius dependence of energy loss in Pb-Pb collisions

ALICE, arXiv:2303.00592

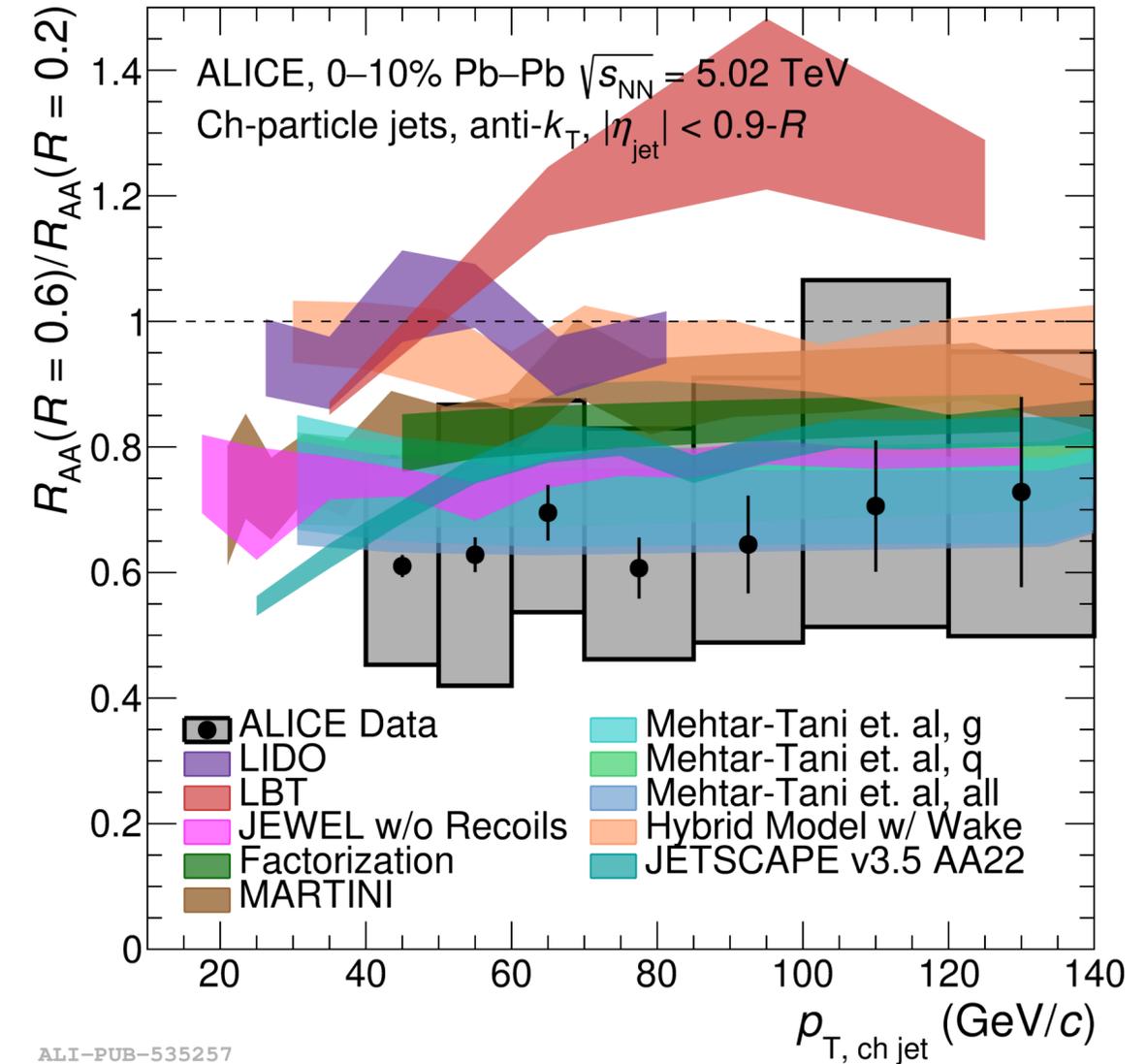


Quantify in- and out-of-cone radiation by measuring jet quenching vs cone radius

Nuclear modification factor



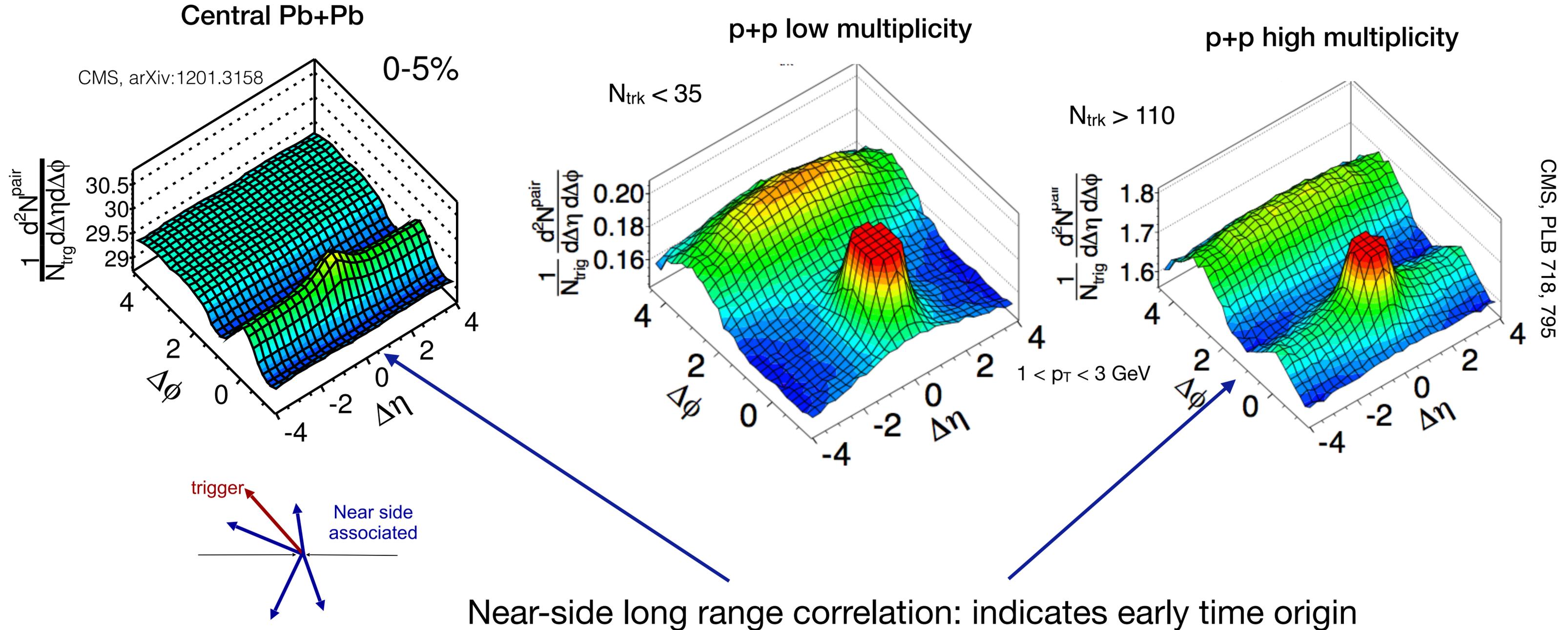
R-dependence: ratio $R=0.6/R=0.2$



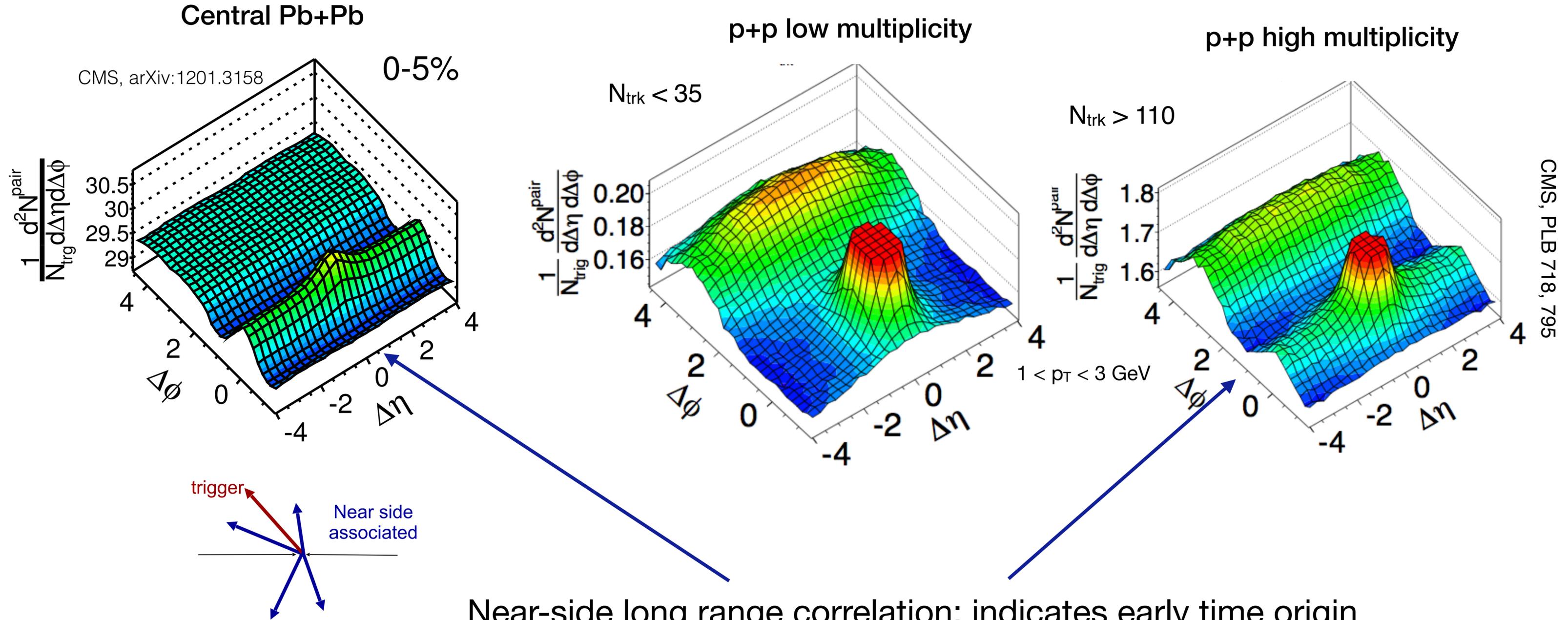
ALI-PUB-535257

- Machine-learning based background subtraction enables jet measurements with R up to 0.6 at $p_T \sim 50$ GeV
- Jet suppression increase with increasing R : wider jets lose more energy

Flow-like effects in pp and Pb-Pb: long range correlations



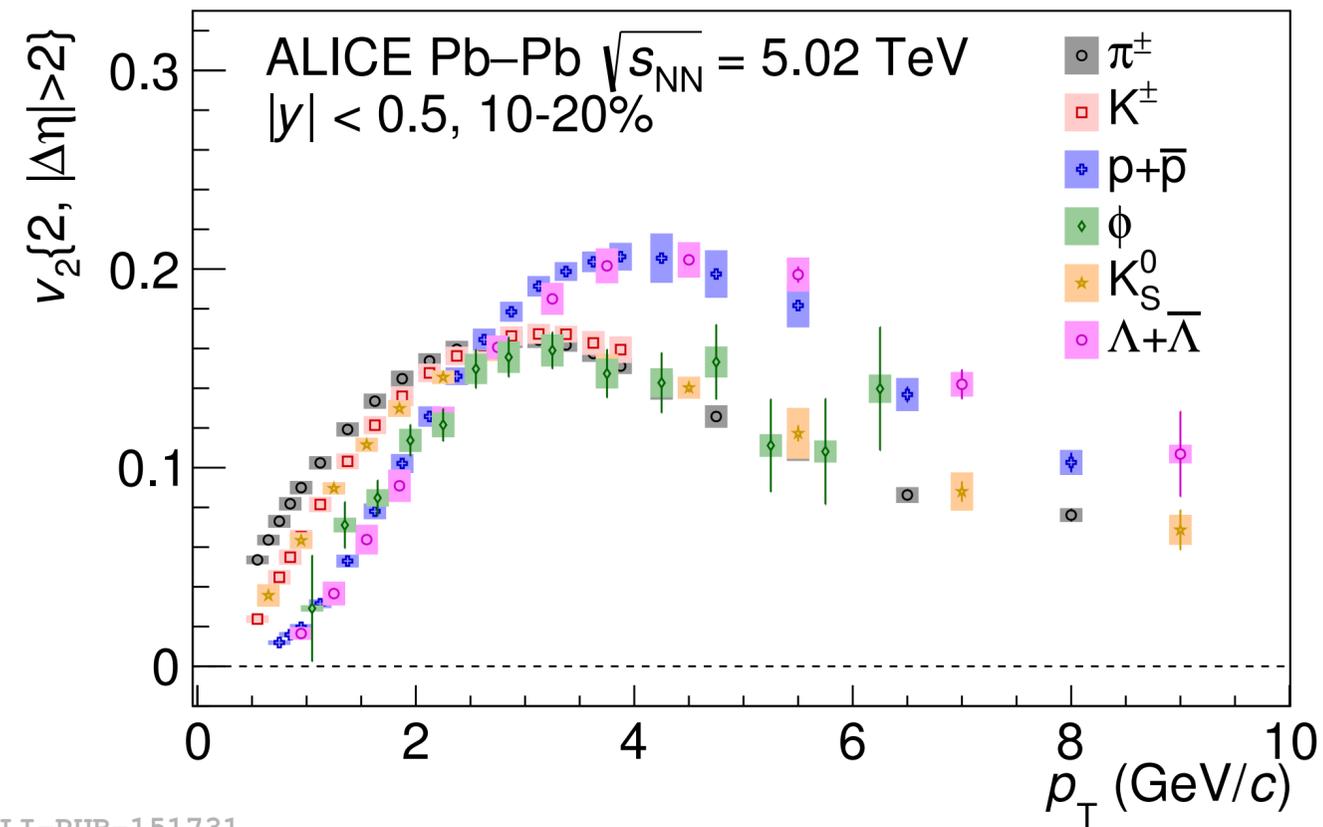
Flow-like effects in pp and Pb-Pb: long range correlations



Near-side long range correlation: indicates early time origin
Seen in high-multiplicity pp and p+Pb events

Elliptic flow in p-Pb

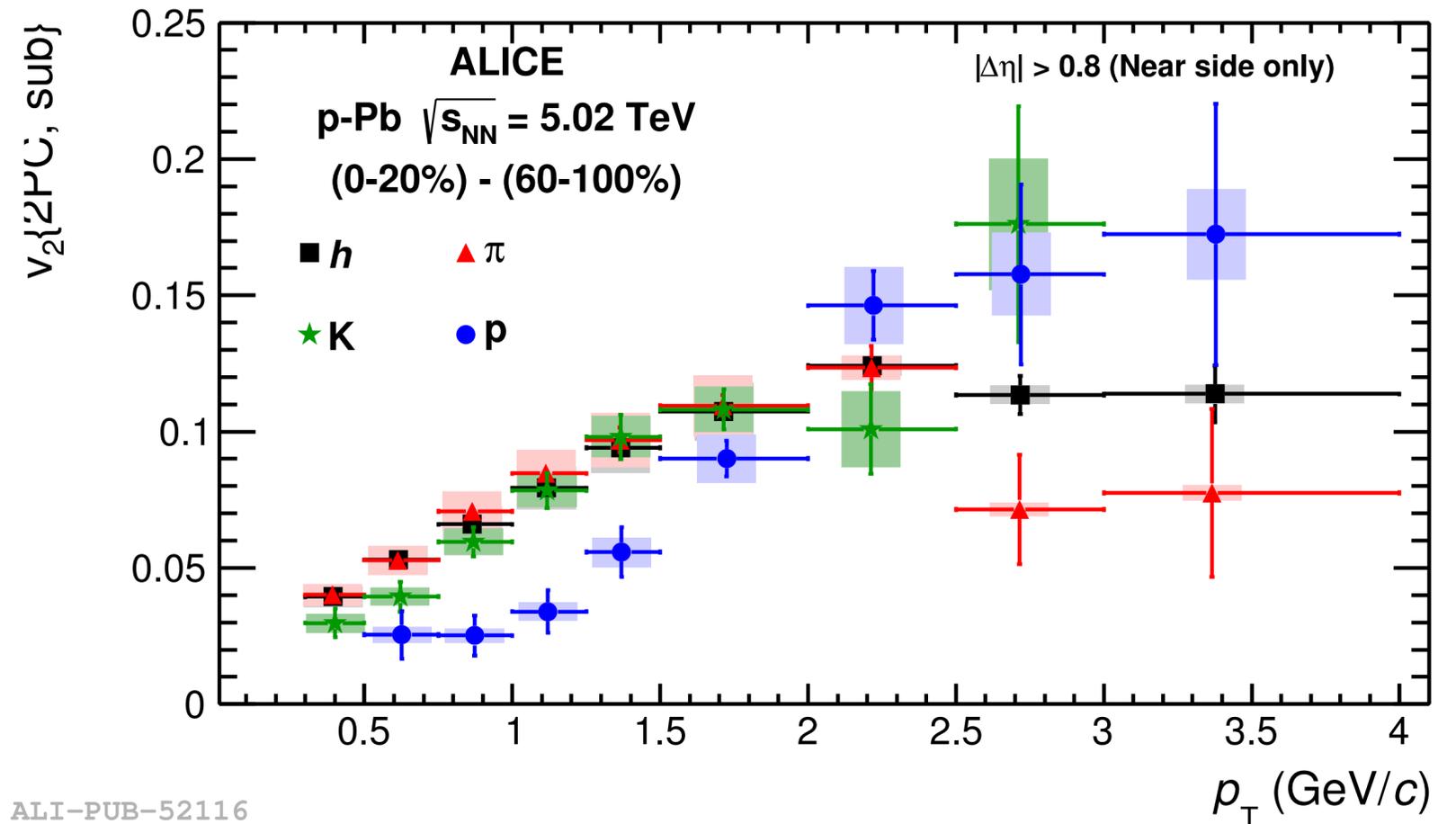
Pb-Pb: elliptic flow v_2



ALI-PUB-151731

Mass-dependence of v_2 measures flow velocity

p-Pb: v_2 from two-particle correlations



ALI-PUB-52116

Similar 'mass ordering' observed for v_2 from two-particle correlations in p+Pb

Is this pressure driven?

Flow effects in small systems

Many aspects of the observed ridge have
a natural explanation in hydrodynamics:

- Long range correlation
- 2- and 3-fold symmetries
- Dependence on initial geometry
- Many-particle correlations
- Particle mass dependence

Why would the system behave as a fluid?
Is there enough time, volume to thermalise?

- Hydrodynamisation (isotropisation) of a dense gluon system?
- Partonic/hadronic rescattering?
- How many scatterings/what density is needed to approximate fluid behaviour?

Limits on hydrodynamic behaviour

Naive expectation: need at least a few collisions for each parton to reach thermal equilibrium and apply hydrodynamic

1) System size: $R > \lambda$

Would not expect azimuthal asymmetries in pp and p-Pb

2) Thermalisation time: $\tau > \frac{\lambda}{v}$

Fits to data: thermalisation times $\tau \approx 0.1-1$ fm/c

Heiselberg and Levy, nucl-th/9812034,
W Lin et al,

pQCD calculation:

$\tau \approx 6.9$ fm/c

Baier et al, PLB 502, 51, PLB 539, 46

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Turns out to be too strict: (viscous) hydro describes non-thermal systems, see next slide

Naive expectations can be 'bypassed' in nature?

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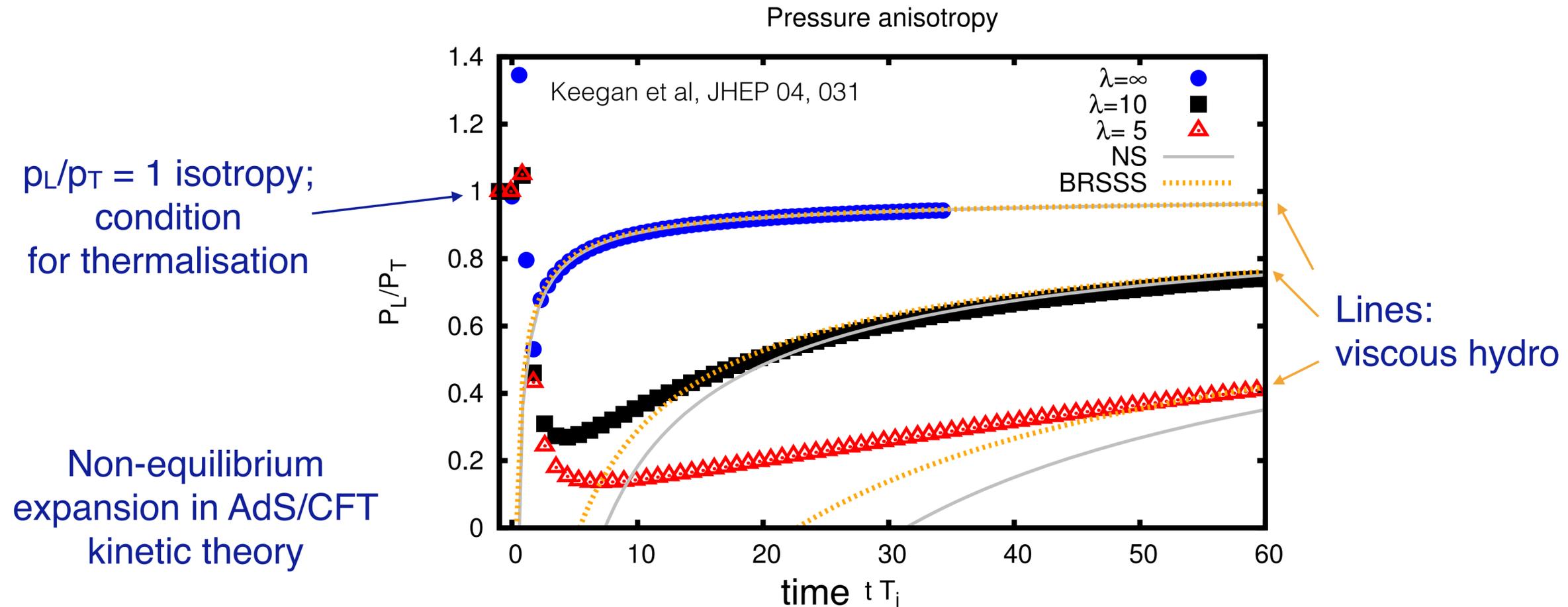
Baier et al, PLB 502, 51, PLB 539, 46

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Naive expectations can be 'bypassed' in nature?

Closely related, since $v \approx c = 1$

Hydrodynamic behaviour in non-thermalised system



Emerging understanding:

Hydrodynamical description valid before thermalisation/isotropisation

Estimate of smallest (possible) system size with fluid behaviour: $r \approx 0.15$ fm

Weller, Romatschke, EPJC 77, 21